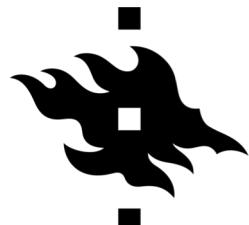


Gravitational Waves from the Early Universe

Chloe Gowling

Supervisors: Mark Hindmarsh and Antony Lewis

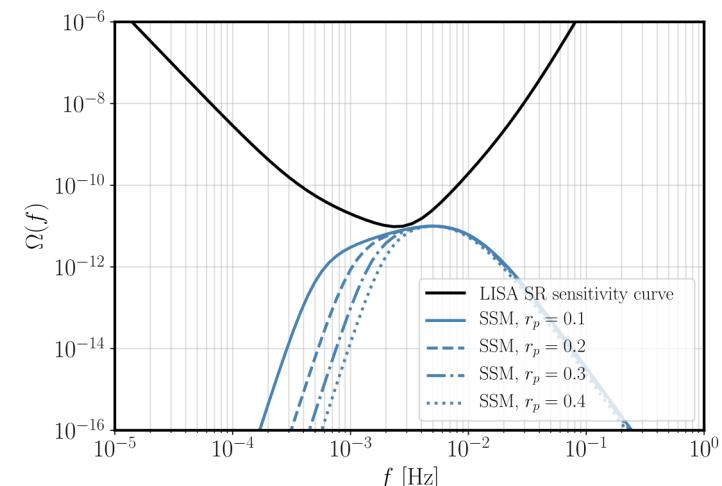
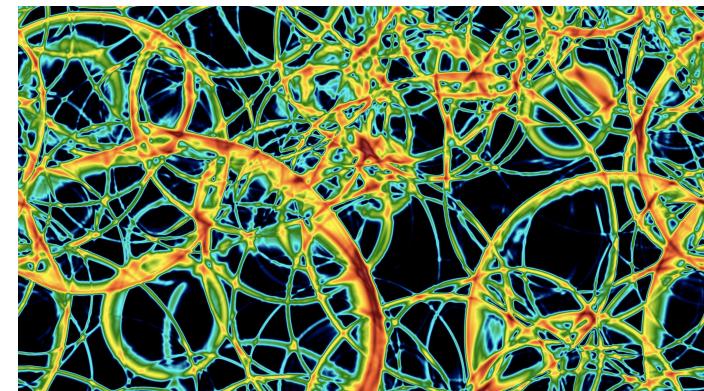
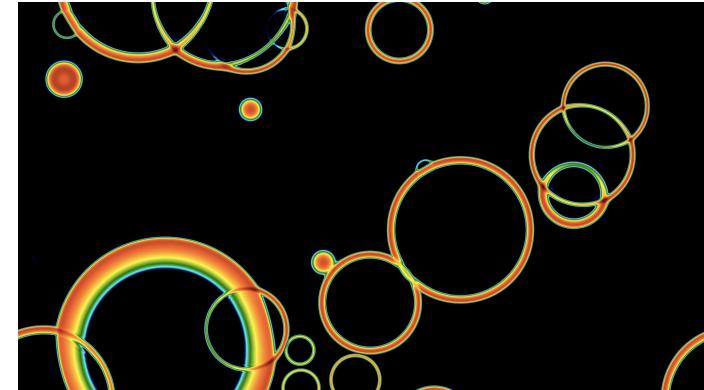


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Overview

- LISA
- Gravitational waves (GWs) from the early universe.
- First order electroweak phase transition.
- The sound shell model.
- Current work
 - Connecting observables to model parameters.
 - How can we extract parameters from a GW power spectrum?



