
Quantum computers: a breakthrough in information processing and application in High Energy Physics

Michele Grossi

IBM Tech Quantum Ambassador
Qiskit Advocate

University of Pavia

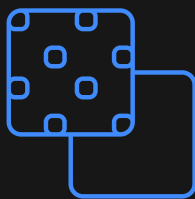


With quantum computers, we will tackle problems in new ways

Model physical
processes
of nature



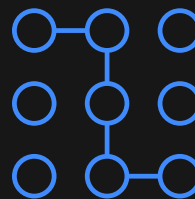
Find better
patterns within
AI/ML processes



Obtain better
optimization
solutions



Perform significantly
more scenario
simulations



Start thinking about quantum computing now

It's moving fast



Talent is scarce



Patents are being
filed



Strategizing takes
time



IBM is defining the future of computing, again

Most advanced hardware supporting the full stack

14 systems
53 qubit system



ibm.co/q-experience

Trusted advisor with the broadest adoption

180k+ users
100B+ executions



qiskit.org

Built the network of experts defining the future together

80+ partners
200+ papers



ibm.co/q-network

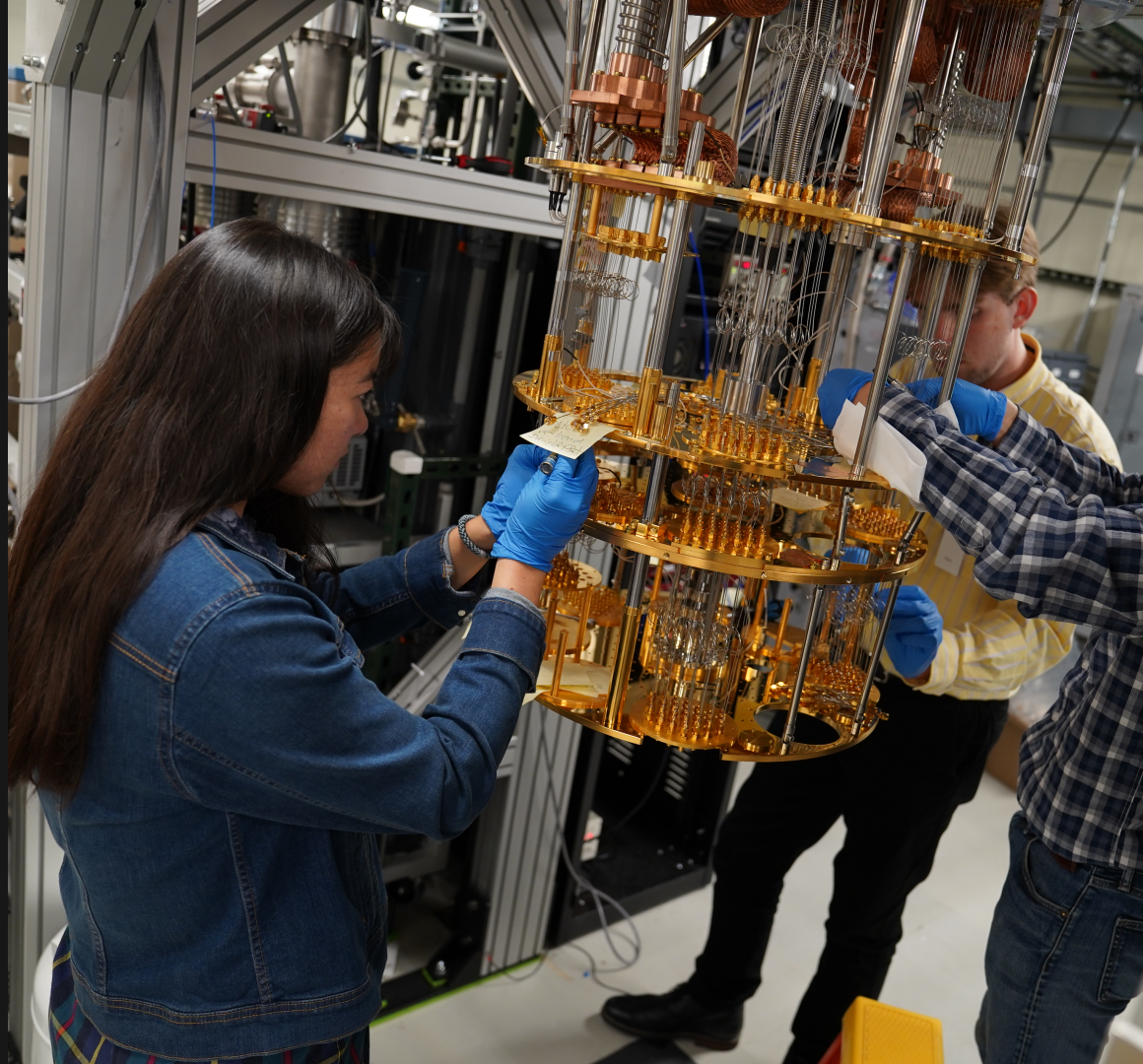
The future belongs to all of us

Giving students hands-
on experience through
internships

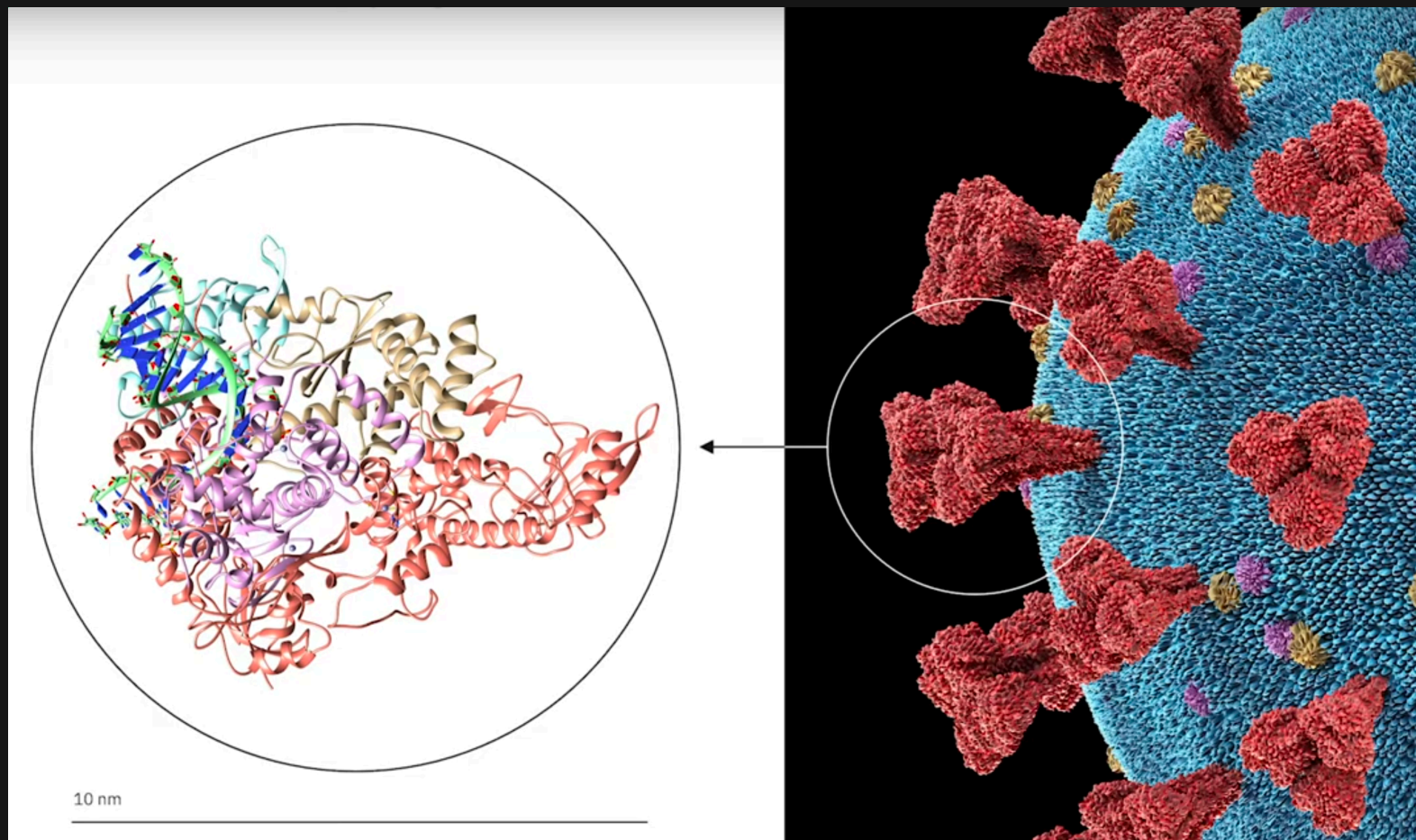
Offering open-access
educational materials

community.qiskit.org/education

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The motivation for quantum computing



