

Understanding template methods for top polarisation measurements

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TOP 2022, Durham, September 6th 2022

Brief outline

- What is top polarisation?
- (A) Template method
- QFT is quantum!
- Right vs wrong templates
- Example for $t\bar{t}$

Based on:

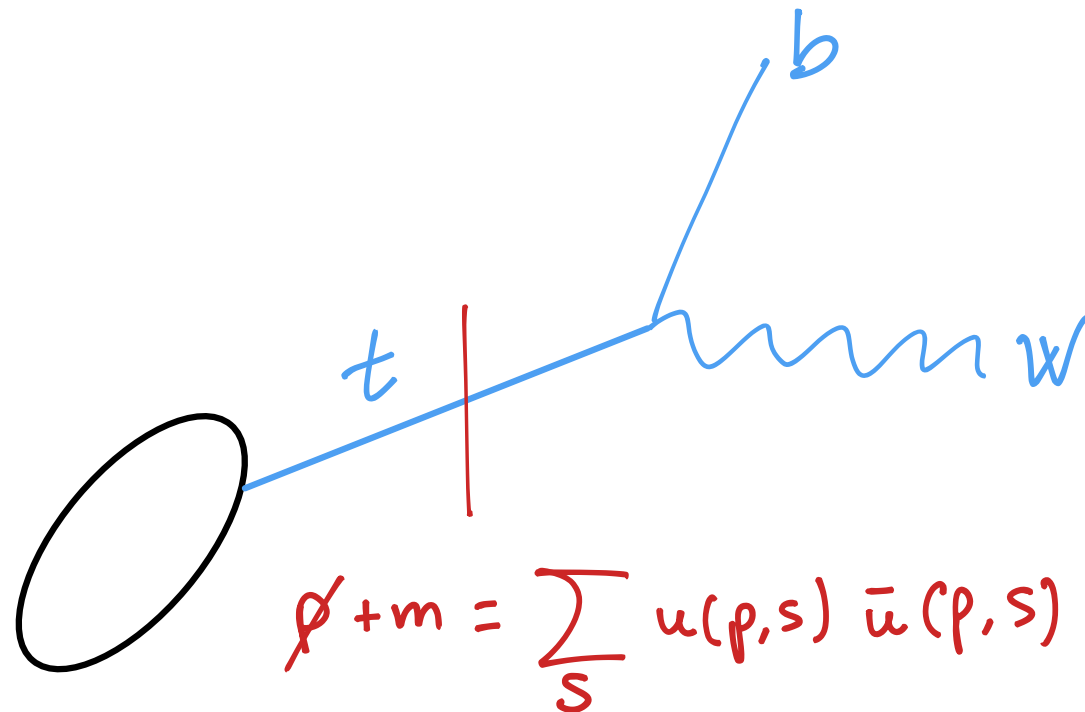
- JAAS, Fiolhais, Martín-Ramiro, Moreno, Onofre, *A template method to measure the $t\bar{t}$ polarisation*, 2111.10394
- JAAS, *Crafting polarisations for top, W and Z*, 2208.00424
- New material and discussions to appear on the proceedings

What is top
polarisation?

What is top polarisation?

The top quark has a mean lifetime of 10^{-25} s. It appears as an intermediate resonance in processes.

Then, what do we mean by 'top polarisation'?



In the NWA, the amplitudes can be decomposed into production $\times$ decay, summed over spins.

$$\mathcal{M} = \sum_{s=\pm 1/2} \mathcal{P}_s \mathcal{D}_s$$

This is the top polarisation that is measured by angular distributions

The production mechanism can be characterised by a spin density operator ρ whose matrix elements are related to production amplitudes $\mathcal{P}_s$

$$\rho_{ss'} = \text{PS} \times \sum_{\text{subprocesses}} \text{PDF} \times \mathcal{P}_s \mathcal{P}_{s'}^*$$

The goal of polarisation measurements is to measure matrix elements of ρ .

For the top quark,

$$\rho = \frac{1}{2} \begin{pmatrix} 1 + P_z & P_x - iP_y \\ P_x + iP_y & 1 - P_z \end{pmatrix} .$$

The polarisation vector $P_i = 2 \langle S_i \rangle$ satisfies $|\vec{P}| \leq 1$ for a physical ρ .