Laser Interferometer Space Antenna

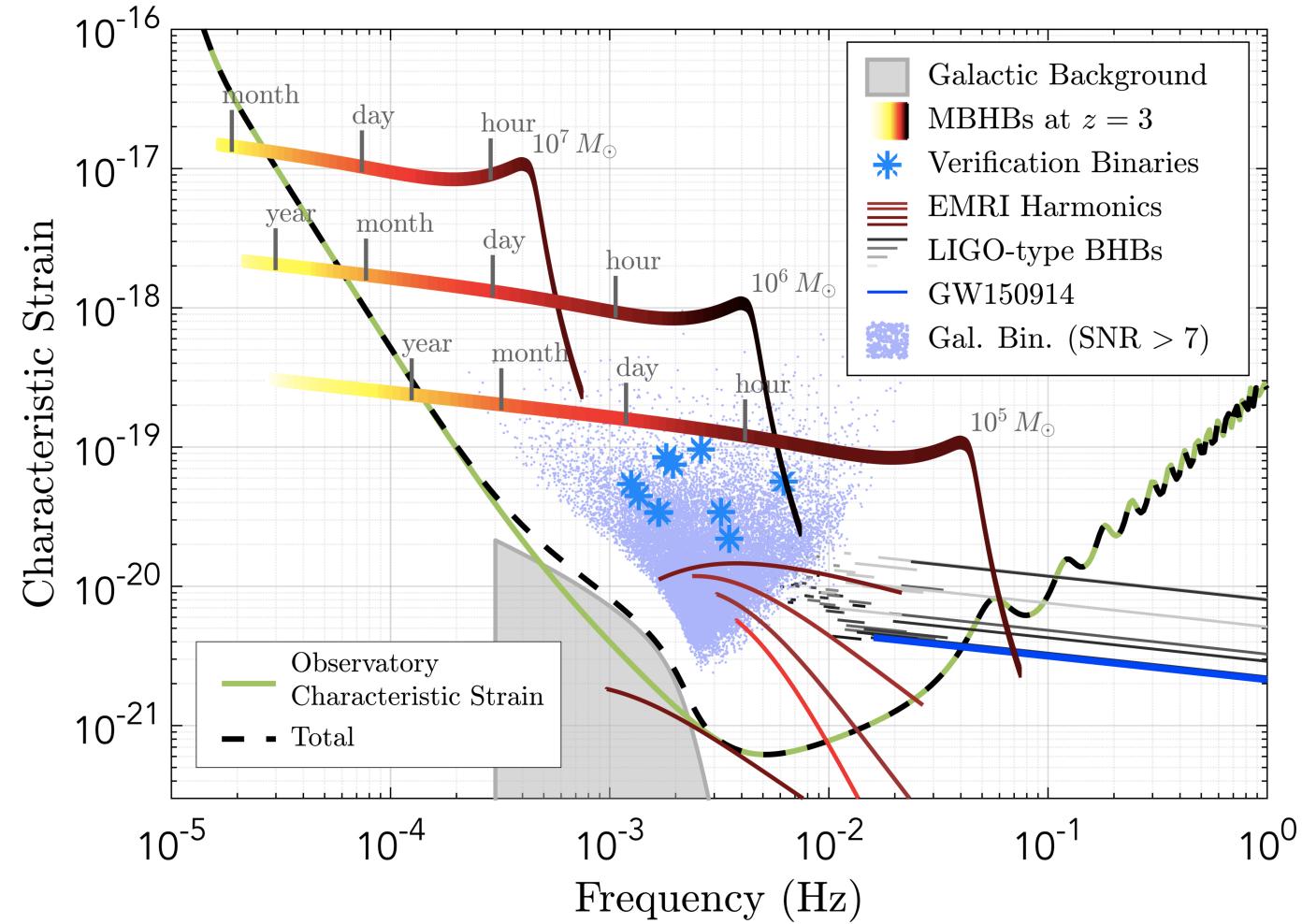
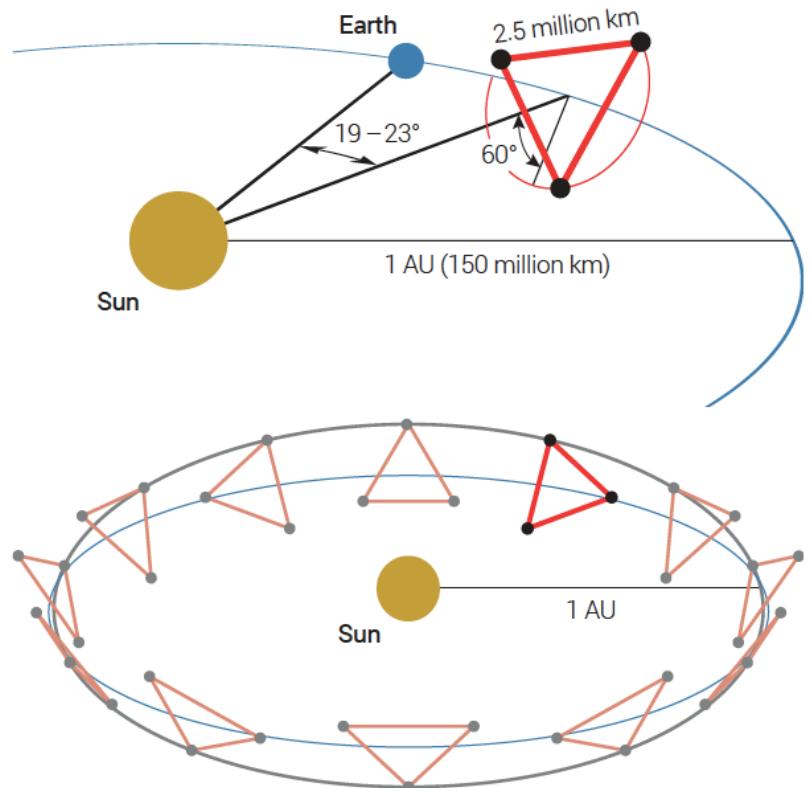
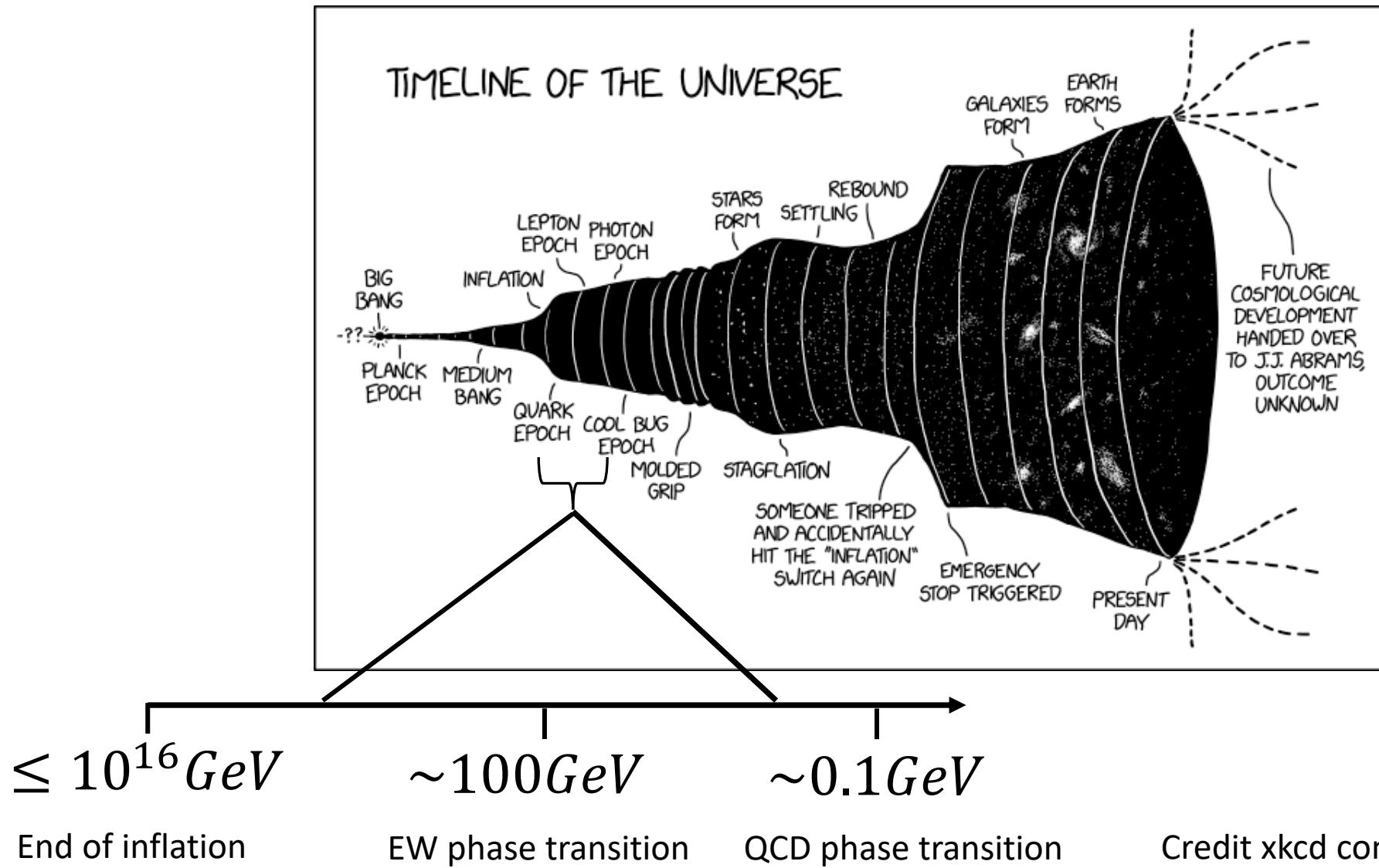


