

UNIVERSITY **OF SUSSEX**

THE ROYAL SOCIETY

Gobol EFT Fits

Top2022 Durham 6/9/2022

Josh McFayden, on behalf of the ATLAS and CMS collaborations

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Review Article Published: 03 October 2019

Multi-messenger astrophysics

Péter Mészáros 🖂, Derek B. Fox, Chad Hanna & Kohta Murase

Nature Reviews Physics 1, 585–599 (2019) Cite this article

6156 Accesses | 49 Citations | 20 Altmetric | Metrics







Multi-messenger particle physics









Multi-messenger particle physics

Global EFT approaches offer a multi messenger approach to searching for new physics









Single process one operator at a time











- Single process one operator at a time
- Single process multiple operators

Josh McFayden | Top2022 | 6/9/2022





ATLAS

^-4

Λ-2

Λ-4

Λ⁻²

 $C_{tq}^{(8)}$

CtG

√s = 13 TeV, 139 fb⁻¹

-0.5









- Single process one operator at a time
- Single process multiple operators
- Multiple processes single operators



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Next steps?

- A more global Top fit
- A fully global fit









- Measurements of several processes
- Including from different experiments
- Many operators
- Complementarity improves sensitivity
- Currently performed mainly in pheno community

Using available experimental data/information



[arXiv:2012.02779]





SMEFit | Top Sector "global" fit SMEFiT framework "proof-of-concept" analysis of top quark production measurements from LHC 8 TeV and 13 TeV data







SMEFit | Top+Higgs+VV fit

- Global fit bounds more stringent for all EFT coefficients than either top- or Higgs-only fit
- Cross-talk of the top and Higgs data leads to significant improvement in sensitivity
 - E.g. $c_{\Phi t}$ and $c^{(-)}_{\Phi Q}$ bounds improved by ~factor 2
- Few operators unconstrained in top-only fit E.g. C_{\$\phi\$}
- In a Higgs-only fit a large number of EFT coefficients are poorly constrained
- Gain information in the global fit by breaking degeneracies
 - Sometimes in unexpected directions in the parameter space (such as for $c_{\phi G}$)





- Global fit using several datasets:
 - LEP+EWPO, Tevatron, LHC Run1, LHC Run 2
- Top-only fit
 - Shows impact of different datasets for individual fits





[arXiv:2012.02779]





LEP+EWPO, Tevatron, LHC Run1, LHC Run 2 Top-only fit

Global fit using several

Fitmaker | Global fit

- Shows impact of different datasets for individual fits
- Marginalised fit

datasets:





[arXiv:2012.02779]





- Global fit using several datasets:
- LEP+EWPO, Tevatron, LHC Run1, LHC Run 2
- Top-only fit
 - Shows impact of different datasets for individual fits
 - Marginalised fit
 - Detailed impact of different measurements to constrain fit directions for some operators



[arXiv:2012.02779]



Global fit

- Shows impact of different datasets
 - All data except top
 - Full global fit
 - Top data alone

In some cases worsened sensitivity More freedom in the fit





<u>[arXiv:2012.02779]</u>





- Some correlations between Top vs Higgs and EWKPO observables.
- Interesting to see e.g. Higgs and Top complementarity
- Fit to subset of operators:

 $\{C_{H\Box}, C_{HG}, C_{HW}, C_{HB}, C_{tH}, C_{bH}, C_{\tau H}, C_{\mu H} C_G \text{ and } C_{tG} \}$

Effect of including both Higgs and top data clearly reduces the allowed parameter space and reduces correlations







- Some correlations between Top vs Higgs and EWKPO observables.
- Interesting to see e.g. Higgs and Top complementarity
- Fit only 2 operators at a time:
 - D fits show what is driving the marginalised constraints
 - Shows contribution of individual measurements
 - STXS on its own not yet very sensitive
 - Including top data helps a lot











QLink to B-anomalies

- EFT already used in interpretation of B-anomalies
- Links between b and top:
 - B-anomalies require NP in semileptonic 4-fermion operators
 - Can also use SMEFT to include tops:

