

YFS Resummation in SHERPA

ALAN PRICE

SUPERVISOR: PROF. FRANK KRAUSS

IPPP, DURHAM UNIVERSITY

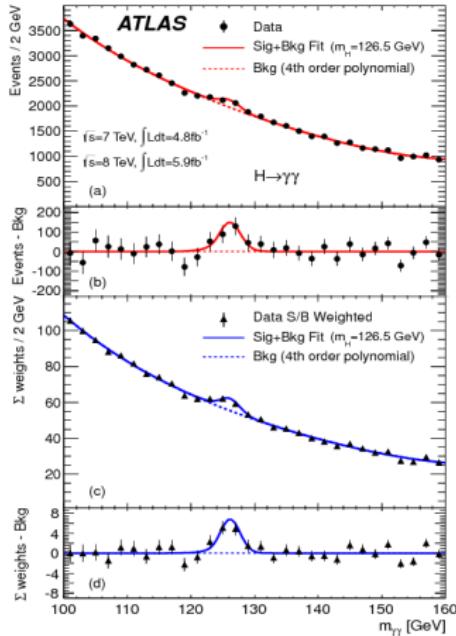
APRIL 3, 2020



INTRODUCTION

1. Motivation
2. YFS Resummation
3. Validation and Results
4. Conclusion

MOTIVATION



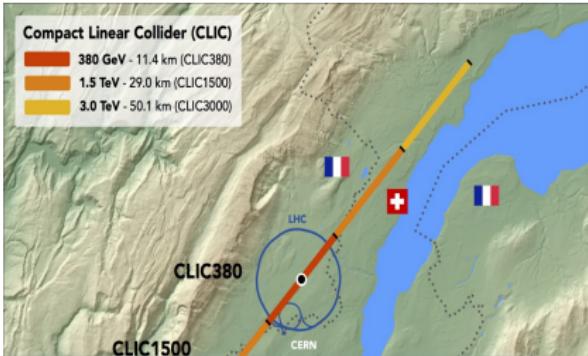
- For precision physics a lepton-lepton collider is desirable
- The discovery of Higgs Boson tells us at what energy to run
- High precision understanding of the Higgs Boson requires below 1% uncertainties in the couplings

MOTIVATION

We must not forget that an e^+e^- will precisely measure EW observables

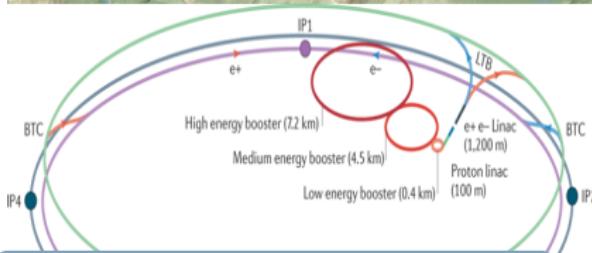
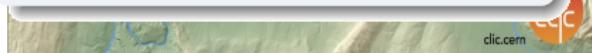
Observable	Where from	Present (LEP)	FCC stat.	FCC syst	Now FCC	Challenge
M_Z [MeV]	Z linesh.	$91187.5 \pm 2.1\{0.3\}$	0.005	0.1	3	QED Corrections
Γ_Z [MeV]	Z linesh.	$2495.2 \pm 2.1\{0.2\}$	0.008	0.1	2	QED Corrections
$R_I^Z = \Gamma_h/\Gamma_I$	$\sigma(M_Z)$	$20.767 \pm 0.025\{0.012\}$	$6 \cdot 10^{-5}$	$1 \cdot 10^{-3}$	12	QED Corrections
N_ν	$\sigma(M_Z)$	$2.984 \pm 0.008\{0.006\}$	$5 \cdot 10^{-6}$	$1 \cdot 10^{-3}$	6	Bhabha scattering (QED)
M_W [MeV]	ADLO	$80376 \pm 33\{6\}$	0.5	0.3	12	QED Corrections
$A_{FB,\mu}^{M_Z \pm 3.5 \text{ GeV}}$	$\frac{d\sigma}{d\cos\theta}$	$\pm 0.020\{0.001\}$	$1.0 \cdot 10^{-5}$	$0.3 \cdot 10^{-5}$	100	

Table: Adapted from Arxiv:1903.09895



CLIC

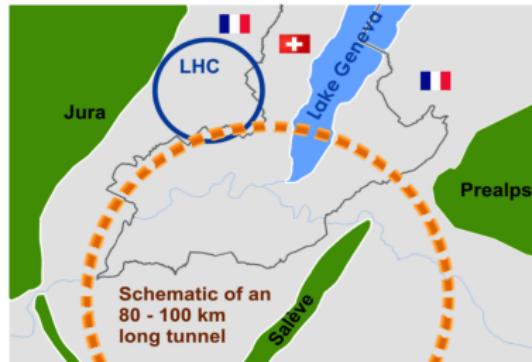
- $\sqrt{s} = 380 \text{ GeV} \rightarrow 3 \text{ TeV}$