Figure credit: K. Danzmann et al. LISA proposal (2017)

Cosmological phase transitions



Frequency today of GWs from the early universe

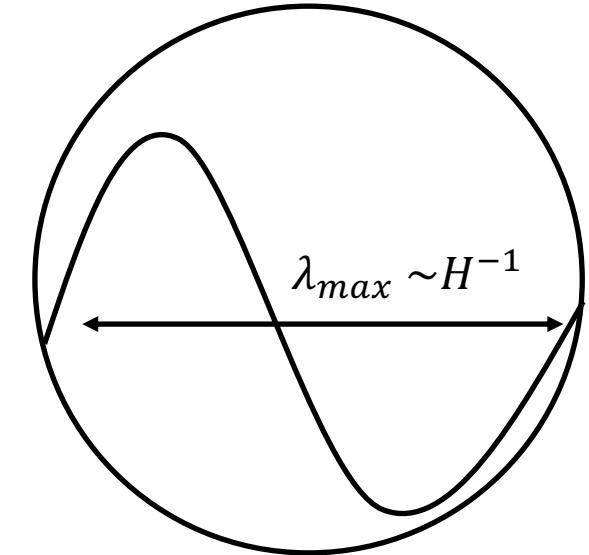
- The minimum frequency of a GW generated at time t

$$f_{min} = H = \frac{1}{t}$$

- Today, taking into account redshift, this corresponds to:

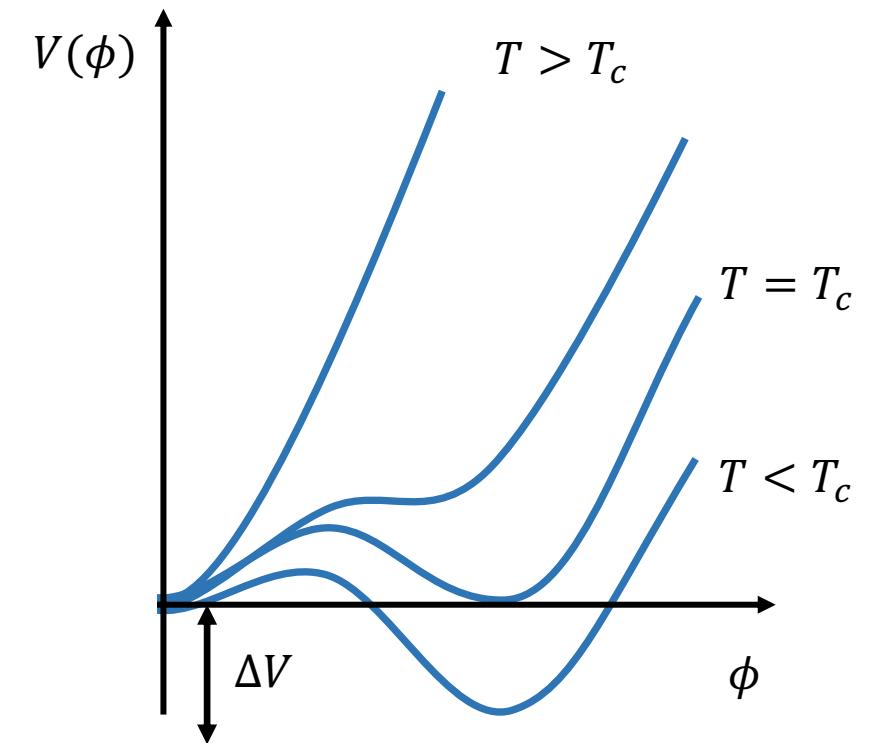
$$f_{min,0} = H_{T,0} = \frac{a}{a_0} H_T$$