 $\bar{s}_L \gamma_\mu b_L \bar{\mu}_L \gamma^\mu \mu_L \longrightarrow \begin{array}{c} O_{lq}^{(1)} = \bar{Q} \gamma_\mu Q \bar{L} \gamma^\mu L, \\ O_{lq}^{(3)} = \bar{Q} \gamma_\mu \tau^I Q \bar{L} \gamma^\mu \tau^I L \end{array}$





[2012.10456]



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Top data particularly complementary wrt B + Zbb for some coefficients



[2012.10456]





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- Top data particularly complementary wrt B + Zbb for some coefficients
- CMS Top+X fit already gives constraints on semileptonic four-fermion operators

┛, ${}^{\mu} \tau^{I} L$



[TOP-19-001]



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- Top data particularly complementary wrt B + Zbb for some coefficients
- CMS Top+X fit already gives constraints on semileptonic four-fermion operators
- Need agreement on EFT flavour assumptions: [LHC EFT WG Meeting]







Global fit from experiments

- What about global fits from the experiments?
- Global "Top+X" fit seem from CMS already
- ATLAS has started looking at first steps to a global EFT fit including Higgs, EW and EW precision measurements
 - Builds on previous efforts to perform fits with multiple datasets (e.g. [ATL-PHYS-PUB-2021-010])
 - Adds LEP and SLC EW precision data for the first time
- LHC Top WG combining ATLAS+CMS data
- LHC EFT WG combining ATLAS+CMS+LEP
 - Including Top data

Wilson coefficient and operator		Affected process group		
	•	LEP/SLD EWPO	ATLAS Higgs	AT electi
$c_{H\Box}$	$(H^\dagger H) \Box (H^\dagger H)$		\checkmark	
c_G	$f^{abc}G^{a u}_\mu G^{b ho}_ u G^{c\mu}_ ho$		\checkmark	
c_W	$\epsilon^{IJK}W^{I u}_{\mu}W^{J ho}_{ u}W^{K\mu}_{ ho}$		\checkmark	
c_{HD}	$\left(H^{\dagger}D_{\mu}H ight)^{*}\left(H^{\dagger}D_{\mu}H ight)$		\checkmark	
c_{HG}	$H^\dagger H G^A_{\mu u} G^{A\mu u}$		\checkmark	
c_{HB}	$H^\dagger H B_{\mu u} B^{\mu u}$		\checkmark	
c_{HW}	$H^\dagger H W^I_{\mu u} W^{I\mu u}$		\checkmark	
c_{HWB}	$H^\dagger au^I H W^I_{\mu u} B^{\mu u}$	\checkmark	\checkmark	
C _{eH}	$(H^{\dagger}H)(\bar{l}_{p}e_{r}H)$		\checkmark	
C _{uH}	$(H^{\dagger}H)(\bar{q}Y_{u}^{\dagger}u\widetilde{H})$		\checkmark	
c_{tH}	$(H^{\dagger}H)(ar{Q}\widetilde{H}t)$		\checkmark	
C _{bH}	$(H^{\dagger}H)(ar{Q}Hb)$		\checkmark	
$c_{Hl}^{(1)}$	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}H)(\bar{l}\gamma^{\mu}l)$	\checkmark	\checkmark	
$c_{HI}^{(3)}$	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}^{I}H)(\bar{l}\tau^{I}\gamma^{\mu}l)$	\checkmark	\checkmark	
C _{He}	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}H)(\bar{e}\gamma^{\mu}e)$	\checkmark	\checkmark	
$c_{Ha}^{(1)}$	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}H)(\bar{q}\gamma^{\mu}q)$	\checkmark	\checkmark	
$c_{H_{q}}^{(3)}$	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}^{I}H)(\bar{q}\tau^{I}\gamma^{\mu}q)$	\checkmark	\checkmark	
C _{Hu}	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}H)(\bar{u}\gamma^{\mu}u)$	\checkmark	\checkmark	
CHd	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}H)(\bar{d}\gamma^{\mu}d)$	\checkmark	\checkmark	
$c_{HO}^{(1)}$	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}H)(\bar{Q}\gamma^{\mu}Q)$	\checkmark	\checkmark	
$c_{\mu 0}^{(3)}$	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}^{I}H)(\bar{Q}\tau^{I}\gamma^{\mu}Q)$	\checkmark	\checkmark	
пQ Снь	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}H)(\bar{b}\gamma^{\mu}b)$	\checkmark		
CHt	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}H)(\bar{t}\gamma^{\mu}t)$	\checkmark	\checkmark	
$\frac{1}{C_{tG}}$	$(\bar{O}\sigma^{\mu\nu}T^At)\tilde{H}G^A_{\mu\nu}$		\checkmark	
C _t w	$(\bar{O}\sigma^{\mu\nu}t)\tau^I \tilde{H} W^I_{}$		\checkmark	
C _t R	$(\bar{Q}\sigma^{\mu\nu}t)\widetilde{H}B_{\mu\nu}$		\checkmark	
$\frac{c_{ll}}{c_{ll}}$	$\frac{(\bar{l}\gamma_{\mu}l)(\bar{l}\gamma^{\mu}l)}{(\bar{l}\gamma^{\mu}l)}$	\checkmark		