Template
method

Template method

If we consider the charged lepton from the top decay to measure the polarisation, its angular distribution [in top rest frame] is

$$\frac{1}{\sigma} \frac{d\sigma}{d\Omega} = \frac{1}{4\pi} \left[1 + \alpha_\ell \vec{P} \cdot \hat{p}_\ell \right]$$

Not templatable

Imagine we want to measure P_z . Then, we can integrate over ϕ to obtain

$$\frac{1}{\sigma} \frac{d\sigma}{d\cos\theta} = \frac{1}{2} [1 + \alpha_\ell P_z \cos\theta]$$

which can be written as

Distribution
for $P_z = 1$

Distribution
for $P_z = -1$

$$\frac{1}{\sigma} \frac{d\sigma}{d\cos\theta} = a_+ \left[\frac{1}{\sigma} \frac{d\sigma}{d\cos\theta} \right]_+ + a_- \left[\frac{1}{\sigma} \frac{d\sigma}{d\cos\theta} \right]_-$$

with

$$a_+ = (1 + P_z)/2 = \rho_{\frac{1}{2}\frac{1}{2}} = \sigma_+/\sigma$$

$$a_- = (1 - P_z)/2 = \rho_{-\frac{1}{2}-\frac{1}{2}} = \sigma_-/\sigma$$

Template
expansion

Template method

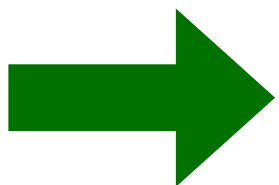
Discussion

- The template fit is well-founded on a theoretical equality

$$\frac{1}{\sigma} \frac{d\sigma}{d\cos\theta} = a_+ \left[\frac{1}{\sigma} \frac{d\sigma}{d\cos\theta} \right]_+ + a_- \left[\frac{1}{\sigma} \frac{d\sigma}{d\cos\theta} \right]_-$$

and therefore the fit coefficients **have a theoretical interpretation**

- Side comment: not clear on which theoretical equality the template fit in ATLAS polarisation measurement **2202.II382** is based.



The theoretical interpretation of the fit coefficients is unclear to me

Template method

Discussion

- The template expansion looks classical:

Sum over $+$, $-$ without interference between polarisations.

- This is possible for the $\cos \theta$ distribution and after integration on ϕ .
Also for the total cross section: σ_+ and σ_- are meaningful.
- Ignoring interference in other distributions leads to results that are not correct. Example: A_{FB}^ℓ

Template method

Reconstruction, detector effects and kinematical cuts make things more complicated.

To see this easily, we rewrite the template expansion [multiply by σ]

$$\frac{1}{\sigma} \frac{d\sigma}{d\cos\theta} = a_+ \left[\frac{1}{\sigma} \frac{d\sigma}{d\cos\theta} \right]_+ + a_- \left[\frac{1}{\sigma} \frac{d\sigma}{d\cos\theta} \right]_- \quad \longrightarrow \quad \frac{d\sigma}{d\cos\theta} = \frac{d\sigma_+}{d\cos\theta} + \frac{d\sigma_-}{d\cos\theta}$$

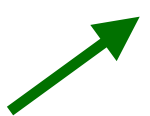
With cuts, etc. this is modified to

The diagram illustrates the modification of the template expansion when cuts and interference are considered. It features two orange boxes: 'Quantities after cuts' and 'Interference'. Three green arrows point from 'Quantities after cuts' to the first three terms of the equation $\frac{d\bar{\sigma}}{d\cos\theta} = \frac{d\bar{\sigma}_+}{d\cos\theta} + \frac{d\bar{\sigma}_-}{d\cos\theta} + \dots$. A fourth green arrow points from 'Interference' to the ellipsis '...', which is circled in orange.

$$\frac{d\bar{\sigma}}{d\cos\theta} = \frac{d\bar{\sigma}_+}{d\cos\theta} + \frac{d\bar{\sigma}_-}{d\cos\theta} + \dots$$