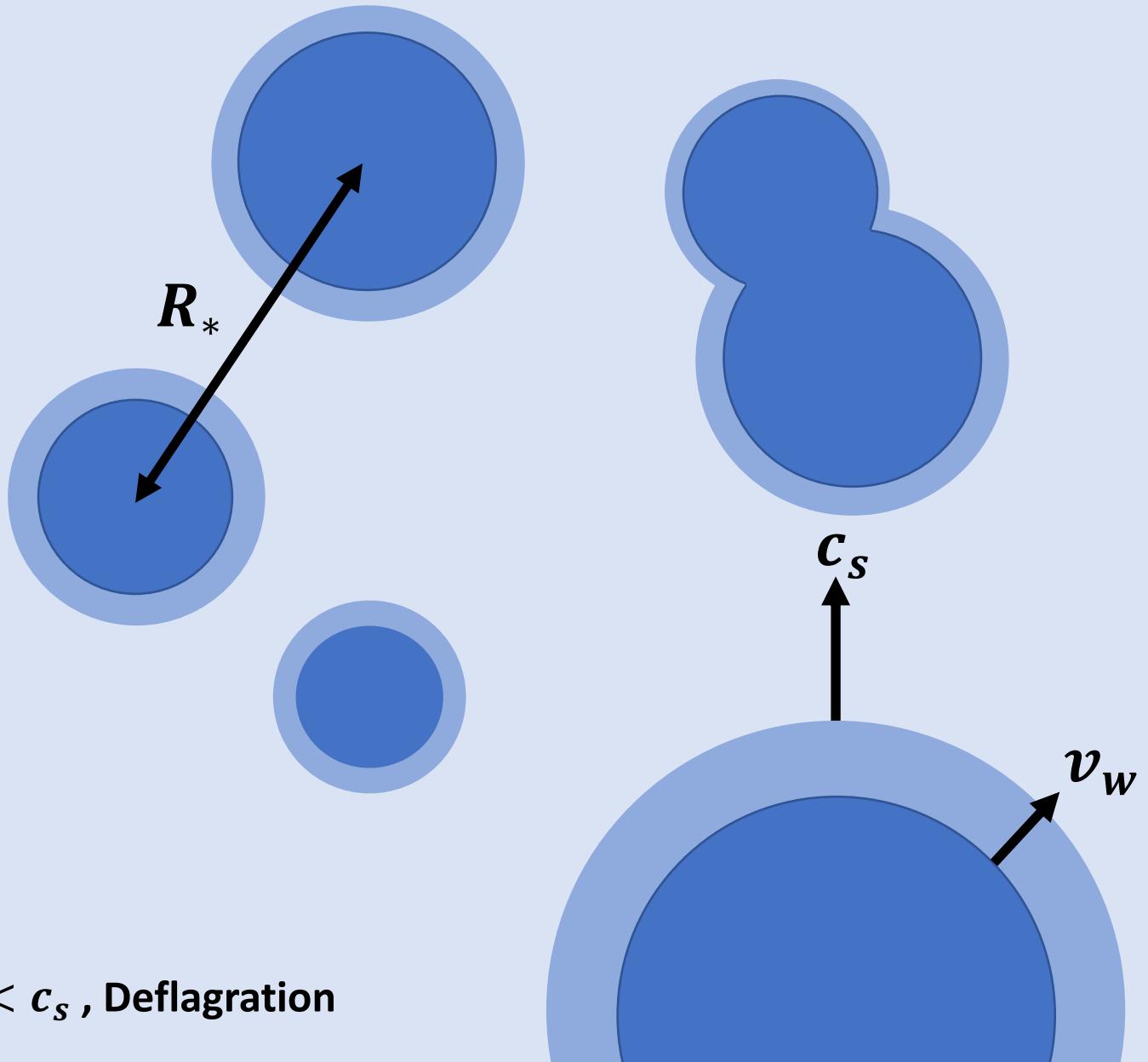
$$f_{min,0} = H_{T,0} = 1.6 \times 10^{-5} \left(\frac{g_*}{100} \right)^{\frac{1}{6}} \left(\frac{T}{100 \text{GeV}} \right) [\text{Hz}]$$



Gravitational waves from the EWPT

- When $T < T_c$ bubbles of broken phase Higgs nucleate.
- Sources of gravitational waves
 - Collision of bubble walls
 - **Sound waves**
 - Other hydrodynamic modes





$v_w < c_s$, Deflagration



Broken phase



Symmetric phase



Sound shell

R_* = mean bubble separation

v_w = wall velocity

c_s = speed of sound

Potential energy in the scalar field is converted into **kinetic energy** of the plasma.

Thermodynamic parameters

- Transition strength - $\alpha \sim \frac{\Delta V}{\rho_{th}}$
- Wall speed- v_w
- Transition rate- β
- Nucleation temperature - T_n

Derived parameters

- Mean bubble separation - $R_* = (8\pi)^{\frac{1}{3}} \frac{v_w}{\beta}$
- Sound shell thickness $\sim |v_w - c_s|$
- Kinetic energy fraction - $K^2(v_w, \alpha)$