The future of computing

Mathematics + Information
Today's computers and HPC

Intelligent
Applications

bits

neurons

qubits

Hybrid Cloud

Secure heterogeneous computational fabric

Biology + Information
AI Systems

Intelligent Automation
Automated programming and AI

Physics + Information
Quantum Systems

“I’m not happy with all the analyses that go with just the classical theory, because nature isn’t classical, dammit, and if you want to make a simulation of nature, you’d better make it quantum mechanical ...”

Richard P. Feynman
Department of Physics,
California Institute of Technology

International Journal of Theoretical Physics,
Vol 21, Nos. 6/7, 1982

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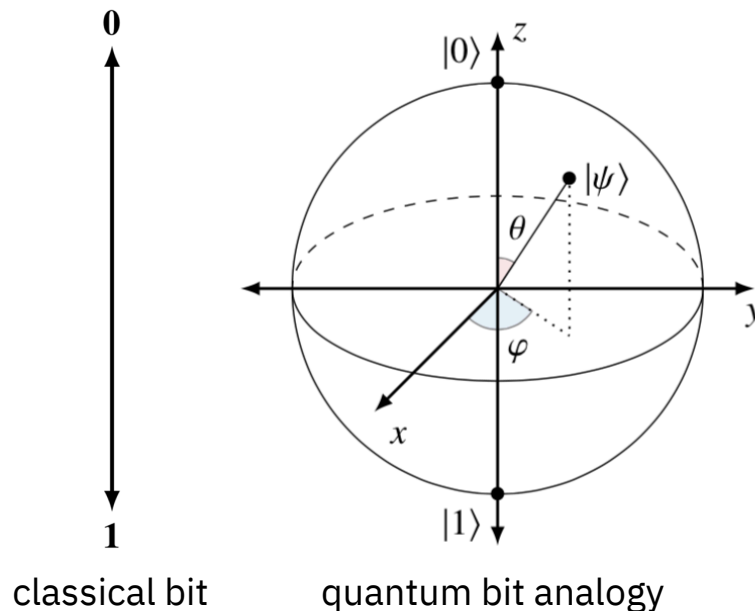
The basics of quantum computing

Bits and qubits

A **classical bit** can be **0** or **1**.

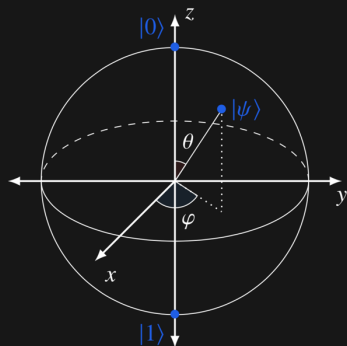
A **quantum bit**, or **qubit**, can take on those values but can represent a combination of **0** and **1** while we are computing.

When we measure a qubit, it becomes **0** or **1** based on probability.



Quantum computing uses essential ideas from quantum mechanics

Measurement



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Measurement is forcing the qubit's state IBM Quantum

$$a |0\rangle + b |1\rangle$$

to $|0\rangle$ or $|1\rangle$ by observing it, where

$|a|^2$ is the probability we will get $|0\rangle$ when we measure

$|b|^2$ is the probability we will get $|1\rangle$ when we measure

Examples

$$\frac{\sqrt{2}}{2} |0\rangle + \frac{\sqrt{2}}{2} |1\rangle$$

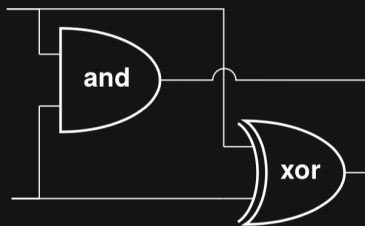
has an equal probability of becoming $|0\rangle$ or $|1\rangle$.

$$\frac{\sqrt{3}}{2} |0\rangle - \frac{1}{2} i |1\rangle$$

has a 75% chance of becoming $|0\rangle$.

Quantum computing uses essential ideas from quantum mechanics

Gates / Operations



Classical logical circuits use operations like **and**, **or**, **not**, **nand**, and **xor**. We also call these gates.

Quantum circuits use reversible gates that change the quantum states of one, two, or more qubits.

$$q_0: |0\rangle \xrightarrow{\text{H}} \text{Measurement} \rightarrow |m_0\rangle = |0\rangle \text{ or } |1\rangle$$

$$q_1: |0\rangle \xrightarrow{\text{H}} \text{Measurement} \xrightarrow{\text{H}} \text{Measurement} \rightarrow |m_1\rangle = |0\rangle \text{ or } |1\rangle$$

$$q_2: |0\rangle \xrightarrow{\text{H}} \text{H} \xrightarrow{\text{Measurement}} |m_2\rangle = |0\rangle$$

$$\begin{array}{l} q_0: |\psi\rangle_0 \xrightarrow{\text{H}} \text{H} \xrightarrow{\text{H}} \text{Measurement} \rightarrow |m_0\rangle = |\psi\rangle_1 \\ q_1: |\psi\rangle_1 \xrightarrow{\text{H}} \text{H} \xrightarrow{\text{H}} \text{Measurement} \rightarrow |m_1\rangle = |\psi\rangle_0 \end{array}$$

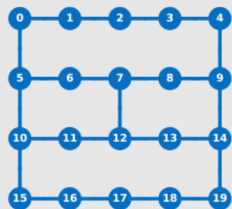
The diagram shows two parallel quantum circuits. The top circuit for q_0 starts with state $|\psi\rangle_0$, followed by three Hadamard (H) gates and a measurement gate, resulting in state $|m_0\rangle = |\psi\rangle_1$. The bottom circuit for q_1 starts with state $|\psi\rangle_1$, followed by three Hadamard (H) gates and a measurement gate, resulting in state $|m_1\rangle = |\psi\rangle_0$. Vertical lines connect the measurement gates of the two circuits, indicating a classical control or correlation between the results.