Template method

In order to cast that sum in terms of normalised quantities, one can define 'efficiencies' $\varepsilon = \bar{\sigma}/\sigma$, $\varepsilon_{\pm} = \bar{\sigma}_{\pm}/\sigma_{\pm}$ and obtain a template expansion

$$\varepsilon \left[\frac{1}{\bar{\sigma}} \frac{d\bar{\sigma}}{d\cos\theta} \right] = \varepsilon_+ a_+ \left[\frac{1}{\bar{\sigma}_+} \frac{d\bar{\sigma}_+}{d\cos\theta} \right] + \varepsilon_- a_- \left[\frac{1}{\bar{\sigma}_-} \frac{d\bar{\sigma}_-}{d\cos\theta} \right] + \Delta_{int}$$


Do experiments ever consider this term?

The term Δ_{int} arises because:

- The integration over ϕ is **not complete** because of cuts.
- The ϕ angle in the top rest frame **is not well reconstructed** and then, integration is not actually over the true ϕ .

In $t\bar{t}$ with dilepton decay, it was seen in [2111.10394](#) that the second effect is way more important.

QFT is
quantum

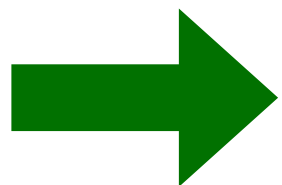
QFT is quantum

The interference term is important for precision measurements.
Fortunately, a SM calculation may be sufficient.

Proposed method: [JAAS, Fiolhais, Martín-Ramiro, Moreno, Onofre, 2111.10394](#)

- Generate statistically independent templates (MC) and SM pseudo-data (also with MC)
- Use SM values for polarisation coefficients
- Use efficiencies from Monte Carlo
- Compute Δ_{int} in the SM as

$$\Delta_{\text{int}} = \varepsilon \left[\frac{1}{\bar{\sigma}} \frac{d\bar{\sigma}}{d\cos\theta} \right] - \varepsilon_+ a_+ \left[\frac{1}{\bar{\sigma}_+} \frac{d\bar{\sigma}_+}{d\cos\theta} \right] - \varepsilon_- a_- \left[\frac{1}{\bar{\sigma}_-} \frac{d\bar{\sigma}_-}{d\cos\theta} \right]$$

