

Resurgence in the Bi-Yang-Baxter Model

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Previous work by Demulder, Dorigoni, Thompson (2016) and Cherman, Dorigoni, Dunne, Ünsal (2014).

- 1 Yang-Baxter deformations
- 2 Uniton solutions
- 3 Dimensional Reduction to Schrödinger System
- 4 Borel Analysis
- 5 Seiberg-Witten Curve

The Principal Chiral Model (PCM)

- Non-linear, **integrable** (i.e. there exists a flat Lax connection) σ -model on the world-sheet.
- Fields are maps from the worldsheet into a Lie group:
 $g : \Sigma \rightarrow G = SU(N)$.

- Action:

$$S = \frac{1}{2\pi t} \int dz^2 \text{Tr} (g^{-1} \partial_+ g g^{-1} \partial_- g) .$$

- $z_{\pm} = t \pm ix$ light-cone coordinates on the worldsheet $\partial_{\pm} = \partial_{z_{\pm}}$.
- $\partial_+ g \in \mathfrak{g}$ Lie algebra valued.
- $G^L \times G^R$ global symmetry.

Yang-Baxter Deformations

- Let R be the Yang-Baxter (YB) Operator on \mathfrak{g} that satisfies

$$[RA, RB] - R([RA, B] + [A, RB]) = [A, B].$$

- On $SU(2)$, in the basis $\sigma_i/\sqrt{2}$:

$$R = \begin{pmatrix} 0 & -1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}.$$

- bi-Yang-Baxter (bYB) deformed PCM:

$$S = \frac{1}{2\pi t} \int dz^2 \text{Tr} \left(g^{-1} \partial_+ g \frac{1}{1 - \eta R - \zeta R^g} g^{-1} \partial_- g \right).$$

- $\text{Ad}_g(u) = gug^{-1}$, $R^g = \text{Ad}_{g^{-1}} \circ R \circ \text{Ad}_g$.
- $\eta, \zeta \in \mathbb{R}$ parameters that squish $S^3 \simeq SU(2)$.
- **preserves integrability** (Klimcik, 2014).

- Parametrise $g \in SU(2)$ by Hopf coordinates:

$$g = \begin{pmatrix} \cos \theta e^{i\phi_1} & i \sin \theta e^{i\phi_2} \\ i \sin \theta e^{-i\phi_2} & \cos \theta e^{-i\phi_1} \end{pmatrix}.$$

- There exists a solution to the equation of motions in terms of a holomorphic function $f(z)$:

$$\phi_1 = \frac{\pi}{2}, \quad \phi_2 = \pi + \frac{i}{2} \log \left(\frac{f}{\bar{f}} \right),$$
$$\cos \theta(|f|^2) = \frac{\sqrt{1 + (\eta - \zeta)^2(|f|^2 - 1)}}{\sqrt{(1 + |f|^2)^2 + (\eta - \zeta)^2(1 - |f|^2)^2}}.$$

- “Uniton”.

Unitons (continued)

- Finite, quantised action configuration:

$$S = \frac{2k}{t(1 + (\zeta + \eta)^2)} S_I,$$

$$S_I = \frac{(1 + (\zeta + \eta)^2)}{2\zeta\eta} [(\zeta - \eta)\arctan(\zeta - \eta) - (\zeta + \eta)\arctan(\zeta + \eta)].$$

- $k \in \mathbb{R}$ order of the polynomial $f(z)$.

→ Analogous complex uniton.

Dimensional Reduction

- Compactify along the x -direction by a twisted boundary condition

$$g(t, x + L) = e^{iH_L} g(t, x) e^{-iH_R}, \quad H_L = H_R = \begin{pmatrix} \frac{\pi}{2} & 0 \\ 0 & \frac{\pi}{2} \end{pmatrix}.$$

→ Integrate out heavy Kaluza-Klein modes.

→ Obtain an effective Quantum Mechanics for θ .

Schrödinger System

- Schrödinger Equation for θ :

$$H = \frac{g^2}{4} p_\theta^2 + \frac{1}{g^2} V(\theta), \quad g = t(1 + (\zeta + \eta)^2),$$

with

$$V(\theta) = \frac{\pi^2 \operatorname{sn}^2(\theta)(1 + (\zeta - \eta)^2 \operatorname{sn}^2(\theta))}{\operatorname{dn}^2(\theta)}, \quad m = \frac{4\zeta\eta}{1 + (\zeta + \eta)^2}.$$

- $\operatorname{sn}(z)$, $\operatorname{cn}(z)$ and $\operatorname{dn}(z)$ are the **Jacobi elliptic functions** depending on the elliptic modulus m .
- $\zeta = 0 \implies m = 0 \implies \operatorname{sn}(z) = \sin(z), \operatorname{dn}(z) = 1$.