The hardware in an IBM Q quantum computing system

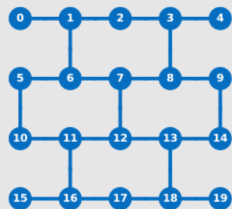


IBM Q quantum devices

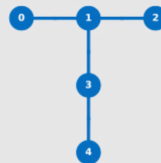
IBM's 10 Quantum Device Lineup



Johannesburg
Poughkeepsie



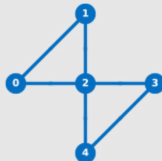
Almaden
Boeblingen
Singapore



Ourense
Valencia
Vigo

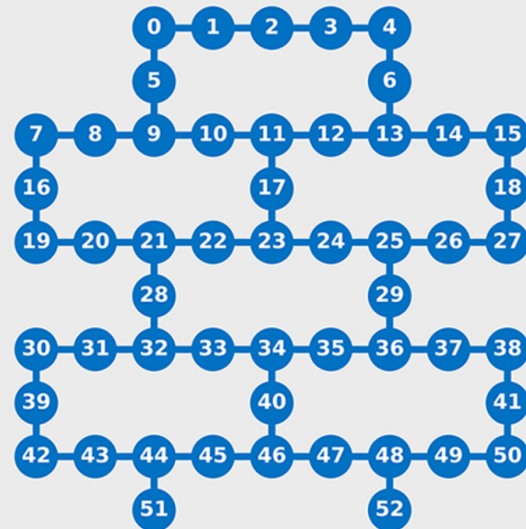


Melbourne



Yorktown

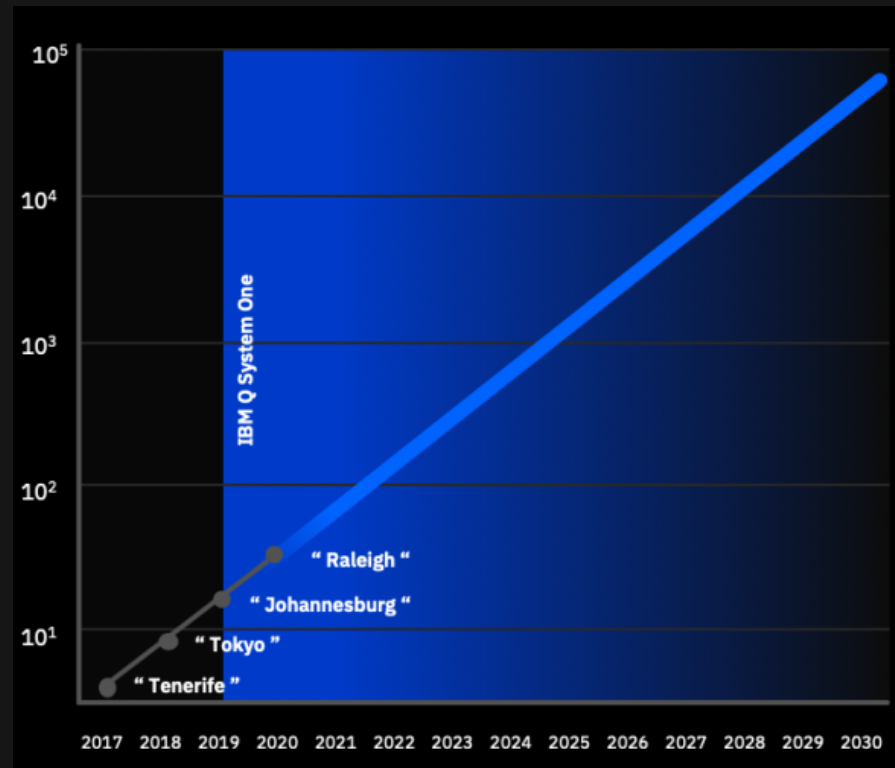
53 Qubit Rochester Device



Quantum volume

Many factors contribute to the performance of the overall system

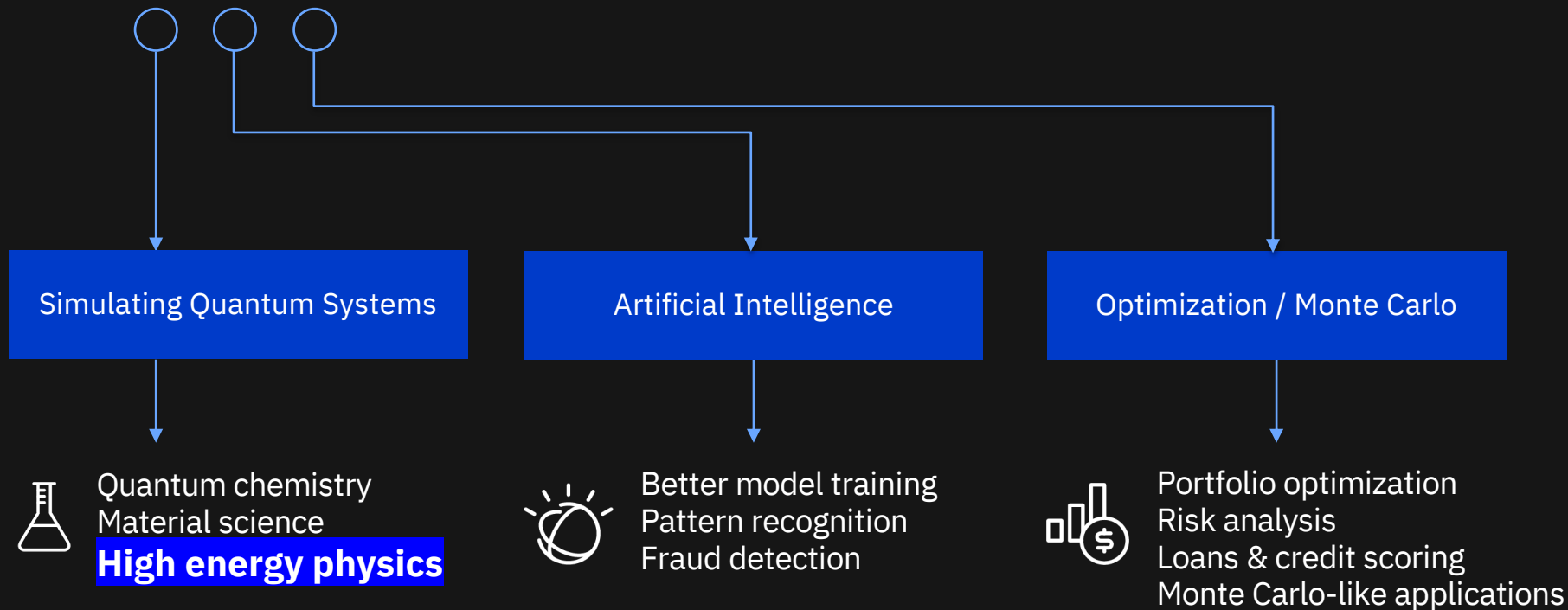
<https://www.ibm.com/blogs/research/2020/01/quantum-volume-32/>





IBM Q Experience is the quantum cloud services and software platform designed to take full advantage of IBM Q systems.

In collaboration with IBM Q Network partners we are driving advancements in quantum software and algorithms
IBM research towards the first use cases with quantum advantage...

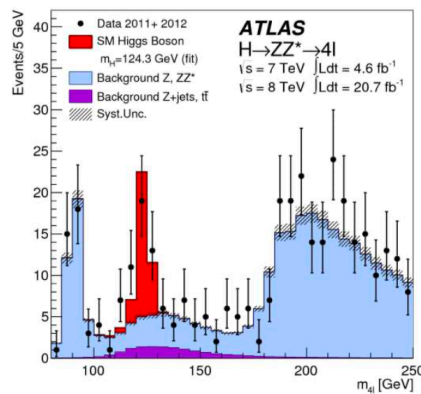
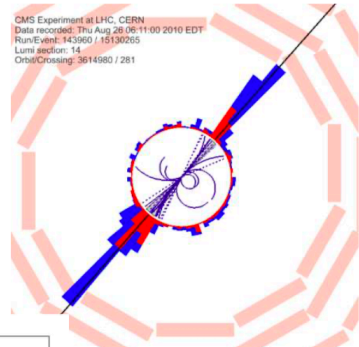


High Energy Physics

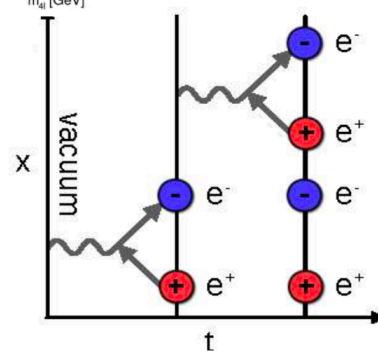
Classification (machine learning and SVM),
select/identify relevant LHC events, reconstruction
of tracks - jets tracking

Quantum-Computing where Time-Evolution: lattice
gauge theory (Schwinger's model and beyond)

Quantum Minimization VQE optimization in lattice
gauge theory



E.A. Martinez et al., nature,
534, 516 (2016)



Simulating lattice gauge theories on a quantum computer

Tim Byrnes* Yoshihisa Yamamoto

2005

Quantum Computation of Scattering
in Scalar Quantum Field Theories

2012

Stephen P. Jordan,^{†§} Keith S. M. Lee,^{†§} and John Preskill ^{§ *}

Atomic Quantum Simulation of $U(N)$ and $SU(N)$ Non-Abelian Lattice Gauge Theories