CEPC

- $\sqrt{s} = 90 \rightarrow 250 \text{ GeV}$

SppC ring (100 km)



FCC

- $\sqrt{s} = 90 \rightarrow 350 \text{ GeV}$

Mandariz

Copyright CERN 2014

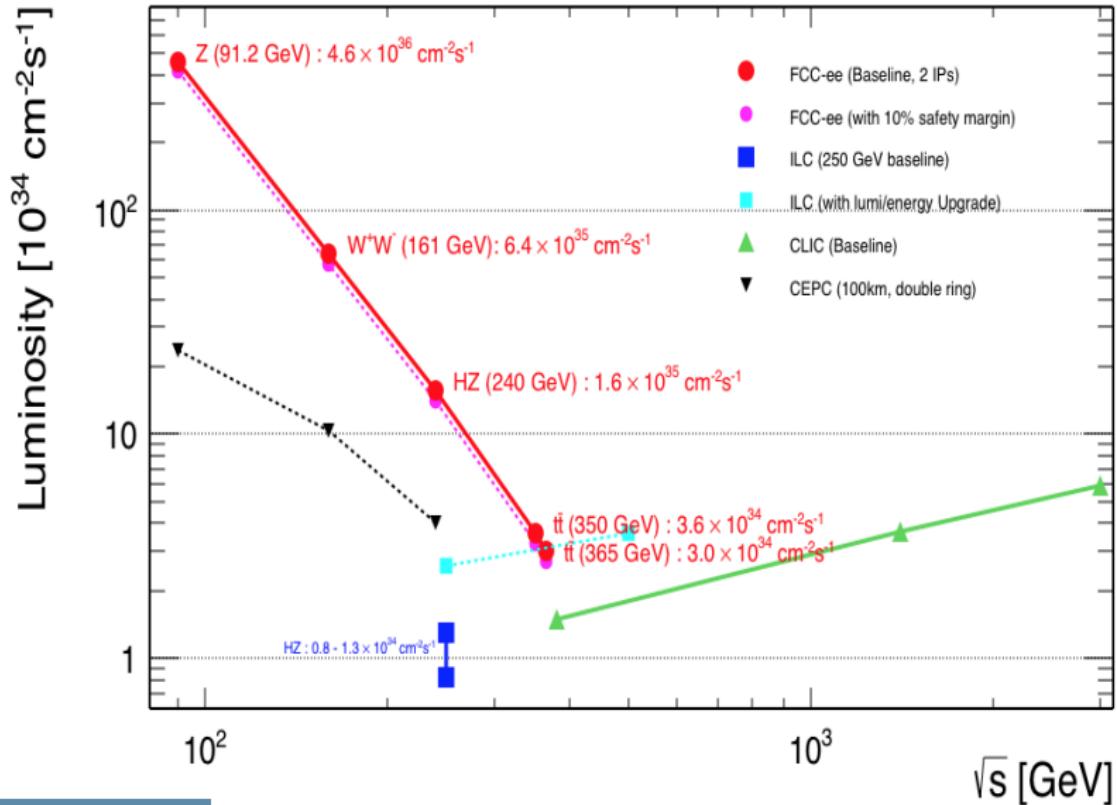


ILC

- $\sqrt{s} = 250 \text{ GeV} \rightarrow 1 \text{ TeV}$

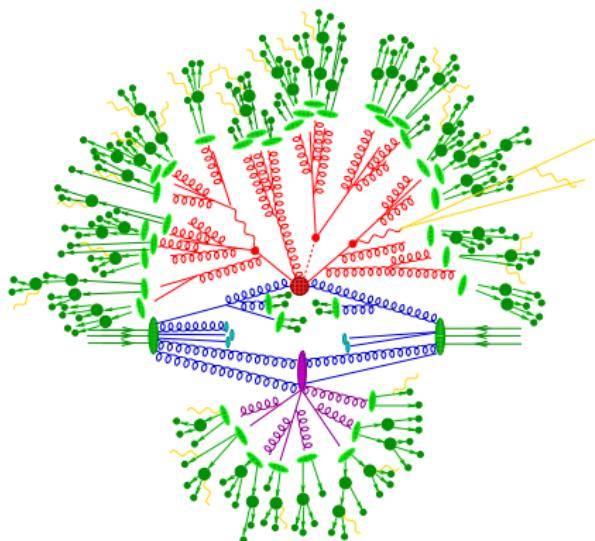


LINEAR VS CIRCULAR



SHERPA EVENT GENERATOR

- Hard Interaction
 - ▶ LO, NLO QCD/EW
- Radiative Corrections
 - ▶ Catani-Seymour based PS, Dire,
 - ▶ YFS QED Resummation
- Multiple Interactions
 - ▶ Sjöstrand-Zijl mode
- Hadronization
 - ▶ Cluster hadronization model
- Hadron Decays
 - ▶ Phase space or EFTs, YFS QED corrections