[ATL-PHYS-PUB-2022-037]





- Sensitivity to numerous EFT operators
- Complementarity between different observables





L-PHYS-PUB-2022-037





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ATL-PHYS-PUB-2022-037





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ATL-PHYS-PUB-2022-037]





- Sensitivity to numerous EFT operators
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//.//	
	-0.05
:	C (1) =U. I
	C =0.2
:	c ^{eu} (31) =0.2
:	aa
	C (38) =1.0
	aa
:	C ¹ (18) =1.0
	0 1.0
	c ⁰ –1 0
	^م –10
	$C^{QQ}_{(8)} = 2.0$
4 4	v Qu
	Y,
16. 2	P



Single operator fits used to assess impact of different datasets

For example O_{Hq} :



<u>[L-PHYS-PUB-2022-037]</u>







Start by looking at Higgs+EW fit

- Linear and quadratic parameterisations
- Expected contribution to sensitivity shown





Start by looking at Higgs+EW fit

- Linear and quadratic parameterisations
- Expected contribution to sensitivity shown

Full fit to EWPO+EW+Higgs

- Results largely compatible with SM
- σ^{0}_{had} , $H \rightarrow \gamma \gamma$, Γ_{Z} , A_{fb} , and $gg \rightarrow H$ giving largest constraints
- In some cases significant constraints come EWPD with important contributions from VH, VBF Higgs and VV production
 - Contributions of the Higgs and EW measurements expected to become more important with larger datasets





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6/9/2022

ATLAS+CMS combination ATLAS \sqrt{s} = 13 TeV, 139 fb $\frac{d\sigma}{dp_T^Z}$ [fb.G6 3I + 4I combination of top measurements and EFT interpretation

Using unfolded data:





p_T^{t,h} [GeV]

ATLAS+CMS combination of top measurements and EFT interpretation

- Using unfolded data
- Using reconstruction-level information:

[LHCtopWG: Ravina+Skovpen]





CHARGEFT WG Efforts

LHC EFT WG already staring work on wider fits

- WG to provide EFT recommendations
- Currently looking at ATLAS+CMS fitting exercise (latest update <u>here</u>)
- Includes Top data!
- Workflow defined
- Rivet used to extract parameterisation
- Input based on YAML
- Fit performed with Roofit



CHARGEFT WG Efforts

- Perform eigenvector decomposition and PCA to identify operator groups
 - Small correlation between groups
- Able to identify expected contributions of each measurement
- Fitting exercise in a good shape Everything in place to perform the different steps in a flexible manner









Coordination/Organisational Challenges General approach

- Combine input measurements & interpret or directly interpret simultaneous fit of measurements?
- Detector level fits vs unfolded data?

Practical difficulties

- Different statistical methods (IBU vs FBU, PL vs toys, ...)
- Proper treatment of statistical and systematic correlations
- Measurements delivered on different timelines
- Interpretations: different assumptions on "backgrounds" \rightarrow EFT effects - Hard without coordination!

What to store?

- Store bootstrap replicas for data estimate stat. correlations
- Stat only results
- Impact of systematic uncertainties:
 - Full likelihood (at least: full ranking and NPs correlations), cov. matrix for each unc. incl. stat only. Metadata.