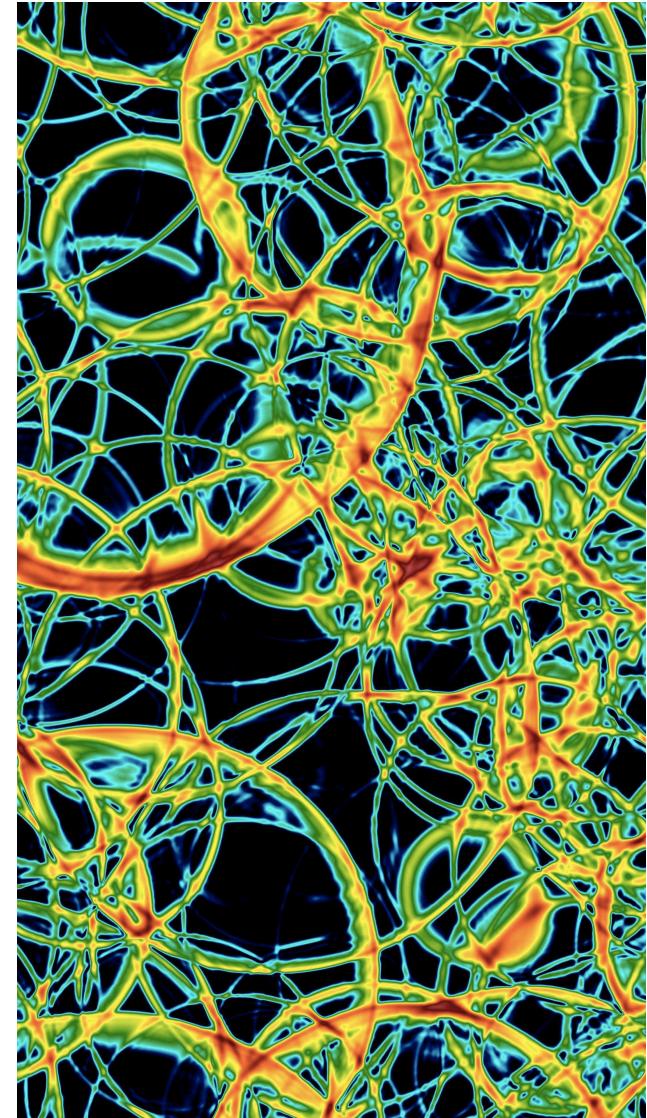
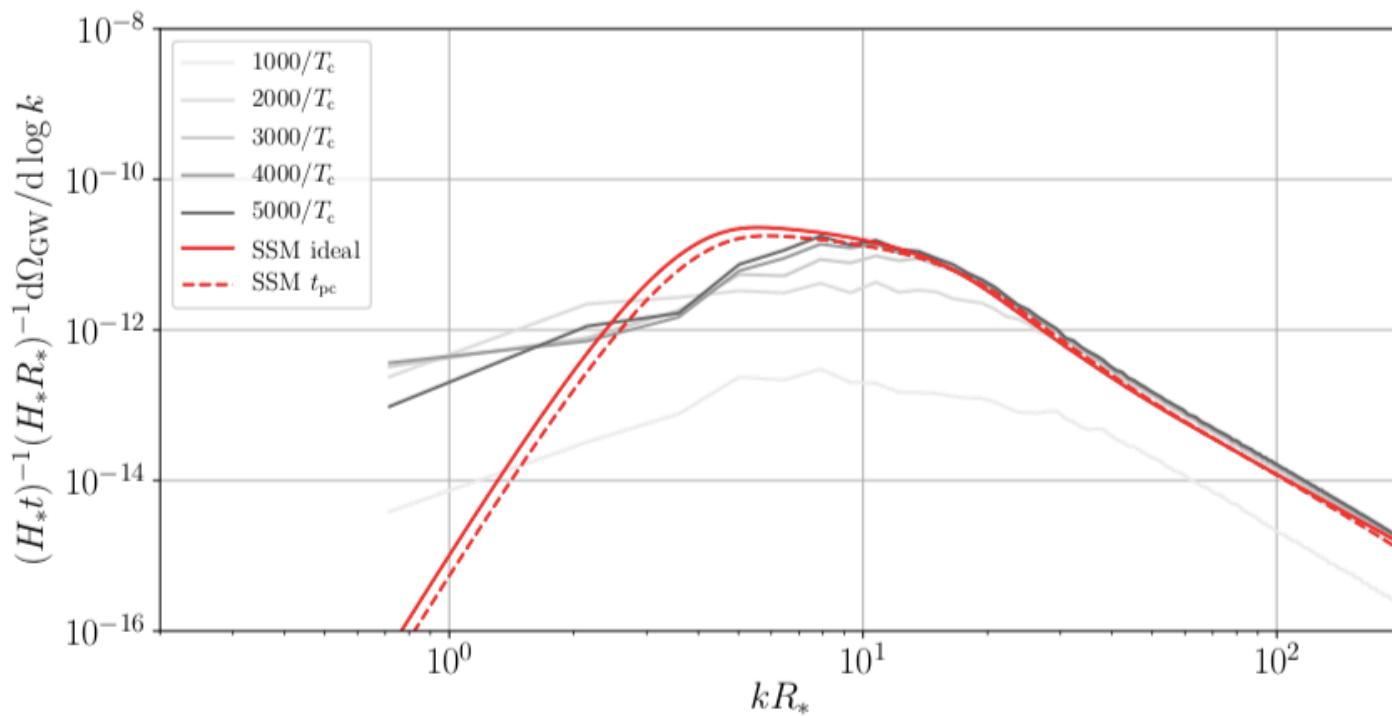


Image credit: David Weir, [arXiv:1705.01783](https://arxiv.org/abs/1705.01783)