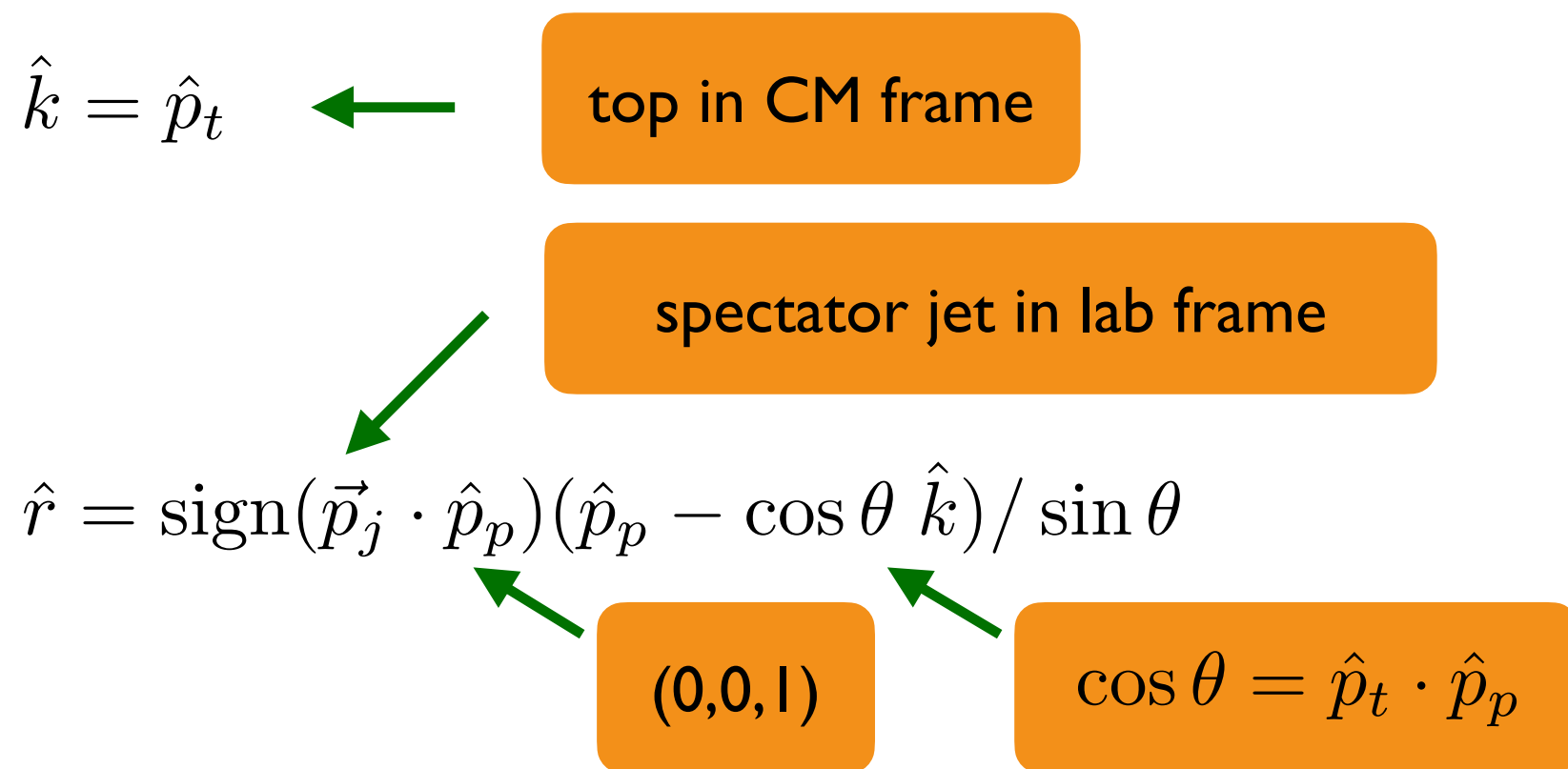


Checked in tt production that Δ_{int} computed in the SM is enough even with large [experimentally excluded] CMDM

QFT is quantum

Simple example for single top production

Axes: helicity (K), transverse (R) and normal (N)

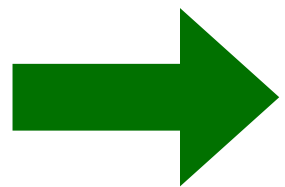


$$\hat{n} = \hat{k} \times \hat{r}$$

QFT is quantum

Simple example for single top production

No detector simulation, working at parton level but applying cuts $p_T \geq 30$ GeV for lepton and b from top decay.



This is sufficient to have large distortions on templates

From fit		K			R			N		
		a_+	a_-	P	a_+	a_-	P	a_+	a_-	P
	True	0.159	0.841	-0.682	0.615	0.385	0.230	0.501	0.499	0.002
	no Δ_{int}	0.160	0.840	-0.680	0.580	0.420	0.160	0.500	0.500	0.000
	with Δ_{int}	0.160	0.840	-0.680	0.616	0.384	0.232	0.500	0.500	0.000

At this level, the difference already is 1/4 of the uncertainty in ATLAS measurement

Right vs wrong
templates

Right vs wrong templates

Q: Is any MC sample

with correct top decay distribution

appropriate as a template?

A: No

To see this subtle point, let us go back to

$$\frac{d\sigma}{d\cos\theta} = \frac{d\sigma_+}{d\cos\theta} + \frac{d\sigma_-}{d\cos\theta} \quad \xrightarrow{\text{cuts}} \quad \frac{d\bar{\sigma}}{d\cos\theta} = \frac{d\bar{\sigma}_+}{d\cos\theta} + \frac{d\bar{\sigma}_-}{d\cos\theta} + \dots$$

The σ_+ and σ_- on the r.h.s. are the + and - components of σ on the l.h.s.

If σ [which will be data in the fit] is SM-like, σ_+ and σ_- must be too.

Therefore, σ_+ and σ_- must be obtained with spin projectors at the matrix-element level.