Signal model

SMEFT@LO or @NLO? Which operators? Linear/quadratic terms. EFT uncs & validity constraints.








- Quite comprehensive global fits exist from the theory/pheno community • The importance of including Top data in these is clear
- global fits
- CMS has published a "global" Top+X fit
- ATLAS has a recent PUB note on a "global" EWPO/Higgs/EW fit
- LHC Top WG exercise looking at unfolded at detector level combinations
- LHC EFT WG fitting exercise includes LEP+ATLAS+CMS EWPO/Higgs/EW/Top fit
- Many coordination/organisational challenges



Now it's time to think about next steps for adding Top data to experimental

Much of the technical machinery is there and existing fits can serve as starting point









Back-ups







- Complicated interplay of many "signals" all allowed to "float" at
 Once.
- Some small processes get very large effective normalisation factors
- One of the complexities of such a fit!



CMS Top+X "global" fit



ATLAS Global EFT fit

Input datasets

Decay channel	Target Production Modes	\mathcal{L} [fb ⁻¹]	Ref.				
$H \rightarrow \gamma \gamma$	ggF, VBF, WH, ZH, tīH, tH	139	[10]				
$H \rightarrow ZZ^*$	ggF, VBF, WH , ZH , $t\bar{t}H(4\ell)$	139	[11]				
$H \rightarrow WW^*$	ggF, VBF	139	[12]	-			
$H \rightarrow \tau \tau$	ggF, VBF, WH, ZH, $t\bar{t}H(\tau_{had}\tau_{had})$	139	[13]	Observable	Measurement	Prediction	Rati
	WH, ZH	139	[14–16]	Γ_Z [MeV]	2495.2 ± 2.3	2495.7 ± 1	$0.9998 \pm$
$H \rightarrow b \bar{b}$	VBF	126	[17]	R^0_{ℓ}	20.767 ± 0.025	20.758 ± 0.008	$1.0004 \pm$
	tĪH	139	[18]	R_c^0	0.1721 ± 0.0030	0.17223 ± 0.00003	$0.999 \pm$
				R_b^0	0.21629 ± 0.00066	0.21586 ± 0.00003	$1.0020 \pm$
				$A_{ m FB}^{0,\ell}$	0.0171 ± 0.0010	0.01718 ± 0.00037	$0.995 \pm$
				$A_{\rm FB}^{0,c}$	0.0707 ± 0.0035	0.0758 ± 0.0012	$0.932 \pm$
				$A_{\rm FB}^{0,b}$	0.0992 ± 0.0016	0.1062 ± 0.0016	$0.935 \pm$
				$\sigma_{\rm had}^{0}$ [pb]	41488 ± 6	41489 ± 5	$0.99998 \pm$

Process	Important phase space requirements	Observable	\mathcal{L} [fb ⁻¹]	Ref.
$pp \rightarrow e^{\pm} \nu \mu^{\mp} \nu$	$m_{\ell\ell} > 55 \text{GeV}, p_{\mathrm{T}}^{\text{jet}} < 35 \text{GeV}$	$p_{\rm T}^{\rm lead. \ lep.}$	36	[19]
$pp \rightarrow \ell^{\pm} \nu \ell^{+} \ell^{-}$	$m_{\ell\ell} \in (81, 101) \mathrm{GeV}$	$m_{\mathrm{T}}^{\mathrm{W}Z}$	36	[20]
$pp \rightarrow \ell^+ \ell^- \ell^+ \ell^-$	$m_{4\ell} > 180 \mathrm{GeV}$	m_{Z2}	139	[21]
$pp \to \ell^+ \ell^- jj$	$m_{jj} > 1000 \text{GeV}, m_{\ell\ell} \in (81, 101) \text{GeV}$	$\Delta \phi_{jj}$	139	[22]

[ATL-PHYS-PUB-2022-037]







ATLAS Global EFT fit

- Not enough information in the measurements to constrain all EFT coeffs
- Use eigenvector decomposition and PCA to identify most sensitive directions
- Choose physics driven groupings for most useful interpretation