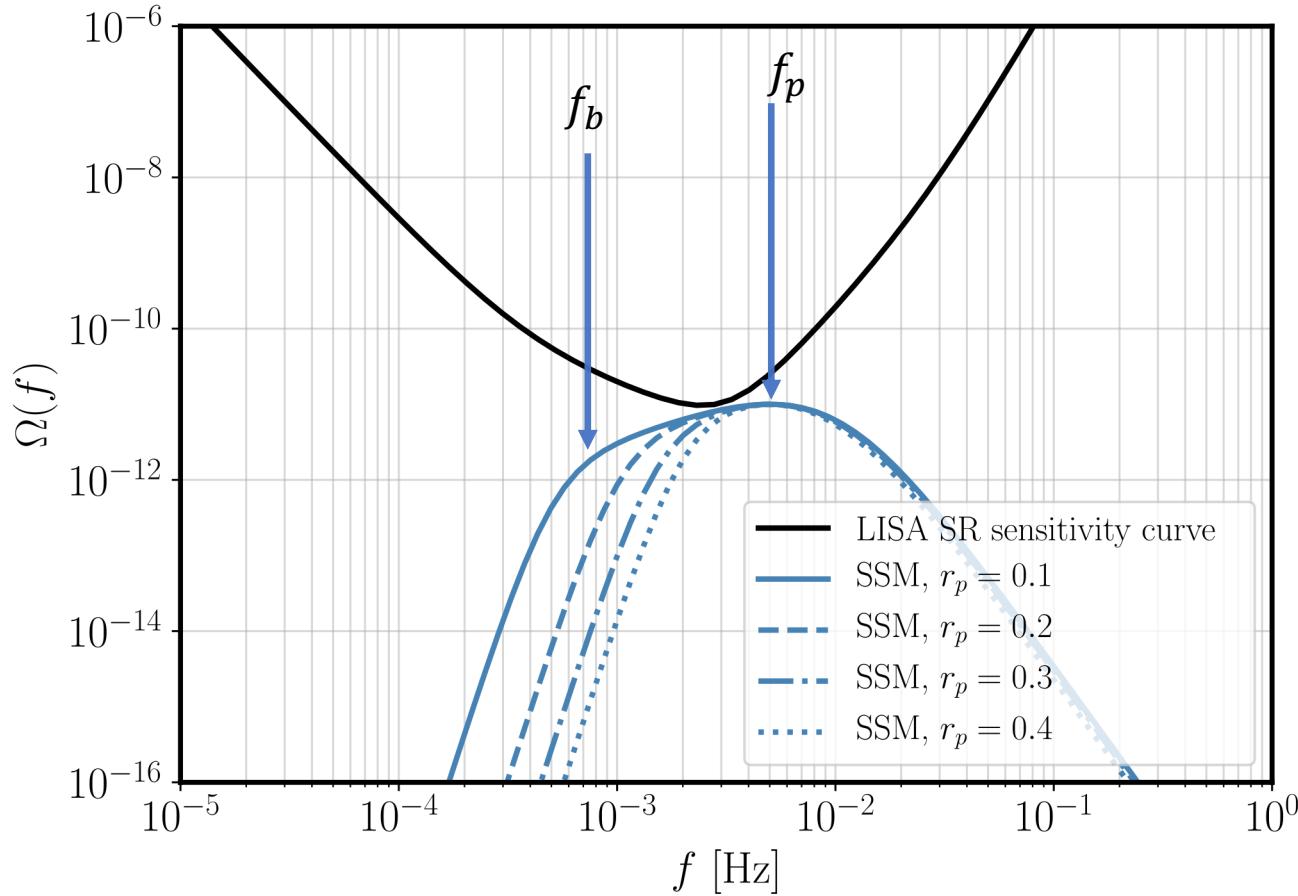
Sound shell model



- GW power spectrum can be computed from the velocity power spectrum of the fluid.
- SSM proposes the velocity PS can be calculated from the sound shells.
- Analytic approximation of simulation results is shown in red.