QED CORRECTIONS AT THE Z POLE

QED corrections (ISR) can be parametrised as,

$$\gamma_{n,m} = \left(\frac{\alpha}{\pi}\right)^n \left(2 \ln \left(\frac{m_Z^2}{m_e^2}\right)\right)^m \quad 0 \leq m \leq n$$

- For future lepton collider $\gamma \leq 10^{-5} \rightarrow$ need to include corrections up to $\mathcal{O}(\alpha^3, L_e^3)$, $L_e = \ln \left(\frac{s}{m_e}\right)$
- Soft and Collinear photons need to be resummed

CROSS SECTION CALCULATION

- The Problem: Emission of soft and/or Collinear photons which can cause divergences
- Solution: Calculate the Cross Section for $e^+ e^- \rightarrow X + n_\gamma$

$$\sigma = \frac{1}{\text{flux}} \int \sum_{i=0}^{n_R} d\Phi_q d\Phi_{k_i} (2\pi)^4 \delta(P - \sum q_f - \sum k_i) \left| \sum_{n_V=0}^{\infty} \mathcal{M}_{n_R}^{n_V + \frac{1}{2}n_R} \right|^2$$

CROSS SECTION CALCULATION

- The Problem: Emission of soft and/or Collinear photons which can cause divergences
- Solution: Calculate the Cross Section for $e^+ e^- \rightarrow X + n_\gamma$

$$\sigma = \frac{1}{\text{flux}} \int \sum_{i=0}^{n_R} d\Phi_q d\Phi_{k_i} (2\pi)^4 \delta(P - \sum q_f - \sum k_i) \left| \sum_{n_V=0}^{\infty} \mathcal{M}_{n_R}^{n_V + \frac{1}{2}n_R} \right|^2$$

YFS VIRTUAL CORRECTIONS

- Lets first consider virtual photon corrections in the soft limit

$$\mathcal{M}_0^0 = \underbrace{M_0^0}_{\text{Born ME}}$$

$$\mathcal{M}_0^1 = \alpha B M_0^0 + M_0^1$$

$$\mathcal{M}_0^1 = \frac{(\alpha B)^2}{2} M_0^0 + \alpha B M_0^1 + M_0^2$$

...

$$\mathcal{M}_0^{n_V} = \sum_{r=0}^{n_V} M_0^{n_V} \frac{(\alpha B)^r}{r!}$$

where B is the virtual infrared factor.

$$B = 2\alpha \Re \int \frac{d^4 k}{k^2} \frac{i}{(2\pi)^2} \left(\frac{2p_1 - k}{2kp_1 - k^2} - \frac{2p_2 - k}{2kp_2 - k^2} \right)^2$$

YFS VIRTUAL CORRECTIONS

Summing to infinity yields,

$$\sum_{n_V=0}^{\infty} \mathcal{M}_0^{n_V} = e^{\alpha B} \sum_{n_V=0}^{\infty} M_0^{n_V}$$

Due to the Abelian nature of QED this can be generalised to n_R real photons,

$$\sum_{n_V=0}^{\infty} \left| \mathcal{M}_{n_R}^{n_V + \frac{1}{2} n_R} \right|^2 = e^{2\alpha B} \sum_{n_V=0}^{\infty} \left| M_{n_R}^{n_V + \frac{1}{2} n_R} \right|^2$$

YFS REAL EMISSIONS

For a single photon emission we have,

$$\frac{1}{2(2\pi)^3} \sum_{n_V=0}^{\infty} \left| M_1^{n_V+\frac{1}{2}} \right|^2 = \tilde{S}(k) \left| M_0^{n_V+\frac{1}{2}} \right|^2 + \sum_{n_V=0}^{\infty} \tilde{\beta}_1^{n_V+1}(k)$$

- Factorisation of real emissions occurs at the amplitude squared level
- Eikonal term $\tilde{S}(k) = -\frac{\alpha}{4\pi^2} \left(\frac{p_1}{p_1 k} - \frac{p_2}{p_2 k} \right)^2$
- $\tilde{\beta}_{n_R}^{n_V+n_R}$ complete IR finite squared matrix element for born process plus n_V virtual and n_R photons
- Note abbreviation used in this talk

$$\tilde{\beta}_{n_R} = \sum_{n_V=0}^{\infty} \tilde{\beta}_{n_R}^{n_V+\frac{1}{2}n_R}$$