2013

D. Banerjee¹, M. Bögli¹, M. Dalmonte², E. Rico^{2,3}, P. Stebler¹, U.-J. Wiese¹, and P. Zoller^{2,3}

2014

Towards Quantum Simulating QCD

Uwe-Jens Wiese

Quantum Simulations of Lattice Gauge Theories
using Ultracold Atoms in Optical Lattices

Erez Zohar J. Ignacio Cirac Benni Reznik

2015

2016

Real-time dynamics of lattice gauge theories with a few-qubit quantum computer

Esteban A. Martinez,^{1,*} Christine Muschik,^{2,3,*} Philipp Schindler,¹ Daniel Nigg,¹ Alexander Erhard,¹ Markus Heyl,^{2,4} Philipp Hauke,^{2,3} Marcello Dalmonte,^{2,3} Thomas Monz,¹ Peter Zoller,^{2,3} and Rainer Blatt^{1,2}

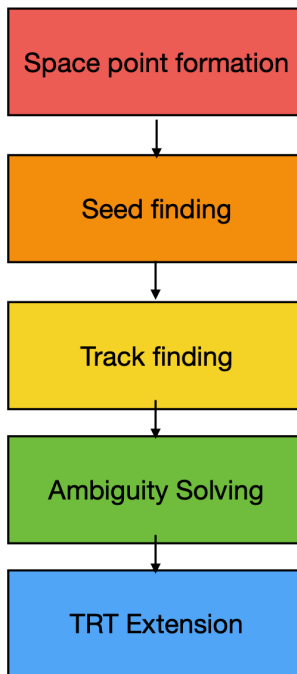
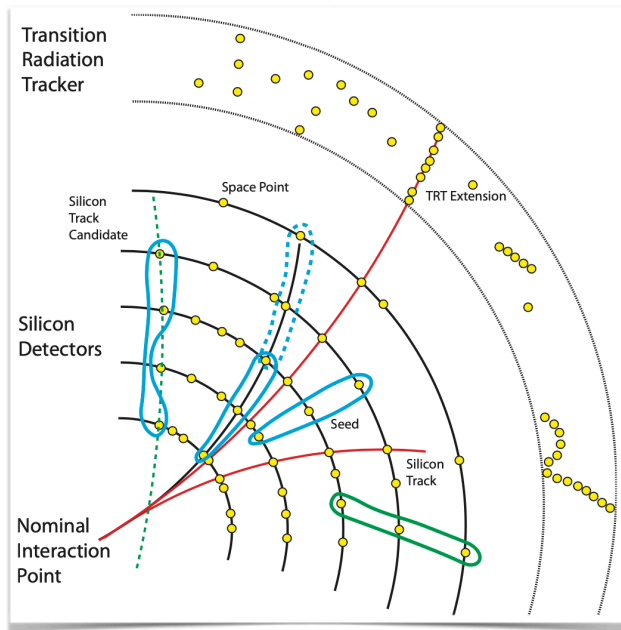
Quantum Sensors for the Generating Functional of Interacting Quantum Field Theories

2017

A. Bermudez,^{1,2,*} G. Aarts,¹ and M. Müller¹

Track Reconstruction – QUBO Algorithm

Multi-step iterative Kalman filter approach



Mapping of QUBO to Hamiltonian

1. Set $x_i = \frac{z_i+1}{2}$, for $z_i \in \{-1, 1\}$
2. Replace z_i by σ_Z^i and $z_i z_j$ by $\sigma_Z^i \otimes \sigma_Z^j$ where

$$\sigma_Z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

→ Combinatorial optimization problem has been translated to ground state problem of Hamiltonian H as known from quantum chemistry

$$\min_{|\psi\rangle} \langle \psi | H | \psi \rangle$$

Quadratic Unconstrained Binary Optimization QUBO

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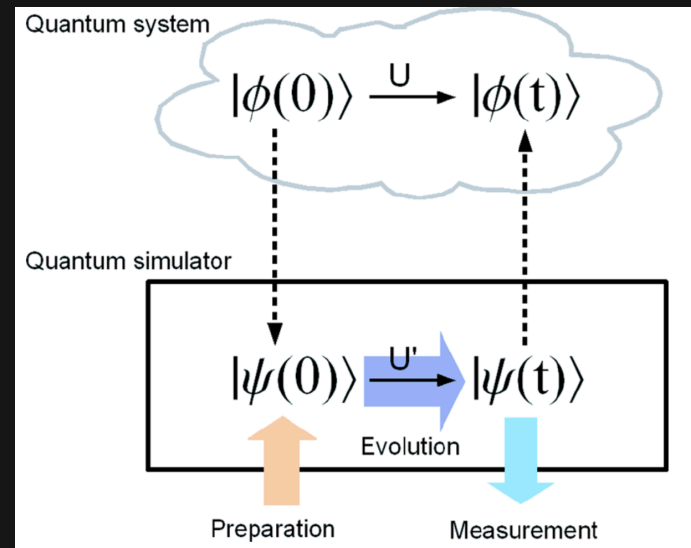
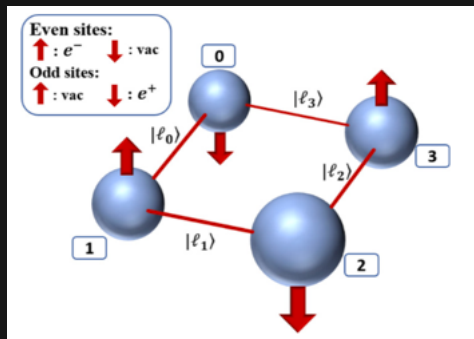
<https://sites.google.com/lbl.gov/hep-qpr>

Quantum Simulation

- Represent the system Hilbert space on the qubit space
- Implement a circuit to simulate the time evolution
- Field on a lattice
- The field amplitude is discretized at every lattice site

Example: **Schwinger** model describes 2D QED with a Dirac fermion

- toy QCD model:
- Charge screening
 - SSB
 - confinement



S. Lloyd, Science **273**, 1073 (1996)

ϕ^4 - quantum field theory

S. Jordan, et al., Science, **336**, 1130 (2012)

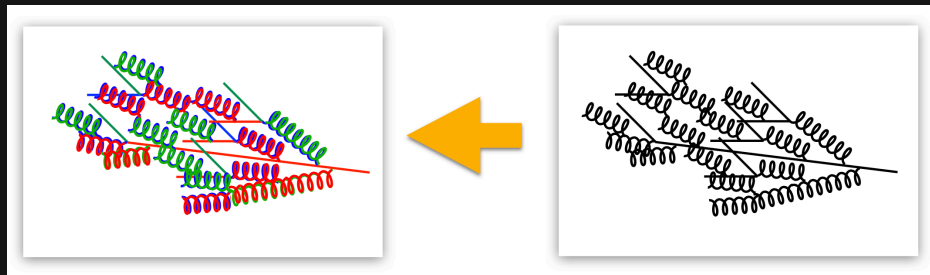
SU(2) non-Abelian gauge field theory in one dimension on digital quantum computers

Natalie Kico, Jesse R. Stryker and Martin J. Savage¹

¹ Institute for Nuclear Theory, University of Washington, Seattle, WA 98195-1550, USA

(Dated: August 20, 2019 - 0:44)

Quantum Computing for Final State Radiation



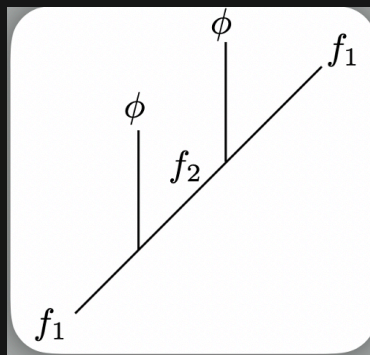
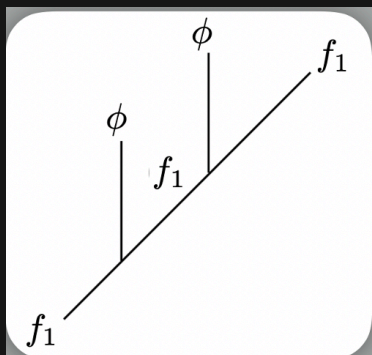
Final State Radiation (FSR) is a complex many-body quantum system. Classic MC simulation cannot capture all quantum effects.