Figure credit: M.Hindmarsh, slides from Nordita workshop 2019

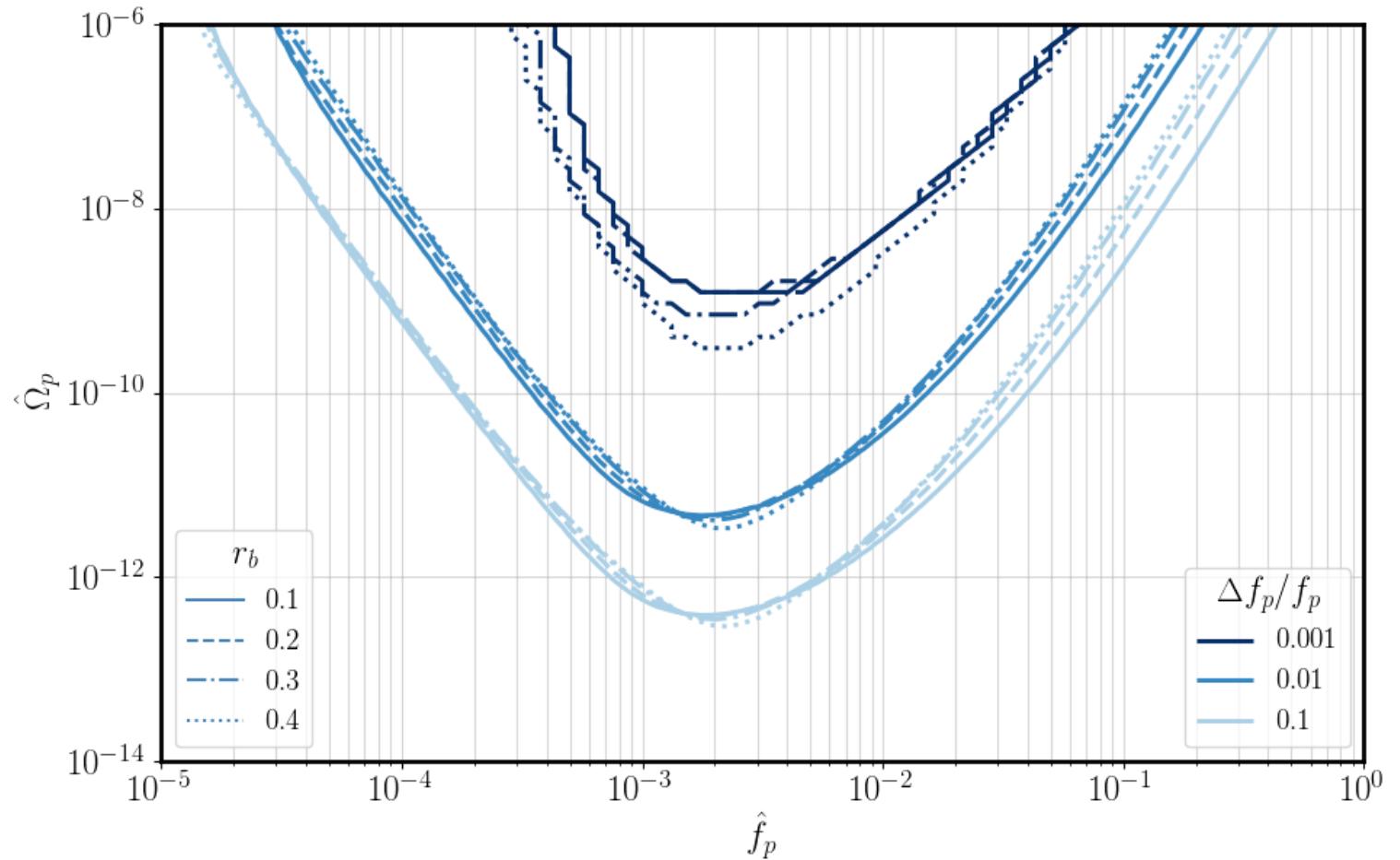
Simplified sound shell model



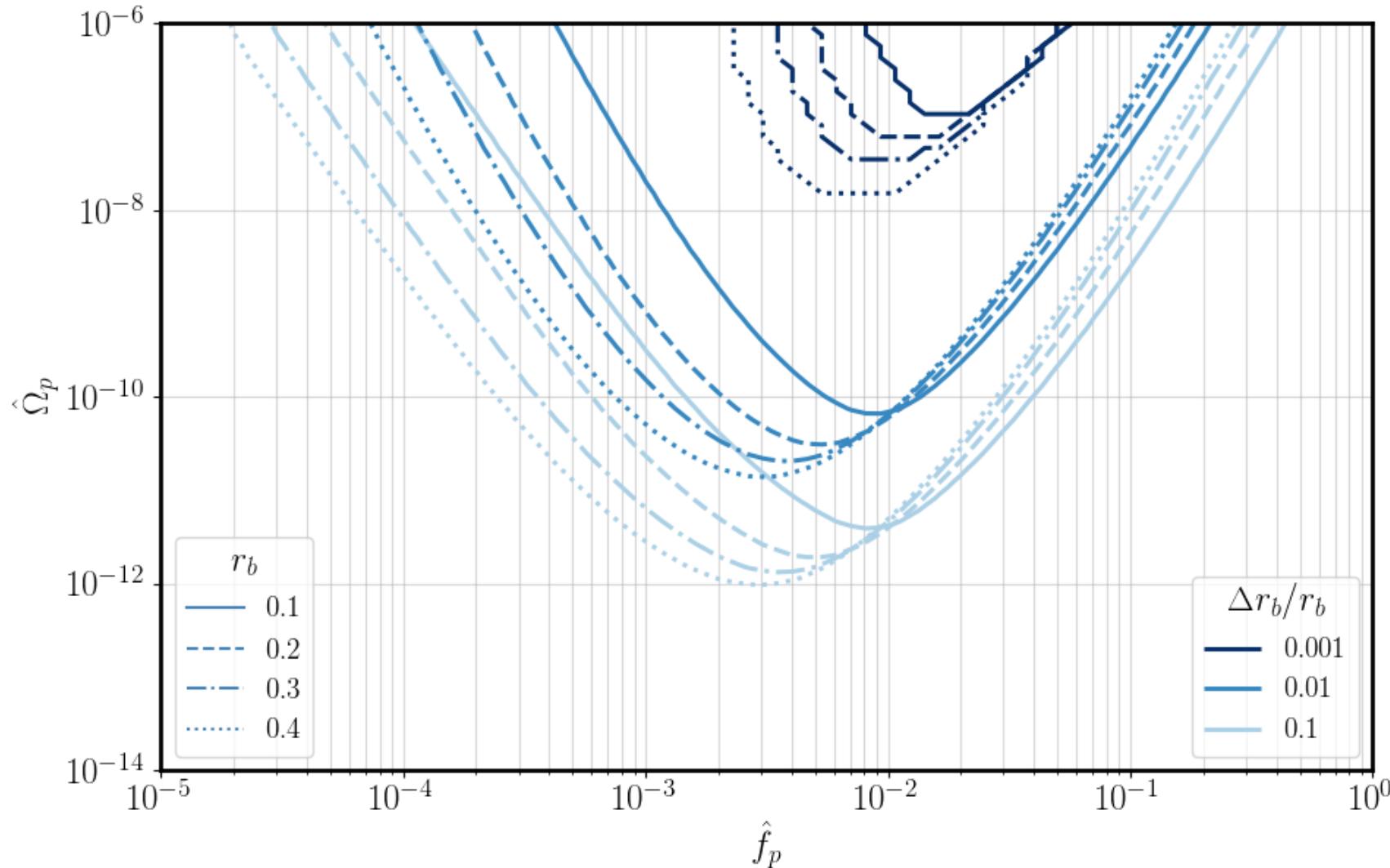
- $\Omega_{GW}(f) = \Omega_p M(s, r_b)$
- $M = s^9 \left(\frac{1 + r_b^4}{r_b^4 + s^4} \right)^2 \left(\frac{5}{5 - m + ms^2} \right)^{\frac{5}{2}}$
- $m = \frac{1 + 9 r_b^4}{1 + r_b^4}, \quad s = \frac{f}{f_p}, \quad r_b = \frac{f_b}{f_p}$
- Observable parameters Ω_p, f_p and r_b

Sensitivity forecasts: f_p

- $\chi^2 = T_{obs} \int \frac{(\Omega_{GW}(f, \theta) - \Omega_n(f, \hat{\theta}))}{\Omega_n^2}$
- $\mathcal{L} = e^{-\frac{1}{2}\chi^2}$
- $F_{ij} = \frac{\partial^2 \ln(\mathcal{L})}{\partial \theta_i \partial \theta_j} = \frac{1}{2} \frac{\partial^2 \chi^2}{\partial \theta_i \partial \theta_j}$
- $\frac{\Delta f_p}{f_p}$ relative uncertainty in f_p