YFS REAL EMISSIONS

For n_R photons summed over all virtual contributions,

$$\begin{aligned} & \left(\frac{1}{2(2\pi)^3} \right)^{n_R} \left| \sum_{n_V=0}^{\infty} M_{n_R}^{n_V + \frac{1}{2}n_R} \right|^2 \\ = & \tilde{\beta}_0 \prod_{i=1}^{n_R} [\tilde{S}(k_i)] + \sum_{i=1}^{n_R} \left[\frac{\tilde{\beta}_1(k_i)}{\tilde{S}(k_i)} \right] \prod_{j=1}^{n_R} [\tilde{S}(k_j)] \\ & + \sum_{\substack{i,j=1 \\ i \neq j}}^{n_R} \left[\frac{\tilde{\beta}_2(k_i, k_j)}{\tilde{S}(k_i)\tilde{S}(k_j)} \right] \prod_{l=1}^{n_R} [\tilde{S}(k_l)] + \dots \\ & + \sum_{i=1}^{n_R} \left[\tilde{\beta}_{n_R-1}(k_1, \dots, k_{i-1}, k_{i+1}, \dots, k_{n_R}) \tilde{S}(k_i) \right] + \tilde{\beta}_{n_R}(k_1, \dots, k_{n_R}) \end{aligned}$$

MASTER EQUATION

This gives us our cross section

$$\sigma = \sum_{n=0}^{\infty} \frac{1}{n!} \int d\Phi_f e^{2\alpha B + 2\alpha \tilde{B}} \prod_{j=1}^n \tilde{S}(k_j) \theta(\Omega; k_j) \left[\tilde{\beta}_0(p_1, p_2; q_1, \dots, q_{n'}) \right. \\ \left. + \sum_{j=1}^n \frac{\tilde{\beta}_1(p_1, p_2; q_1, \dots, q_{n'}; k_j)}{S(k_j)} \right. \\ \left. + \sum_{\substack{j, l=1 \\ j \neq l}}^n \frac{\tilde{\beta}_2(p_1, p_2; q_1, \dots, q_{n'}; k_j, k_l)}{S(k_j)S(k_l)} + \dots \right]$$

The exponentiation of real emissions gives a factor

$$\tilde{B} = -\frac{1}{8\pi^2} \int \frac{d^3 k}{k^0} \Theta(\Omega, k) \left(\frac{p_1}{p_1 k} - \frac{p_2}{p_2 k} \right)^2$$

IR FINITE ME

$$\tilde{\beta}_0^0 = M_0^0 M_0^{0*}$$

$$\tilde{\beta}_0^1 = M_0^1 M_0^{0*} + M_0^0 M_1^{0*}$$

$$\tilde{\beta}_1^1 = \frac{1}{2(2\pi)^3} M_0^{\frac{1}{2}} M_0^{\frac{1}{2}*} - \tilde{S}(k) M_0^0 M_0^{0*} = \frac{1}{2(2\pi)^3} M_0^{\frac{1}{2}} M_0^{\frac{1}{2}*} - \tilde{S}(k) \tilde{\beta}_0^0$$

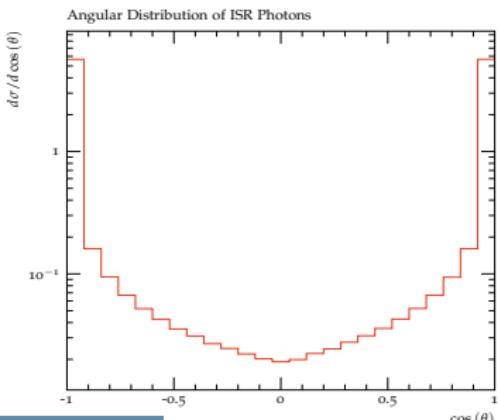
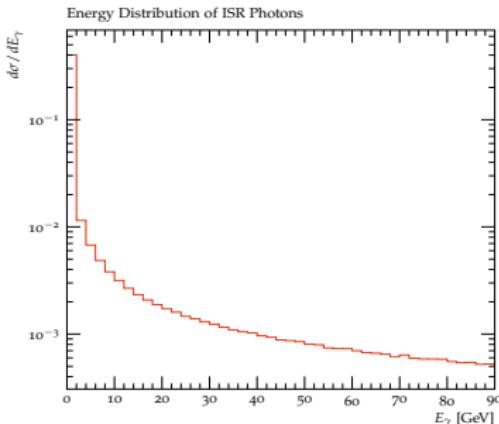
- These are the infrared subtracted squared matrix elements read, up to $\mathcal{O}(\alpha)$
- They have been implemented in SHERPA up to $\mathcal{O}(\alpha^2)$ for ISR