Parton shower models are implemented using classical Markov Chain MC (MCMC) algorithms to efficiently generate high multiplicity radiation patterns.

Perhaps quantum tools can be used to incorporate quantum degrees of freedom

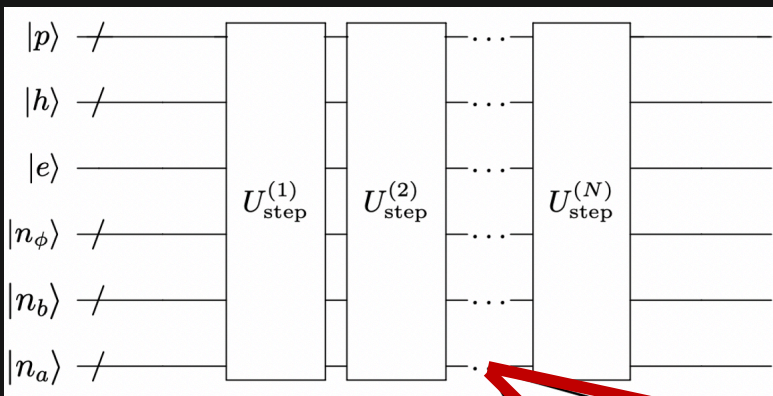
Ref: 1904.03196 D.Provasoli, C. Bauer, W. de Jong

Like the SM Higgs when $g_{1,2} \sim m/v$ and $g_1 = g_2 = 0$



$$\mathcal{L} = \bar{f}_1 i(\not{\partial} + m_1) f_1 + \bar{f}_2 (i\not{\partial} + m_2) f_2 + (\partial_\mu \phi)^2 \\ + g_1 \bar{f}_1 f_1 \phi + g_2 \bar{f}_2 f_2 \phi + g_{12} [\bar{f}_1 f_2 + \bar{f}_2 f_1] \phi$$

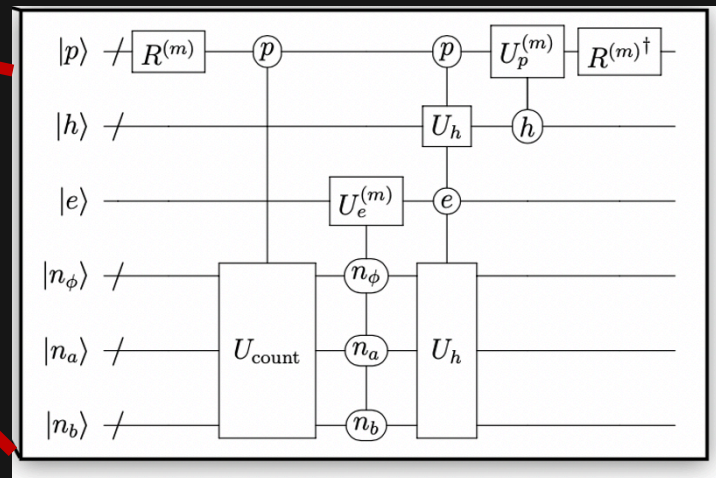
Quantum Computing for FSR: algorithm



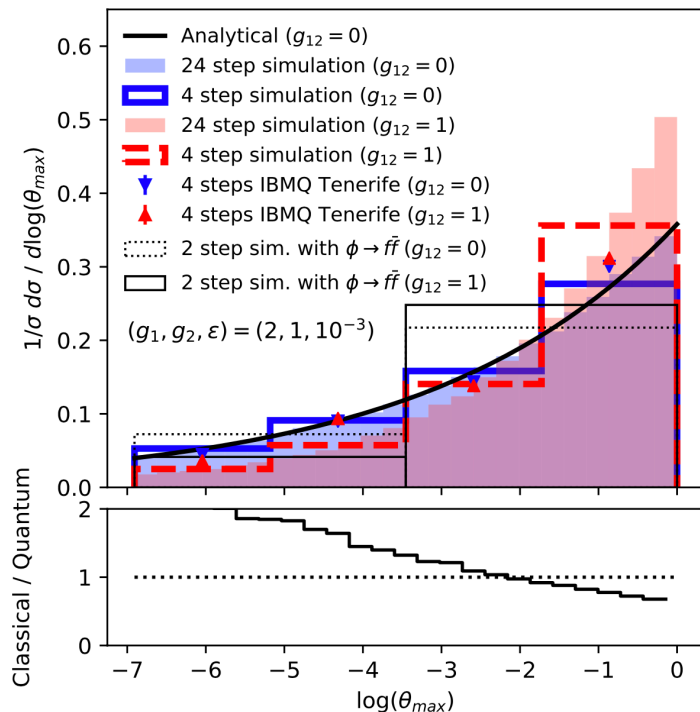
The *algorithm* is able to keep track of amplitudes and not probabilities, it samples from the full probability distribution in polynomial time

Measurement: normalized differential cross section for $\log \theta_{\text{max}}$ and the number of emissions. Interference effects are turned on ($g_{12} = 1$) and off ($g_{12} = 0$), where the classical simulations/calculations are expected to agree with the quantum simulations and measurements.

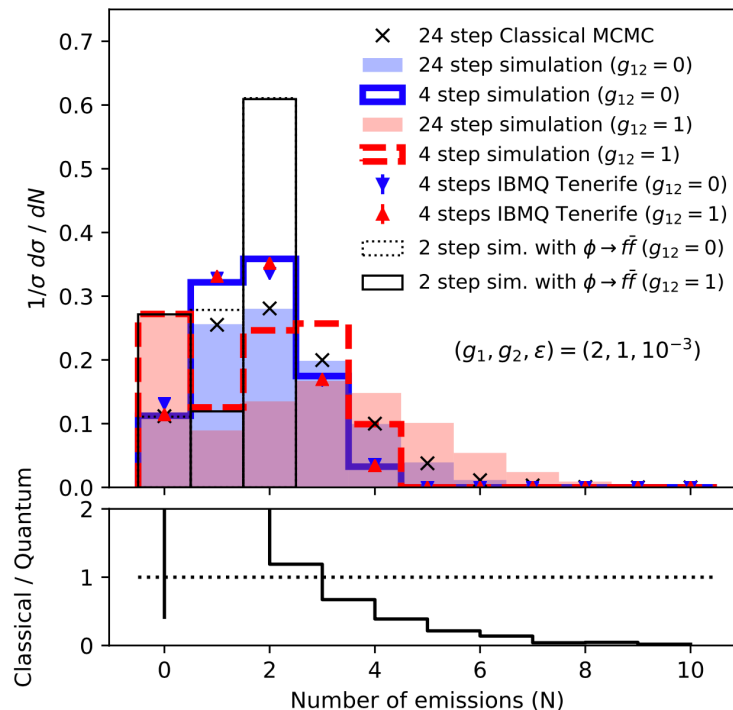
Register	Purpose	# of qubits
$ p\rangle$	Particle state	$3(N + n_I)$
$ h\rangle$	Emission history	$N \lceil \log_2(N + n_I) \rceil$
$ e\rangle$	Did emission happen?	1
$ n_\phi\rangle$	Number of bosons	$\lceil \log_2(N + n_I) \rceil$
$ n_a\rangle$	Number of f_a	$\lceil \log_2(N + n_I) \rceil$
$ n_b\rangle$	Number of f_b	$\lceil \log_2(N + n_I) \rceil$