- Use the WKB method to obtain a perturbative expansion for the ground state energy:

$$E(g^2) = \sum_{n=0}^{\infty} a_n g^{2n} \quad a_n = a_n(\zeta, \eta).$$

- This series is divergent: $a_n \sim n!$. **asymptotic expansion.**
- Borel transform:

$$B[E](g^2) = \sum_{n=0}^{\infty} \frac{a_n}{n!} g^{2n}.$$

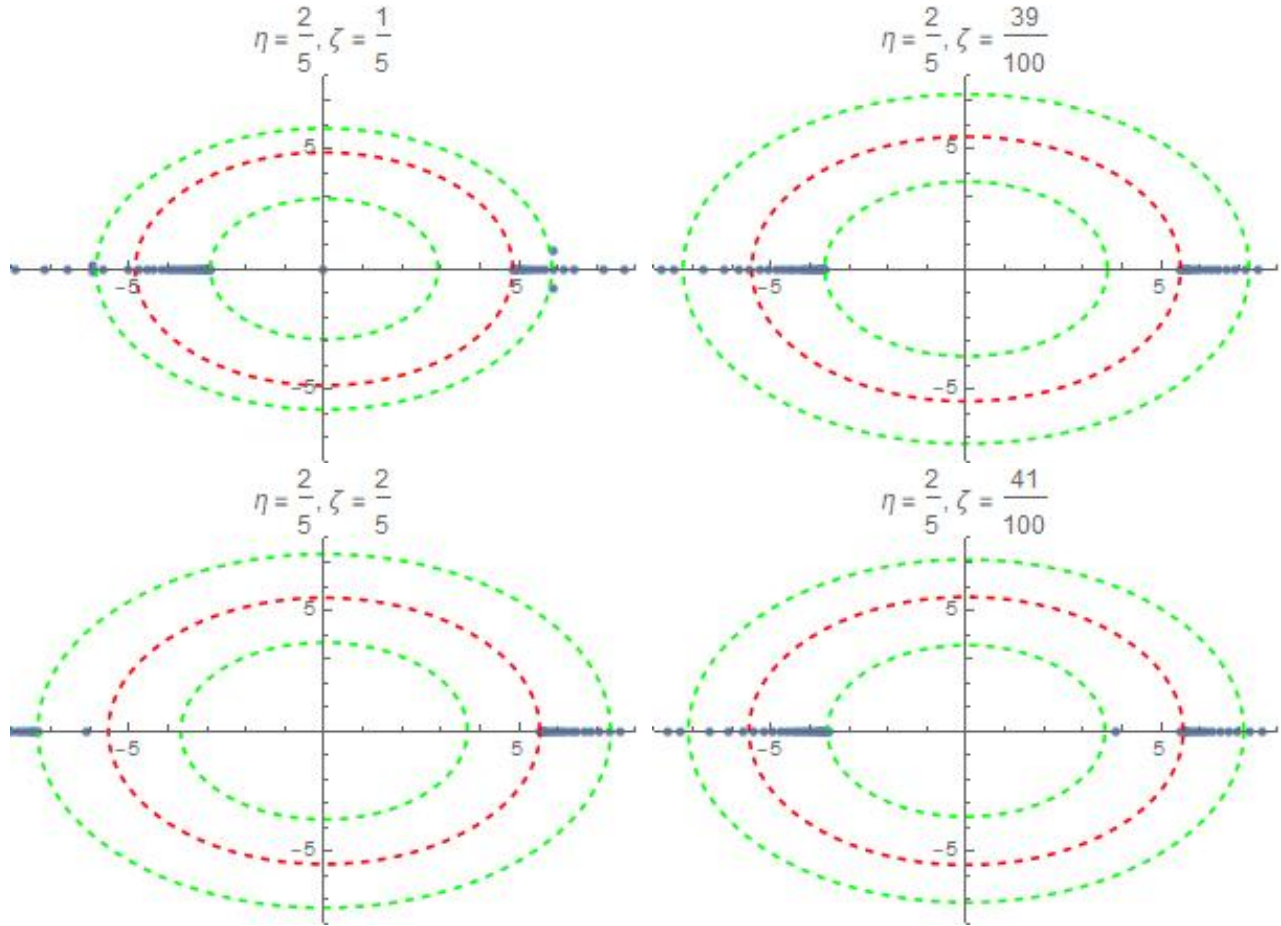
- Borel resummation

$$S[E](g^2) = \int_0^{\infty} dz B[E](g^2) e^{-z/g^2} \sim E(g^2).$$

Borel Analysis (Continued)