THE ALGORITHM

1. First the phasespace is rearrange by a change of variable $v = 1 - \frac{s'}{s}$
2. v is then generated by standard Monte-Carlo Methods e.g Importance Sampling
3. If $v < \epsilon$ no photons are generate. For $v \geq \epsilon$ the multiplicity is generated according to a Poissonian distribution with average $\mu = \frac{2\alpha}{\pi} \log\left(\frac{s}{m_e^2}\right) \log\left(\frac{v}{\epsilon}\right)$
4. The n photon momenta are distributed according to the Eikonal $\tilde{S}(k)$ and then rescaled to ensure 4-momentum conservation
5. The perturbative parts can be calculated separately using ME generator
6. Give everything to the master equation to calculate the cross section

THE POWER OF YFS

- $\tilde{\beta}$ are infrared finite and are calculated perturbatively *order by order*
 - ▶ Non trivial for other methods e.g Structure function
- Photons are explicitly created
 - ▶ Allows us to calculate properties of soft photons e.g Energy, Angular distribution
 - ▶ Allows for decays e.g $\gamma \rightarrow e^+ e^-$



VALIDATION

- For LEP there were dedicated MC codes that implemented YFS for specific processes
 - KKMC [Jadach et.al, ARXIV:9912214](#)
 - KoralWW [Jadach et.al, ARXIV:0012094](#)
 - YFSWW [Jadach et.al, ARXIV:0012094](#)

VALIDATION

- For LEP there were dedicated MC codes that implemented YFS for specific processes
 - KKMC [Jadach et.al, ARXIV:9912214](#)
 - KoralWW [Jadach et.al, ARXIV:0012094](#)
 - YFSWW [Jadach et.al, ARXIV:0012094](#)

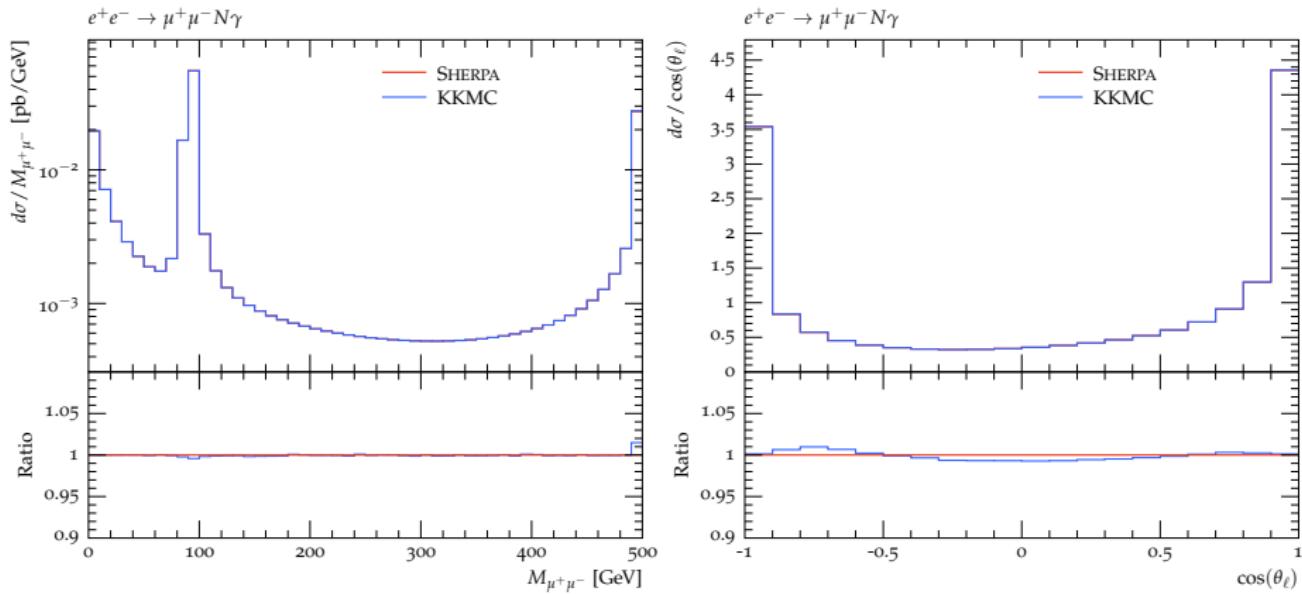
VALIDATION

YFS $\mathcal{O}(\alpha^0)$	KKMC 500 GeV	SHERPA 500 GeV
$e^+e^- \rightarrow \mu^+\mu^-$	1.748(3) pb	1.746(3) pb
$e^+e^- \rightarrow \nu_\mu\bar{\nu}_\mu$	1.6235(8) pb	1.626(3) pb
$e^+e^- \rightarrow \nu_\tau\bar{\nu}_\tau$	1.6235(8) pb	1.626(3) pb
$e^+e^- \rightarrow jj$	19.94(2) pb	19.98(8) pb