Right vs wrong templates

Example: single top / helicity axis / cut $p_T \geq 30$ GeV for extra b

SM w/cut

Norm. distributions unchanged by cut

Cut keeps $\Delta_{\text{int}} = 0$

$$\varepsilon \left[\frac{1}{\bar{\sigma}} \frac{d\bar{\sigma}}{d\cos\theta} \right] = \varepsilon_+ a_+ \left[\frac{1}{\bar{\sigma}_+} \frac{d\bar{\sigma}_+}{d\cos\theta} \right] + \varepsilon_- a_- \left[\frac{1}{\bar{\sigma}_-} \frac{d\bar{\sigma}_-}{d\cos\theta} \right] + \Delta_{\text{int}}$$

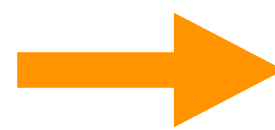
Fit returns $a_+ \varepsilon_+ / \varepsilon = 0.286$ $a_- \varepsilon_- / \varepsilon = 0.714$

SM eff: $\varepsilon = 0.292$

Templates with SM amplitudes
and spin projectors in matrix
element

$$\varepsilon_+ = 0.524$$

$$\varepsilon_- = 0.245$$



$$a_+ = 0.160$$

$$a_- = 0.840$$



Templates with SM production
kinematics and modified decay

$$\varepsilon'_+ = \varepsilon$$

$$\varepsilon'_- = \varepsilon$$



$$a_+ = 0.286$$

$$a_- = 0.714$$



Right vs wrong templates

Discussion

- This example has been chosen so that $\Delta_{\text{int}} = 0$ for simplicity and to have the template distributions unaltered from parton level.
- It clearly shows that to fit SM (on the l.h.s.) we need **SM polarised templates** on the r.h.s.
- Because single top kinematics is SM-like, SM polarised templates must be used in measurements.
- Using templates with **incorrect production mechanism**, leads to [quite?] different ε and therefore, to [quite?] biased results.
- In addition, the normalised distributions are different [not in the previous example].
- In the lucky event that measurements deviate significantly from the SM, we should think about closure tests using SM polarised templates.

Right vs wrong templates

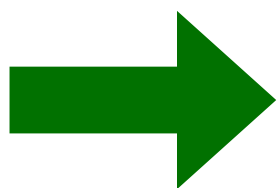
CAR: Custom Angle Replacement

JAAS, 2208.00424

Method to replace the decay kinematics of top, W and Z to tune desired polarisations [just noticed that the same method was used by ATLAS, PolManip].

In template measurements for top, it can be used to **augment samples**. In the previous example:

- Take a small $t_L\bar{t}_L$ polarised sample [50 K] to parameterise production kinematics.
- Use CAR to obtain large [1 M] $t_L\bar{t}_L, t_R\bar{t}_R$ statistically-independent samples with new top decays in less than a minute.



Especially useful for multi-top production, generation beyond LO, etc.

Example for tt

Example for $t\bar{t}$

Top pair  two polarisations $t_L\bar{t}_L, t_L\bar{t}_R, t_R\bar{t}_L, t_R\bar{t}_R$

JAAS, Fiolhais, Martín-Ramiro, Moreno, Onofre, 2111.10394

- Simulation with Delphes.
- Reconstruction of dilepton decay by scan over neutrino momenta and fit of top and W masses.
- Three axes: K (helicity), R (transverse) and N (normal)
- No systematics study. A comparison with unfolding methods should be made by experiments.
- Background not considered: it is quite small for the dilepton decay [different flavour leptons / cut on same-flavour dilepton invariant mass]
- Fit results obtained by fluctuating templates assuming # events corresponding to 36.1 fb^{-1} .

Example for $t\bar{t}$

As mentioned, there are two effects: distortion of the [normalised] templates and different efficiency factors for the different samples

Efficiency
factors



Sample		ε	
SM		0.174	
CMDM		0.174	
Template	ε		
	K-axis	R-axis	N-axis
LL	0.177	0.175	0.178
RR	0.178	0.176	0.178
LR	0.160	0.188	0.171
RL	0.182	0.161	0.171

Values are similar but it is crucial to plug in the correct ones!