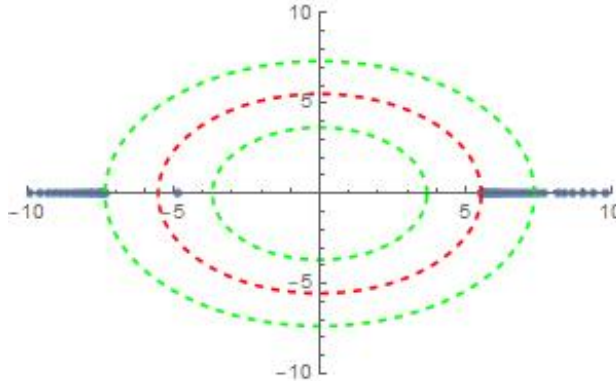
- Poles in the Borel plane obstruct this integral along certain rays and are due to **non-perturbative** effects: Unitons!
- Non-perturbative effects go as e^{-2S_i/g^2} .
- Use Bender-Wu package (Sulejmanpasic and Ünsal, 2016) to obtain 150-200 orders of perturbation theory (~ 10 minutes per plot).
- Use Padé approximant to find poles and accumulation points.

Poles in the Borel Plane (1)

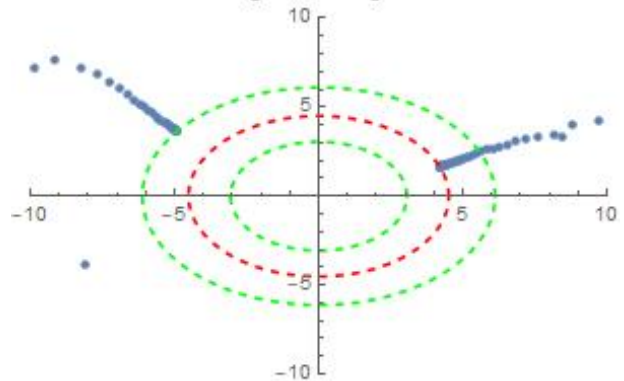


Poles in the Borel Plane (2)

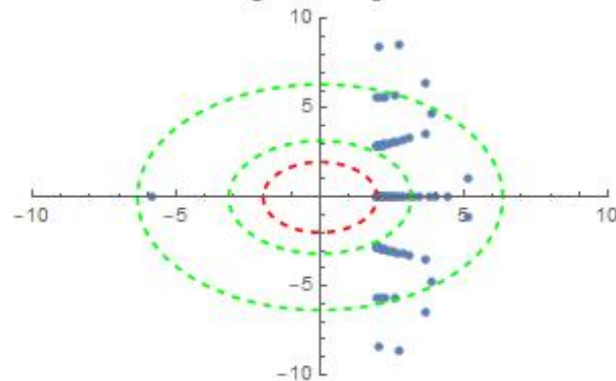
$$\eta = \frac{2}{5}e^{i0}, \zeta = \frac{2}{5}e^{i0}$$



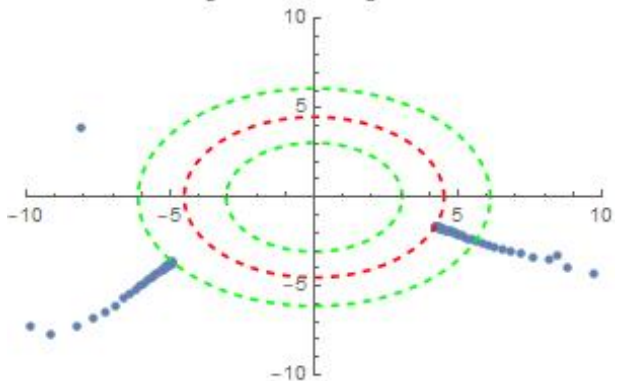
$$\eta = \frac{2}{5}e^{\frac{i\pi}{4}}, \zeta = \frac{2}{5}e^{\frac{i\pi}{4}}$$



$$\eta = \frac{2}{5}e^{\frac{i\pi}{2}}, \zeta = \frac{2}{5}e^{\frac{i\pi}{2}}$$



$$\eta = \frac{2}{5}e^{\frac{1}{2}i(3\pi)}, \zeta = \frac{2}{5}e^{\frac{1}{2}i(3\pi)}$$



Seiberg-Witten (SW) Curve

- Consider the quadratic differential of the low-energy effective theory of $\mathcal{N} = 2$ Seiberg-Witten theory.
- $SU(2)$ with $N_f = 4$ and pairwise equal masses.
- Trigonometric limit of $SU(2) \times SU(2)$ elliptic quiver theory with $N_f = 2$.
- Coincides with our WKB-differential!
- BPS states = Unitons?
- Connection with integrable Hitchin system through Gaudin Models??

Questions