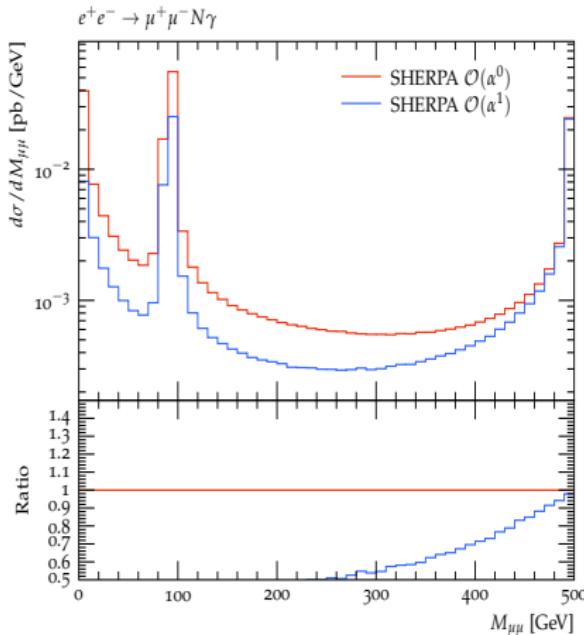
YFS $\mathcal{O}(\alpha^1)$	KKMC 500 GeV	SHERPA 500 GeV
$e^+e^- \rightarrow \mu^+\mu^-$	1.024(3) pb	1.020(3) pb
$e^+e^- \rightarrow \nu_\mu\bar{\nu}_\mu$	0.8072(7) pb	0.808(2) pb
$e^+e^- \rightarrow \nu_\tau\bar{\nu}_\tau$	0.8072(7) pb	0.808(2) pb
$e^+e^- \rightarrow jj$	10.54(1) pb	10.21(4) pb

YFS $\mathcal{O}(\alpha^2)$	KKMC 500 GeV	SHERPA 500 GeV
$e^+e^- \rightarrow \mu^+\mu^-$	1.080(3) pb	1.082(6) pb
$e^+e^- \rightarrow \nu_\mu\bar{\nu}_\mu$	0.8604(7) pb	0.860(3) pb
$e^+e^- \rightarrow \nu_\tau\bar{\nu}_\tau$	0.8604(7) pb	0.860(3) pb
$e^+e^- \rightarrow jj$	11.18(2) pb	11.10(8) pb

DIFFERENTIAL DISTRIBUTIONS

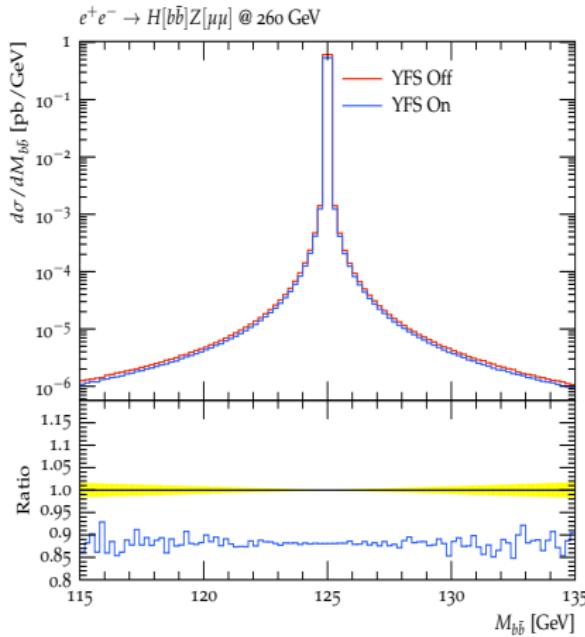


DIFFERENTIAL DISTRIBUTION



- The large deviation between α^0 and higher order is expected and is in agreement with semi-analytical approximation
- $\frac{\mathcal{O}(\alpha^1)}{\mathcal{O}(\alpha^0)} \approx 0.5$ at the z-pole
- $\frac{\mathcal{O}(\alpha^2)}{\mathcal{O}(\alpha^1)} \approx 0.98$ at the z-pole