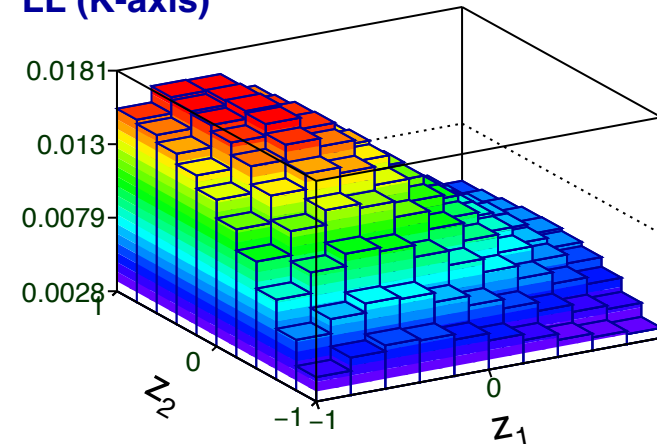
Example for tt

Also: distortion of [normalised] templates. Example: K axis

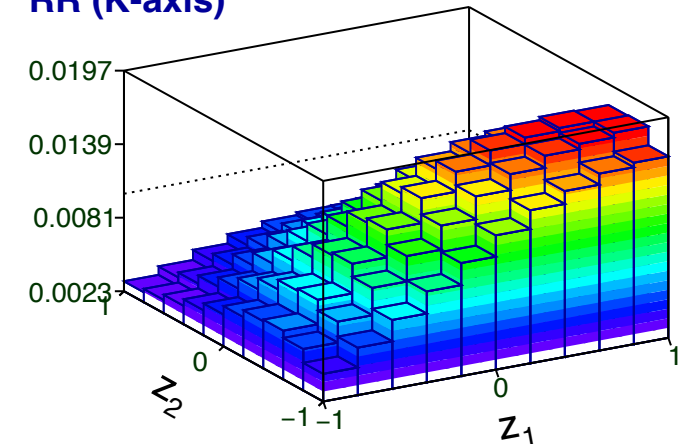
Normalised templates for
K axis after simulation

$$z_1 = \cos \theta_{\ell+} \quad z_2 = \cos \theta_{\ell-}$$

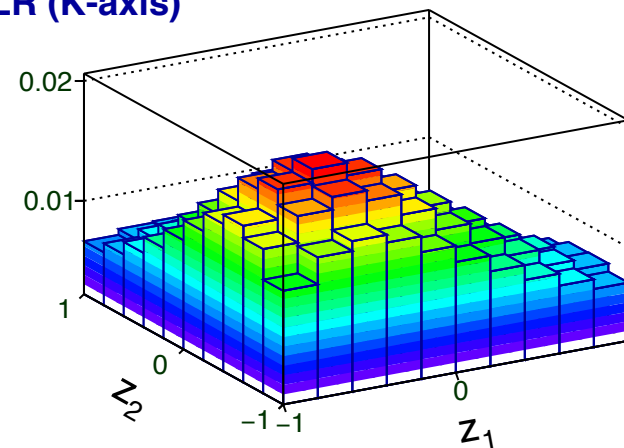
LL (K-axis)



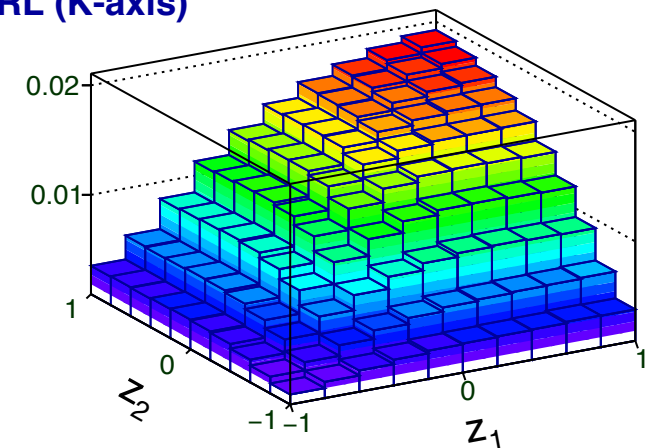
RR (K-axis)



LR (K-axis)



RL (K-axis)