no interference
with interference



angle of maximum emission



number of emissions

Quantum Generative Adversarial Networks

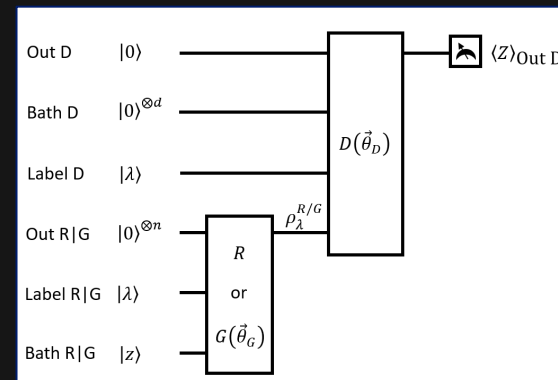
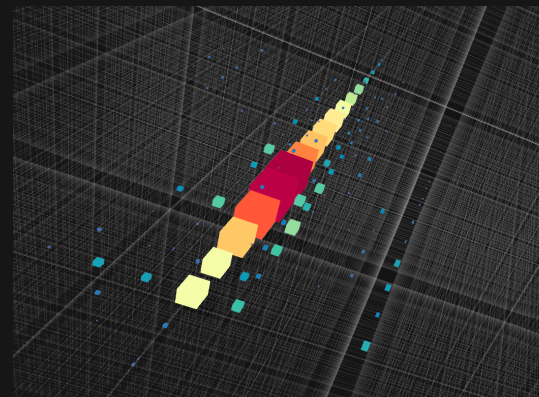
Use classical **Generative Adversarial Networks** to **simulate detector response**

- Replace Monte Carlo simulation

Quantum GAN can have **larger representational power**

- Different hybrid classical-quantum algorithms for generative models exist

Train a quantum GAN to generate **few-pixels image**

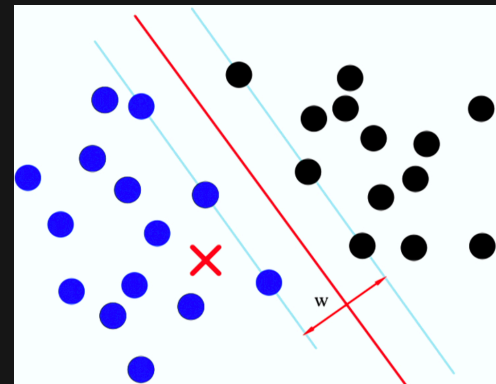


From Classics SVM to QSVM

Primal Problem

$$L_P = \underbrace{\frac{1}{2} \|\vec{w}\|^2}_{\text{minimize this to maximize the margin}} - \sum_{i \in T} \alpha_i \underbrace{[k_i(\vec{w} \cdot \Phi(\vec{x}_i) + b) - 1]}_{\text{constraints}}$$

Lagrange multipliers class labels $\{+1, -1\}$



From Classics SVM to QSVM

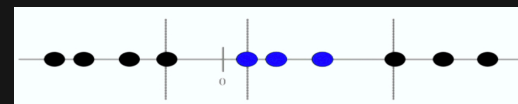
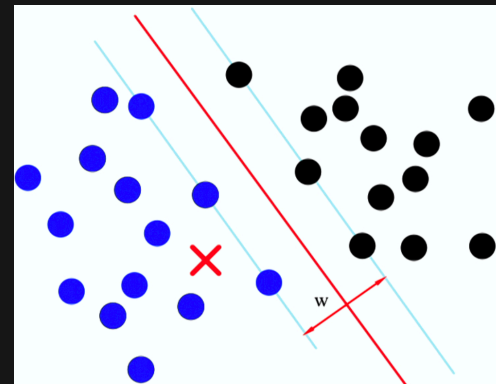
Primal Problem

$$L_P = \underbrace{\frac{1}{2} \|\vec{w}\|^2}_{\text{minimize this to maximize the margin}} - \sum_{i \in T} \alpha_i \underbrace{[k_i(\vec{w} \cdot \Phi(\vec{x}_i) + b) - 1]}_{\text{constraints}}$$

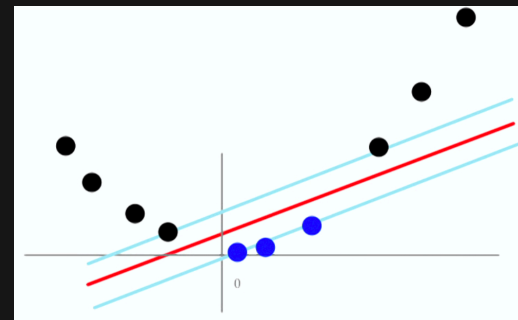
Lagrange multipliers
class labels $\{+1, -1\}$

Dual Problem is useful because dot products can
Be replaced by a **kernel** function

$$L_D(\alpha) = \sum_{i \in T} \alpha_i - \frac{1}{2} \sum_{i,j \in T} k_i k_j \alpha_i \alpha_j \underbrace{\Phi(\vec{x}_i) \cdot \Phi(\vec{x}_j)}_{K(\vec{x}_i, \vec{x}_j)}$$

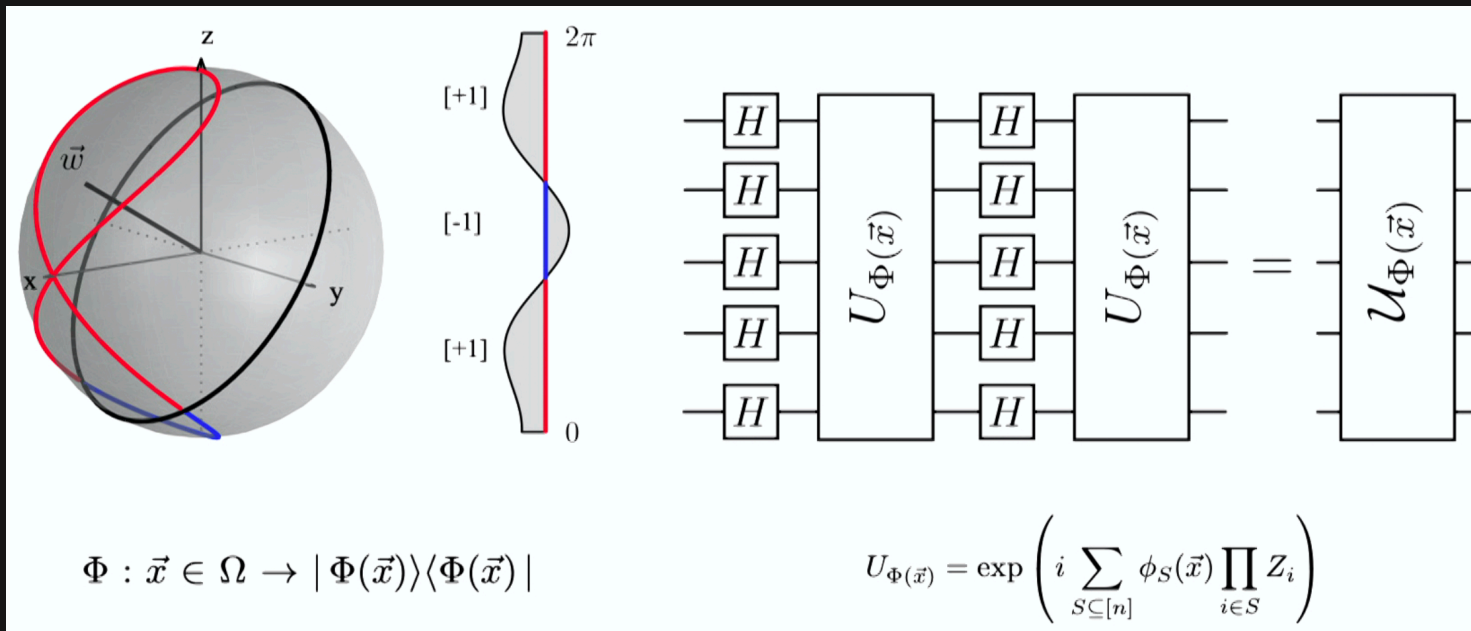


Choose the right feature map



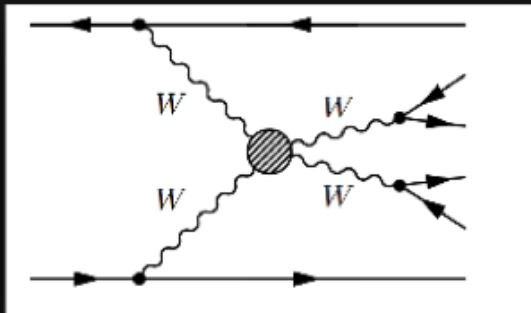
Supervised learning with quantum enhanced feature space