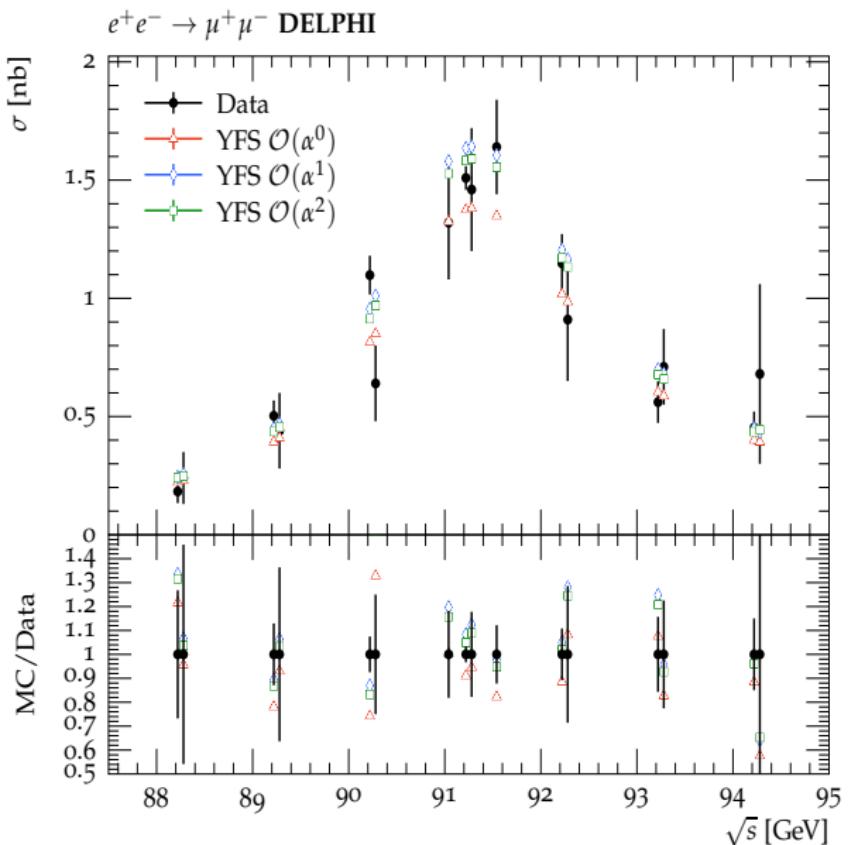
HIGGS-STRÄHLUNG



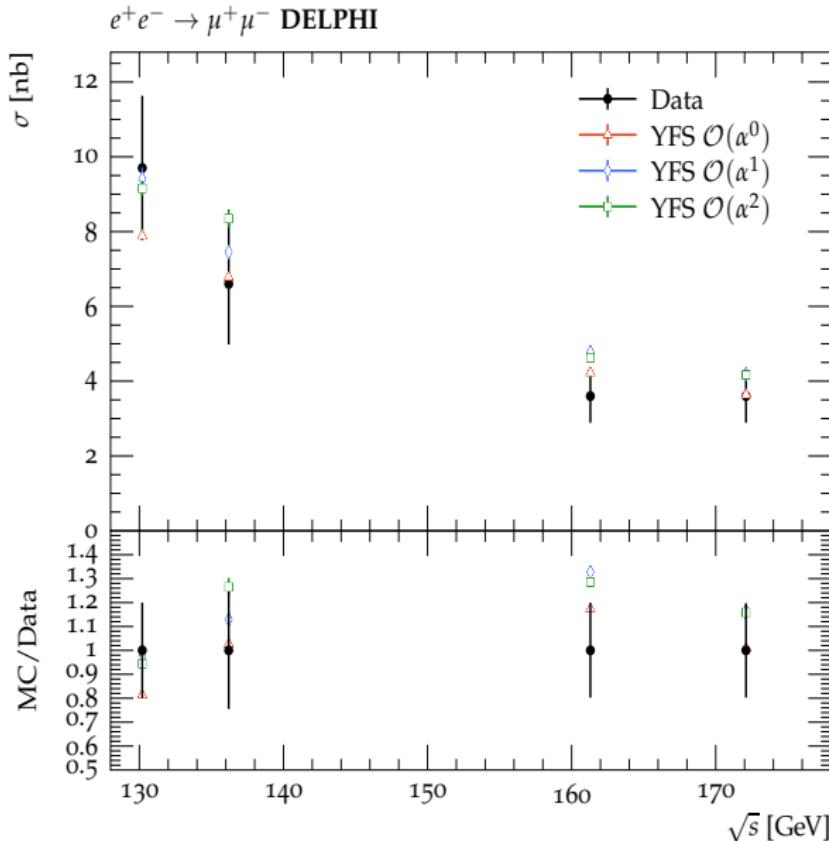
- At the Higgs pole ISR contributes a $\approx 10\%$ effect

	ZH-Pole 240 GeV
$\mathcal{O}(\alpha^0)$	201.46 fb
$\mathcal{O}(\alpha)$	234.59 fb
$\mathcal{O}(\alpha^2)$	233.23 fb