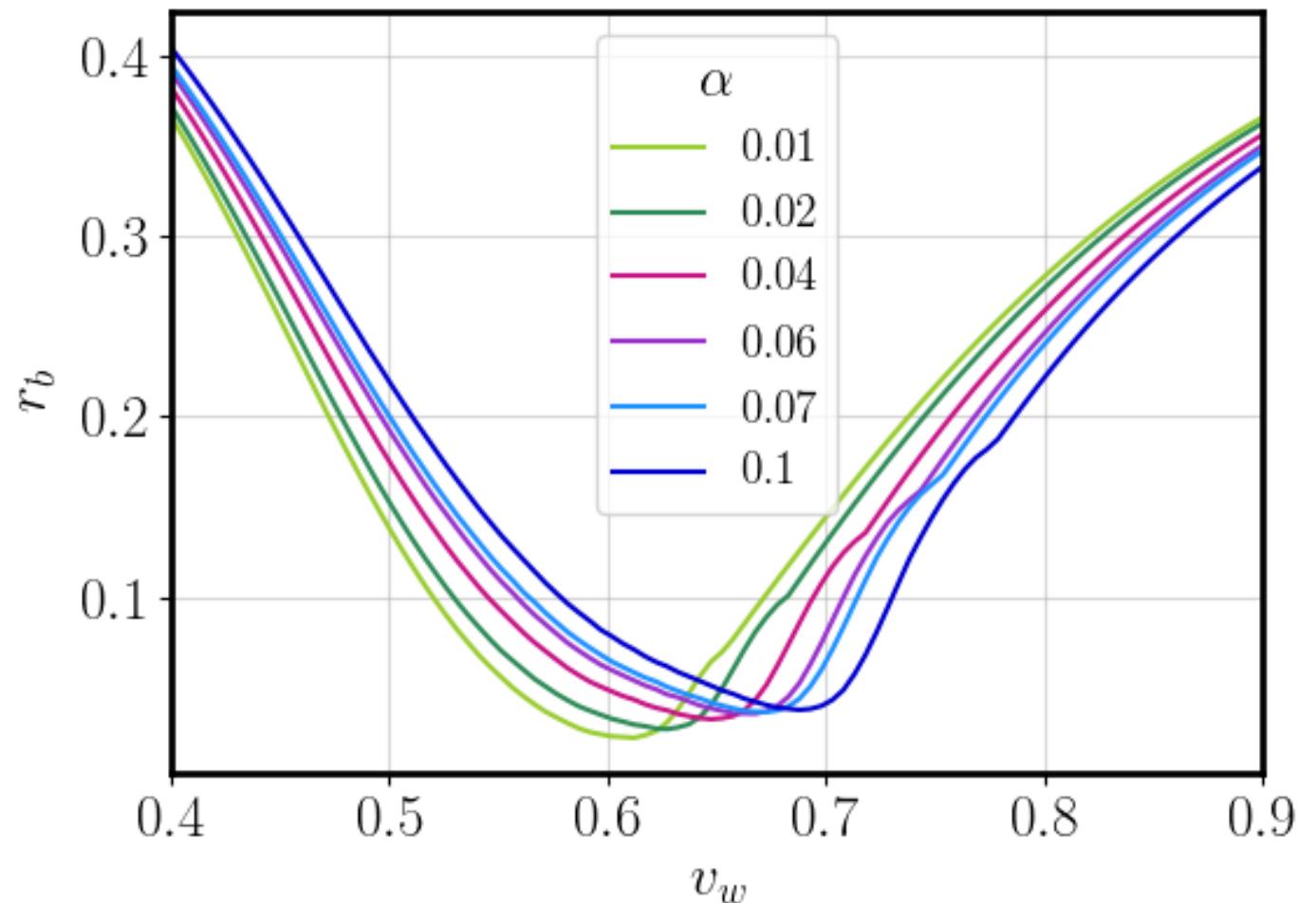
Sensitivity forecasts: r_b



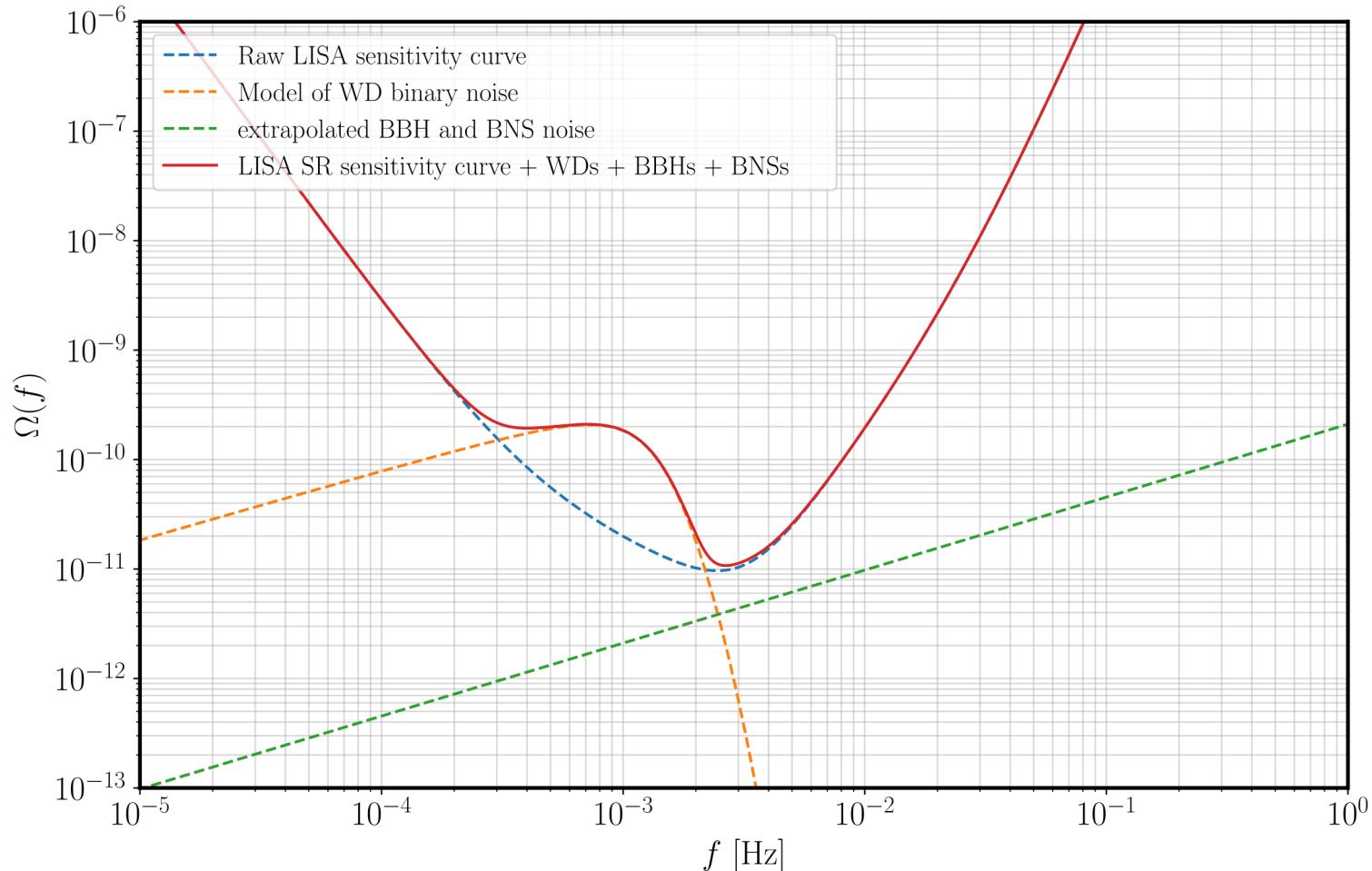
- $r_b = \frac{f_b}{f_p}, \quad 0 < r_b < 1$
- Larger r_b corresponds to a narrower power spectrum.

Connection to thermodynamic parameters

- $\Omega_p(v_w, \alpha, H_n R_*)$
- $f_p(v_w, \alpha, H_n R_*)$
- $r_b(v_w, \alpha)$
- fixed T_n



Astrophysical noise



White dwarf binaries : arXiv:1703.09858v2, Compact binaries: PhysRevLett.120.091101

Next steps

- Better understanding of the sources of noise
 - Astrophysical noise, white dwarfs, BBHs and BNSs
- Further investigate the relationship between (Ω_p, f_p, r_p) and thermodynamic parameters $(v_w, \alpha, H_n R_*, T_n)$.
- Perform Monte Carlo analysis on for thermodynamic parameters.