- I**
Map the data
- II**
Apply short-depth circuit
- III**
Measure in canonical basis
- IV**
Assign label



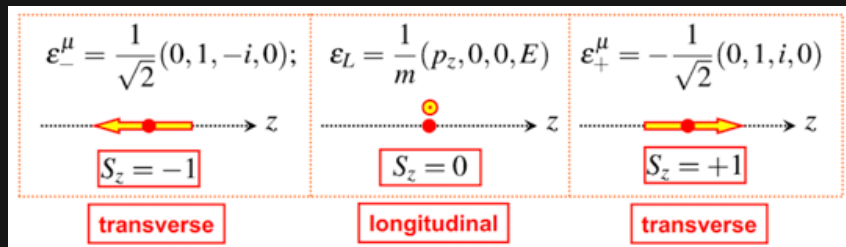
Michele Grossi

Vector Boson Scattering



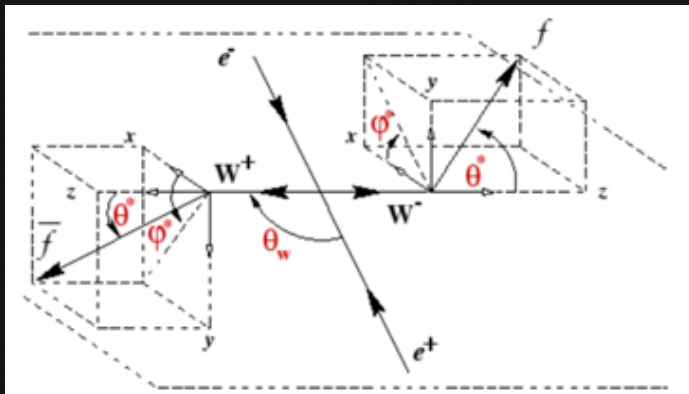
- Unitarity at high energies requires presence of the SM Higgs boson
- Sensitivity grows with energy of vector bosons
- Self-interaction of heavy gauge bosons
- Search for anomalous quartic-gauge-boson couplings

- The cross-section and angular distribution of longitudinal polarisations are particularly sensitive to beyond standard model (BSM) physics
- Boson polarisation can be measured as angular distributions of particles produced in the decay process



$$\frac{1}{\sigma} \frac{d\sigma}{d\cos\theta}(W^\pm \rightarrow l^\pm \nu) = \frac{3}{4} f_0 \sin^2 \theta + \frac{3}{8} f_R (1 \pm \cos \theta)^2 + \frac{3}{8} f_L (1 \mp \cos \theta)^2$$

Vector Boson Scattering

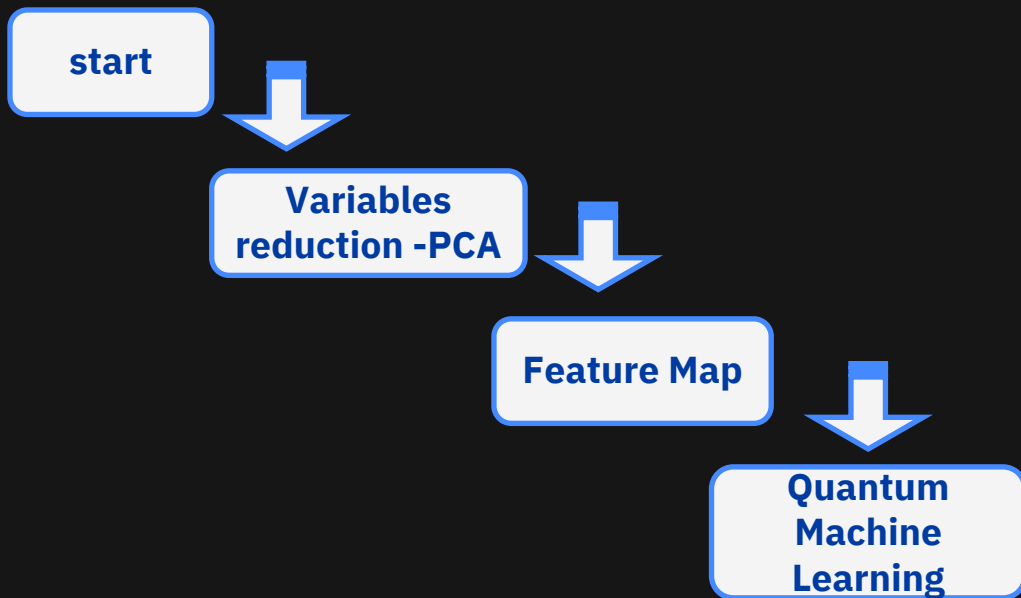


- Events with first solution closer to truth solution (1)
- Events with second solution closer to truth (0)
- Events with negative discriminant are discarded

$$\underbrace{(p_{\ell L}^2 - E_{\ell}^2) p_{\nu L}^2}_{a} + \underbrace{(m_W^2 p_{\ell L} + 2 p_{\ell L} \vec{p}_{\ell T} \vec{p}_{\nu T}) p_{\nu L}}_b + \underbrace{\frac{m_W^4}{4} + (\vec{p}_{\ell T} \vec{p}_{\nu T})^2 + m_W^2 \vec{p}_{\ell T} \vec{p}_{\nu T} - E_{\ell}^2 \vec{p}_{\nu T}^2}_c = 0$$

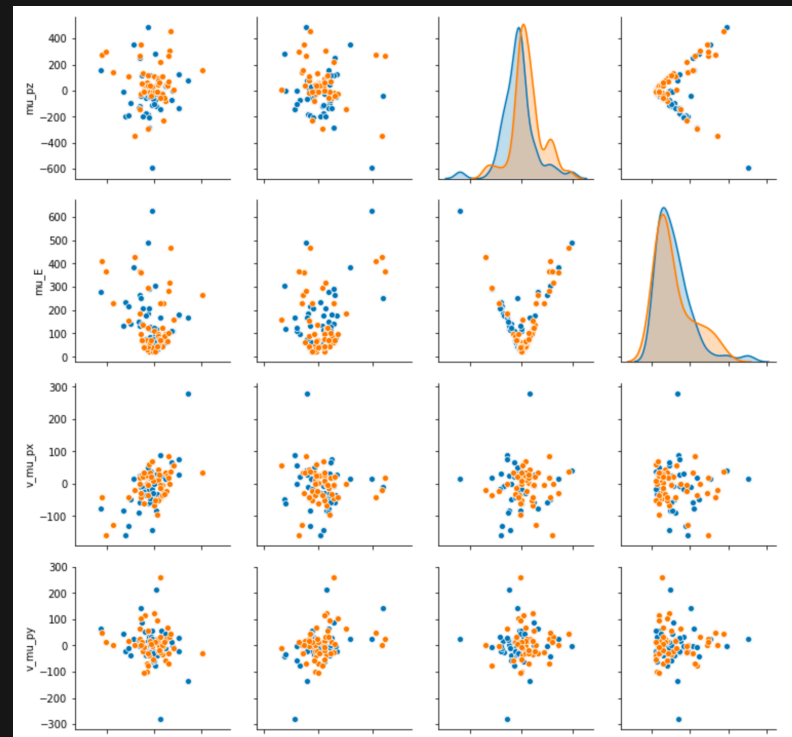
$$p_{\nu L} = \frac{-b \pm \sqrt{b^2 - 4ac}}{2a}$$