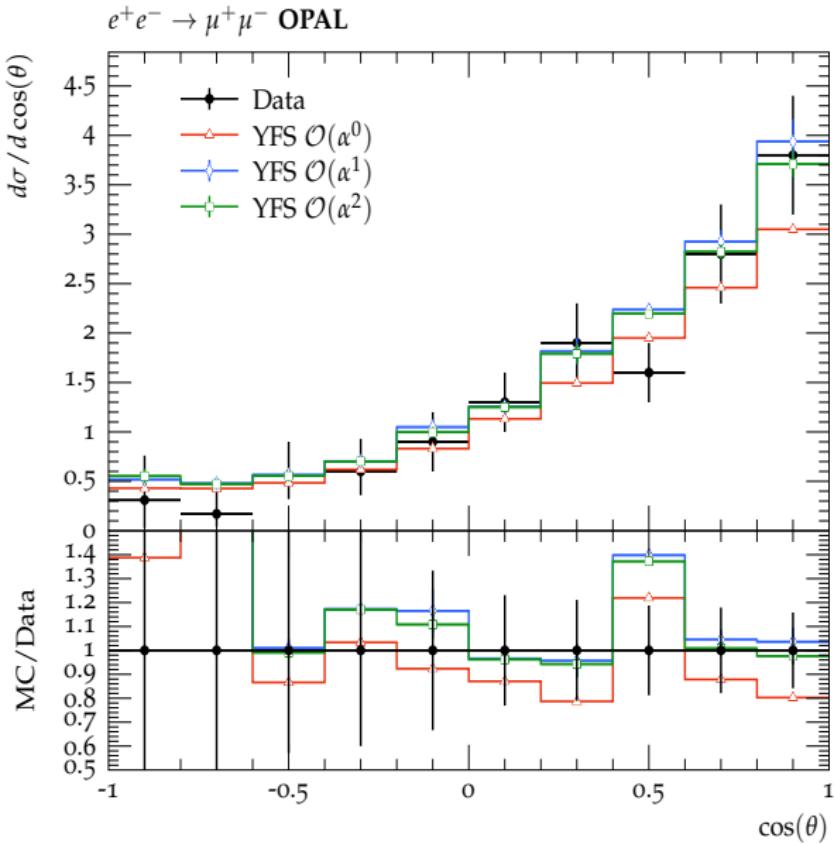
COMPARISON TO DATA



COMPARISON TO DATA



COMPARISON TO DATA



CONCLUSION

- Interface with other SHERPA modules e.g ME Generator, PS, Hadronization, UFO...
- Fully automated for arbitrary processes
- Higher order corrections have been hard coded for ISR
- YFS has been fully implemented for FSR up to NNLO QCD and NLO_{EW} [Krauss et.al, ARXIV:1809.10650](#)
- Matching of YFS resummation to NLO_{EW} is under-way
 - EW Loops are now provided in full by OPENLOOPSS [Buccioni et.al, ARXIV:1907.13071](#)
 - Real emission ME provided by COMIX/AMEGIC
 - Main problem is lack of massive leptons in the loops

The SHERPA 2.2 event generator framework

User Inputs	Matrix Elements	Parton Showers	Soft Physics	Interfaces/Outputs
Initial Beams <ul style="list-style-type: none">• collider setup• PDFs (built-in, LHAPDF)• beam spectra	Matrix Element Generators <ul style="list-style-type: none">• AMEGIC• COMIX• CS subtraction	CS-Shower (default) <ul style="list-style-type: none">• dipole shower• fully massive• QED splittings DIRE <ul style="list-style-type: none">• hybrid dipole-parton shower algorithm• fully massive	Hadronisation <ul style="list-style-type: none">• AHADIC: a cluster fragmentation model• interface to Pythia string fragmentation	Output Formats <ul style="list-style-type: none">• HepMC• LHEF• Root Ntuple
Parameters/Models <ul style="list-style-type: none">• FeynRules/UFO• couplings• masses• variations• shower settings• non-perturbative parameters	1-loop Amplitudes <ul style="list-style-type: none">• OpenLoops• Recola• GoSam• BLHA 	 	Hadron Decays <ul style="list-style-type: none">• decay tables for hadronic resonances• dedicated form-factor models, e.g. τ, B, Λ• spin correlations• YFS QED corrections• partonic channels	Interfaces <ul style="list-style-type: none">• RIVET analyses• C++/Python ME access• MCgrid• integration into ATLAS/CMS
Physics Process <ul style="list-style-type: none">• parton level• perturbative order (QCD/EW)• selectors• matching/merging• partonic decays	Matching and Merging Automated MC@NLO style matching Multijet-merging algorithms <ul style="list-style-type: none">• based on truncated showers• tree-level and one-loop matrix elements: MEPS@LO and MEPS@NLO• approximate electroweak corrections NNLO QCD with parton showers <ul style="list-style-type: none">• selected processes only		 	Code/Docu <ul style="list-style-type: none">• HepForge• GitLab• online documentation <p>sherpa.hepforge.org</p> <p>gitlab.com/sherpa-team/sherpa</p>