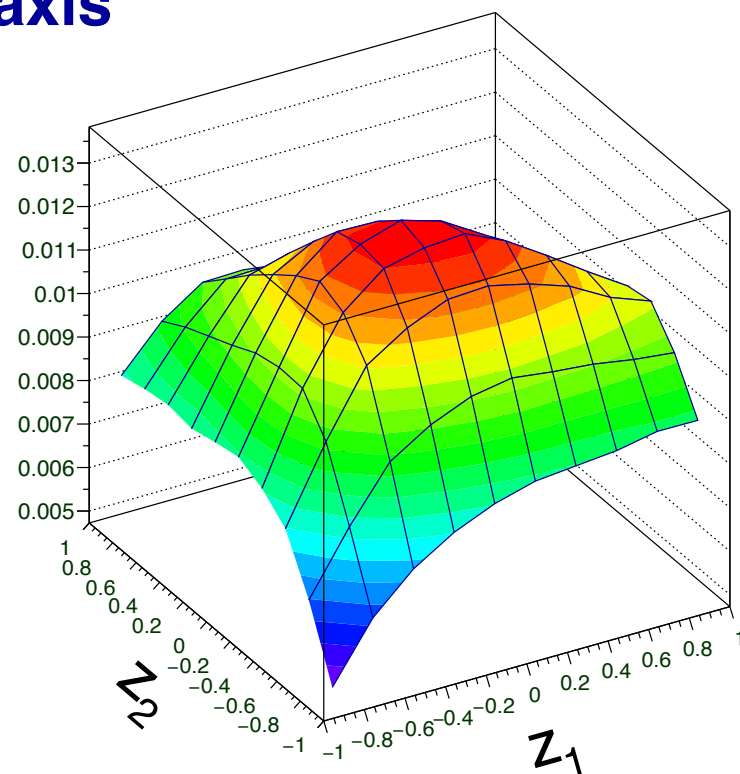
Example for $t\bar{t}$

The interference term Δ_{int} is small but not negligible

SM distribution for K axis
after simulation

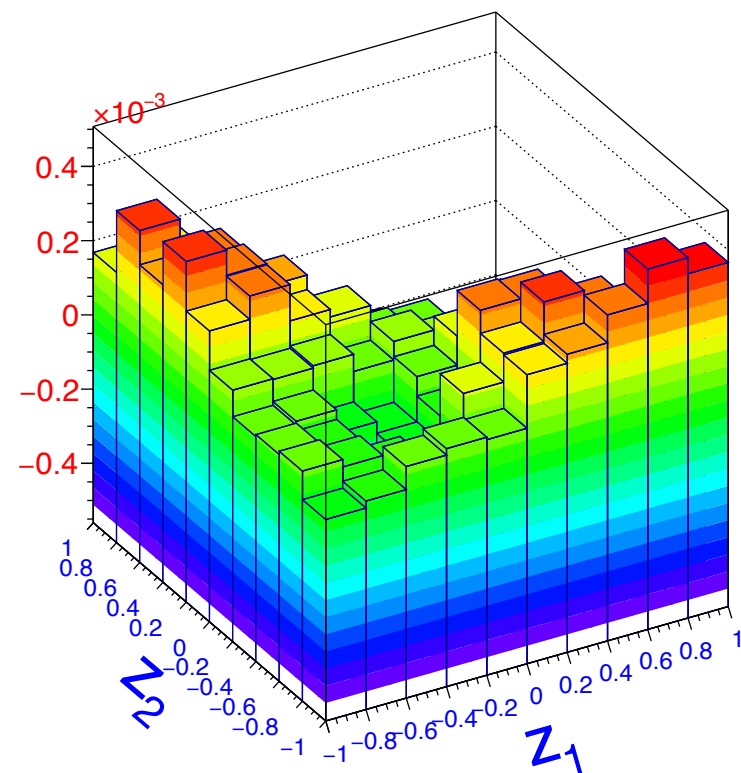
$$z_1 = \cos \theta_{\ell+} \quad z_2 = \cos \theta_{\ell-}$$

K-axis



Δ_{int} for K axis after
simulation

K-axis



Omitting Δ_{int} we get results with a bias of the order of the statistical uncertainty or larger

