

A Bogomol'nyi equation for Magnetic Skyrmions

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What are topological solitons?

What are they not?

The Bogomol'nyi argument

Chiral magnets and magnetic skyrmions

The bogomol'nyi argument for magnetic skyrmions

Solving the bogomol'nyi equation for magnetic skyrmions

Conclusions

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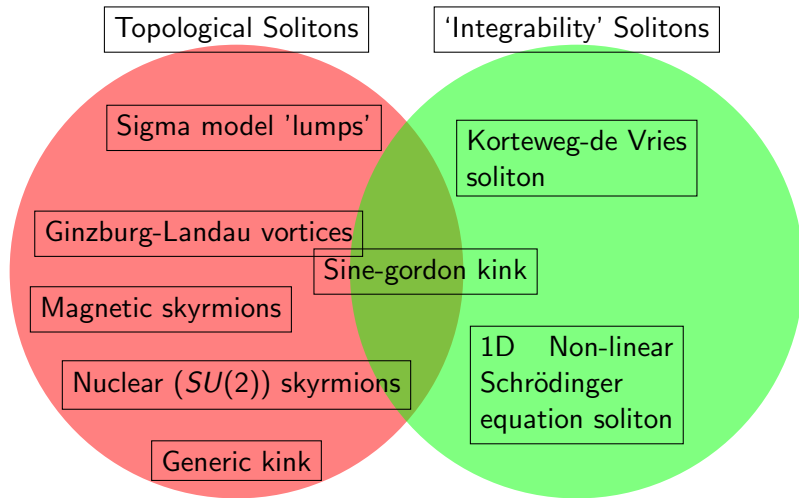
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- ▶ These are 'conserved charges' of the theory- no corresponding symmetry!

What are they not? A side-note on bad naming



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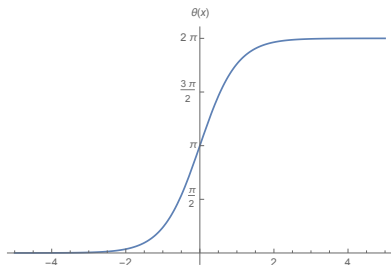
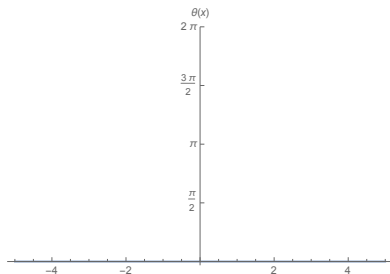
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 - ▶ It gives us a lower bound on the energy, related to the topology
 - ▶ It gives us an equation to solve to find minimizers- which is in general easier to solve than the Euler-Lagrange equation

Example: the 1D kink

We have the static energy functional for a S^1 -valued field $\theta(x)$:

$$E = \int \left\{ \frac{1}{2} \theta'(x)^2 + V(\theta) \right\} dx$$



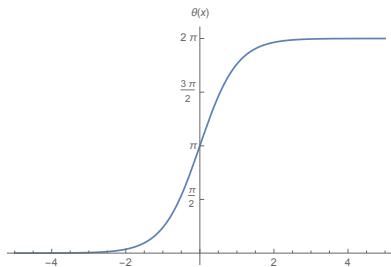
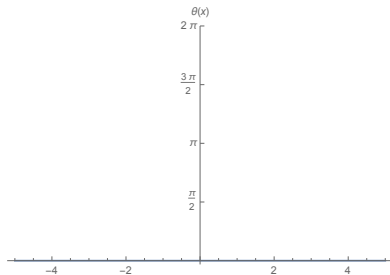
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after some blackboard work:

$$E = \int \frac{1}{2} (\theta'(x) - \sqrt{2V})^2 dx + Q \int_0^{2\pi} \sqrt{2V} d\theta$$



Chiral magnets and magnetic skyrmions

Take a magnet with a crystal structure that breaks parity. This allows an antisymmetric 'Dzyaloshinskii-Moriya' exchange term in the microscopic hamiltonian:

$$E[\{\mathbf{S}_i\}] = -J \sum_{\langle ij \rangle} \mathbf{S}_i \cdot \mathbf{S}_j + \sum_{\langle ij \rangle} \mathbf{D}_{ij} \cdot (\mathbf{S}_i \times \mathbf{S}_j) - \sum_i \mathbf{B} \cdot \mathbf{S}_i$$

Going to the continuum level, this gives the following energy functional:

$$E[n] = \int_{\mathbb{R}^2} \left(\frac{1}{2} (\nabla \mathbf{n})^2 + k \mathbf{n} \cdot (\nabla \times \mathbf{n}) + h_z (1 - n_3) + h_a (1 - n_3^2) \right)$$

The bogomol'nyi argument for magnetic skyrmions

Two steps:

1. We first reinterpret the model as a gauged sigma model with covariant derivative $D_i \mathbf{n} = \partial_i \mathbf{n} + \mathbf{A}_i \times \mathbf{n}$, $\mathbf{A}_i = -k \mathbf{e}_i$, and associated field strength $\mathbf{F}_{12} = k^2 \mathbf{e}_3$:

$$E[n] = \int_{\mathbb{R}^2} \left(\frac{1}{2} (D\mathbf{n})^2 + V(\mathbf{n}) \right)$$

2. When $h_a = -\frac{k^2}{2}$, $h_z = k^2$, we can rewrite the above as $E[n] = \int_{\mathbb{R}^2} \left(\frac{1}{2} (D_1 \mathbf{n} + \mathbf{n} \times D_2 \mathbf{n})^2 + \mathbf{n} \cdot (\partial_1 \mathbf{n} \times \partial_2 \mathbf{n}) + b.t. \right)$

Solving the bogomol'nyi equation

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For the case considered, where $\mathbf{A}_j = -k\mathbf{e}_j$, this becomes:

$$\partial_{\bar{z}} w = \frac{i}{2} k w^2 \implies \partial_{\bar{z}} \frac{1}{w} = -\frac{i}{2} k$$

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Which is solved by

$$w(z, \bar{z}) = -\frac{i}{2} k \bar{z} + f(z)$$

with $f(z)$ a smooth function of z . For finite degree, $f(z)$ must be a ratio of two polynomials and the degree of the solution can be found from the degrees of these polynomials

Conclusions and further work

- ▶ Reference: BBS, Calum Ross and Bernd J Schroers, 'Magnetic Skyrmions at Critical Coupling'
- ▶ This method from mathematical physics can be used to find explicit analytical solutions of the Euler-Lagrange equations for 'real world' topological solitons.
- ▶ What's more, some of the more outlandish solutions have now been seen numerically away from this 'critical point' (upcoming paper with Kiselev, Kuchkin and collaborators)
- ▶ On the theory side, these solutions could be used to investigate the dynamics of skyrmions, or their quantization. No progress yet however.