Vector Boson Scattering – QSVM Flow

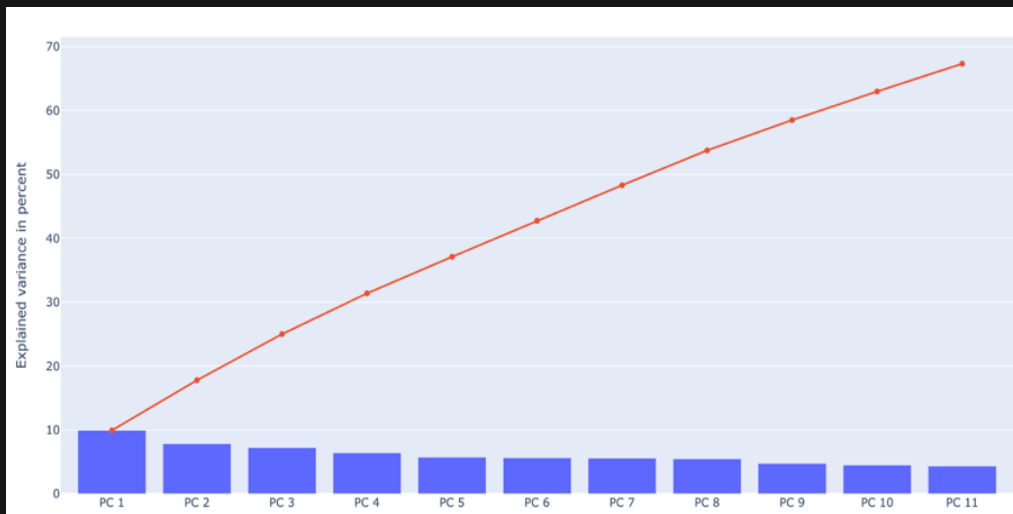


Given the current status of NISQ device we are limited in the number of features to map for our problem. To reduce the dimension of this space we apply the PCA analysis. We need to find balance between number of PCA and total variance explained

Ref: M. Grossi et al in preparation

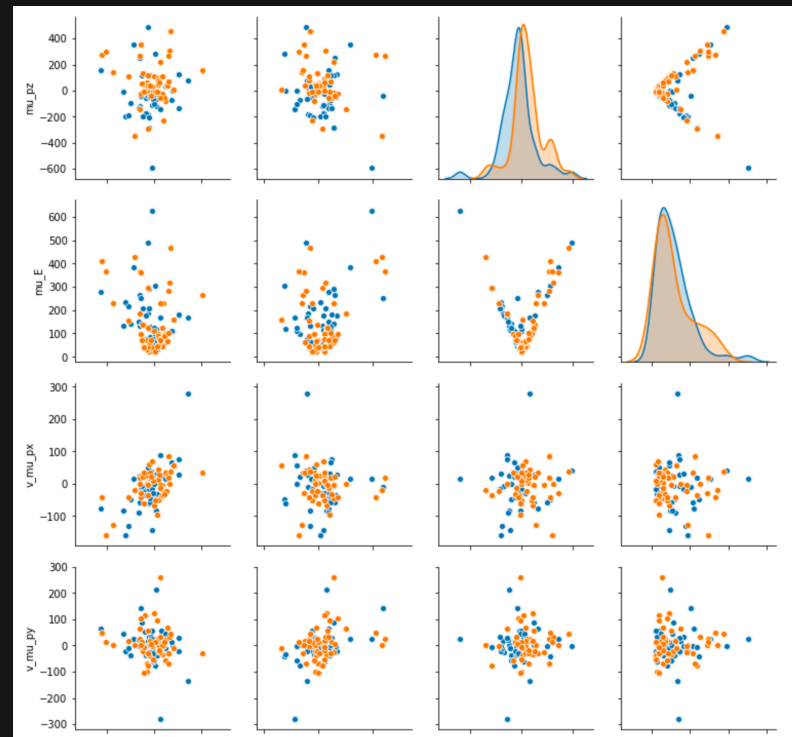


Vector Boson Scattering – QSVM Flow

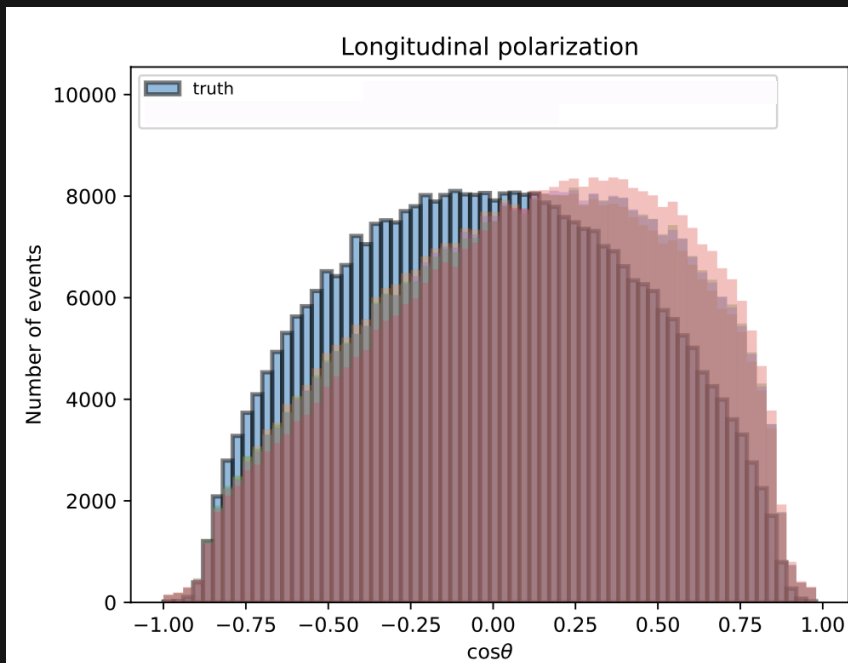
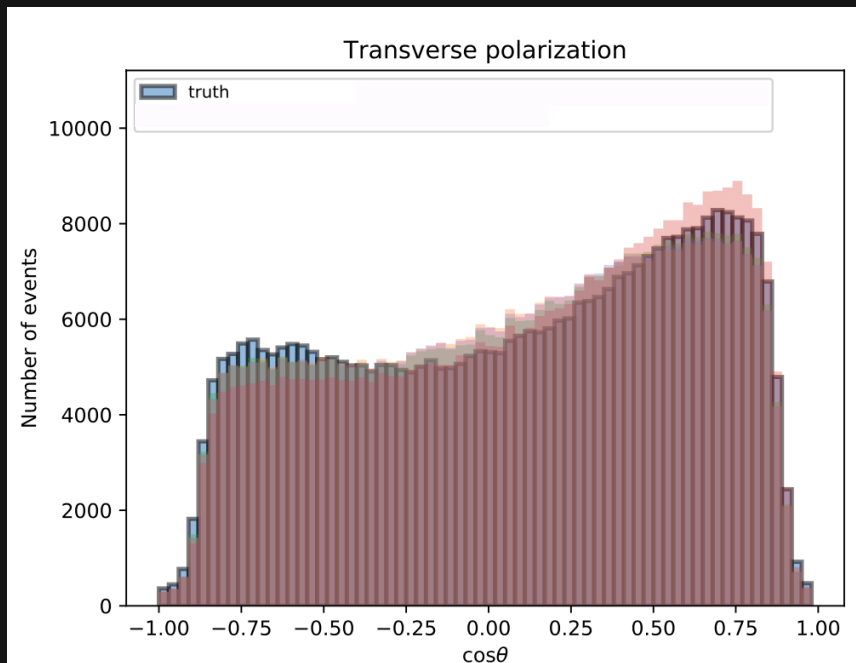


Given the current status of NISQ device we are limited in the number of features to map for our problem. To reduce the dimension of this space we apply the PCA analysis. We need to find balance between number of PCA and total variance explained

Ref: M. Grossi et al in preparation



Vector Boson Scattering – Angular Distribution



Thank you!

*There is a long road ahead,
but quantum algorithms are
very promising for modelling
high energy scattering
processes.*

Michele Grossi

IBM Quantum

An abstract graphic on the right side of the slide, consisting of several overlapping, elongated, rounded rectangular shapes in a light blue color, arranged in a vertical, slightly tilted stack.