Example for $t\bar{t}$

Fit results with statistical uncertainties for 36.1 fb⁻¹

K	SM		CMDM	
	Prediction	Fit	Prediction	Fit
a_{LL}	0.335 ± 0.001	0.337 ± 0.006	0.349 ± 0.001	0.350 ± 0.006
a_{RR}	0.336 ± 0.003	0.330 ± 0.005	0.349 ± 0.001	0.339 ± 0.005
a_{LR}	0.165 ± 0.003	0.167 ± 0.007	0.151 ± 0.001	0.175 ± 0.007
a_{RL}	0.165 ± 0.002	0.160 ± 0.004	0.151 ± 0.001	0.131 ± 0.004
C_{kk}	0.340 ± 0.002	0.340 ± 0.019	0.394 ± 0.004	0.383 ± 0.019
P_t	0.001 ± 0.002	-0.014 ± 0.008	0.000 ± 0.001	-0.058 ± 0.008
$P_{\bar{t}}$	0.001 ± 0.002	0.000 ± 0.008	0.001 ± 0.002	0.033 ± 0.008
R	SM		CMDM	
	Prediction	Fit	Prediction	Fit
a_{LL}	0.258 ± 0.001	0.254 ± 0.006	0.290 ± 0.002	0.291 ± 0.006
a_{RR}	0.259 ± 0.002	0.264 ± 0.006	0.289 ± 0.002	0.290 ± 0.006
a_{LR}	0.242 ± 0.001	0.236 ± 0.006	0.210 ± 0.001	0.210 ± 0.006
a_{RL}	0.241 ± 0.002	0.241 ± 0.006	0.211 ± 0.001	0.201 ± 0.006
C_{rr}	0.036 ± 0.002	0.041 ± 0.019	0.159 ± 0.002	0.170 ± 0.019
P_t	0.0004 ± 0.0005	0.015 ± 0.010	-0.001 ± 0.004	-0.011 ± 0.010
$P_{\bar{t}}$	0.002 ± 0.002	0.006 ± 0.010	-0.001 ± 0.003	0.008 ± 0.009
N	SM		CMDM	
	Prediction	Fit	Prediction	Fit
a_{LL}	0.333 ± 0.001	0.329 ± 0.004	0.358 ± 0.001	0.363 ± 0.004
a_{RR}	0.334 ± 0.002	0.329 ± 0.004	0.358 ± 0.001	0.352 ± 0.004
a_{LR}	0.166 ± 0.001	0.164 ± 0.004	0.142 ± 0.0003	0.138 ± 0.004
a_{RL}	0.167 ± 0.002	0.169 ± 0.004	0.142 ± 0.001	0.136 ± 0.004
C_{nn}	0.336 ± 0.002	0.325 ± 0.010	0.433 ± 0.002	0.442 ± 0.010
P_t	0.002 ± 0.001	0.005 ± 0.009	-0.001 ± 0.002	-0.014 ± 0.009
$P_{\bar{t}}$	0.000 ± 0.002	-0.005 ± 0.008	0.000 ± 0.001	-0.009 ± 0.009

Example for $t\bar{t}$

Discussion

- The template method presented allows to access quantities that are not directly measured yet [individual a 's].
- The fit converges well, the only deviations found may have a statistical origin [samples have ~ 200 K events]
- The calculation of Δ_{int} is valid not only for the SM but also for a sample with a large CMDM that yields quite different polarisation coefficients.

End

Top Quark Production Asymmetries A_{FB}^t and A_{FB}^ℓ Edmond L. Berger,² Qing-Hong Cao,¹ Chuan-Ren Chen,² Jiang-Hao Yu,³ and Hao Zhang^{2,4}¹*Department of Physics and State Key Laboratory of Nuclear Physics and Technology, Peking University, Beijing, 100871, China*²*High Energy Division, Argonne National Laboratory, Argonne, Illinois 60439, USA*³*Department of Physics and Astronomy, Michigan State University, East Lansing, Michigan 48824, USA*⁴*Illinois Institute of Technology, Chicago, Illinois 60616-3793, USA*

(Received 23 November 2011; published 14 February 2012)