ATL-PHYS-PUB-2022-037

·					0.8 -0.45 0.02 0.01 -0.23 0.3 0.01 0.02 -0.02 -0.04 -0.06 0.01 -0.06
0.02	2 -0.06 -0	0.61 -0.02			-0.01 0.66 -0.38 0.21 -0.01 0.01 -0.01 0.01 0.01
0.01	0.05 0.1	.79 0.03			-0.01 0.51 -0.29 0.16 -0.01 0.03
4] V	-0.01				0.29 -0.2 0.09 0.14 0.03 0.61 -0.55 0.35 -0.05 0.05 0.11 0.12 -0.02 0.12
i] V 0.01	-0.01 0.0	.01			0.23 0.51 0.24 0.47 0.08 0.59 0.11 0.21 -0.05 -0.01 0.01 0.03 0.02 0.03 0.01 0.01
i] V	-0.04 0.0	.01	0.01		-0.02 -0.02 -0.06 -0.09 0.01 -0.03 0.56 0.24 -0.17 -0.03 0.03 0.05 -0.54 0.09 -0.54
1 V	0.05		-0.02 0.01		-0.32 -0.24 0.17 0.32 0.04 0.01 -0.13 0.29 0.75 0.01 -0.15 0.04 -0.15 0.01
l] V	-0.28 0.0	.02	0.79 -0.33 -0.19 -0.09 -0.11 0.03 0.02		-0.05 -0.05 -0.03 -0.02 0.02 -0.04 0.14 0.21 0.03 0.03 -0.06 -0.04 0.14 -0.02 0.14
9] V	0.69 -0	0.06 0.01 -0.01 0.01	0.39 -0.17 -0.1 -0.05 -0.06 0.02 0.01	0.01	0.08 0.09 0.05 0.02 -0.05 0.07 -0.25 -0.34 0.01 -0.05 0.08 0.07 -0.23 0.04 -0.23
0]	0.66 -0	0.03 0.03	-0.05 0.04 0.02 0.01 0.01	-0.02 0.01 -0.01	-0.07 -0.09 -0.05 -0.08 0.08 -0.07 0.38 0.43 -0.04 0.07 -0.12 -0.13 0.27 -0.05 0.27
1] -0.0	2 0.01			0.95 -0.2 0.22 0.11	0.01 0.01 0.01 0.02 0.02 0.02 0.01 0.01
2]		-0.05 -0.01 -0.01	0.02 0.01 0.02 0.01	0.01	-0.01 -0.01 0.01 -0.03 0.03 0.07 -0.03 -0.02 0.94 -0.31 0.05 0.05 0.01
3] 1			0.01	0.02 -0.01 -0.01	0.01 -0.02 -0.01
4]	-0.01 -0	0.03 0.83 -0.04 0.08 0.01	0.01 -0.01 0.01 0.12 0.09 0.09 0.06 0.03 0.02 0.03 0.01 0.01		0.02 0.05 -0.03 -0.03 0.04 -0.07 -0.02 0.1 0.02 -0.11 0.17 0.43 0.05 0.05 0.05 -0.02
5] -0.0	1 0.04 -0.01 0.0	.02 -0.47 0.07	0.01 0.02 0.04 -0.01 0.03 -0.09 -0.07 -0.04 -0.03 -0.02 -0.02 -0.01 -0.01	-0.01	0.07 0.14 -0.07 -0.1 0.04 -0.12 -0.02 0.19 0.08 -0.13 0.19 0.75 0.09 0.12 0.08 -0.02
6]	-0.01 -0.02	-0.02 0.02 0.13 0.01	-0.16 -0.22 -0.21 -0.03 -0.11 0.04 -0.01 -0.01 -0.01		0.03 0.05 -0.28 0.01 0.87 -0.02 -0.06 -0.07 0.04 0.01 -0.06 -0.04 -0.04 -0.04 -0.01 -0.01
7] -0.0	1 -0.02 0.02	0.04 -0.03 -0.02	-0.08 -0.02 -0.15 0.01 -0.06 0.04 0.02 0.02 0.02 0.01 0.01	-0.01	0.25 0.51 -0.21 -0.5 -0.19 0.02 0.07 0.51 0.07 -0.03 -0.19 -0.02 0.05 -0.02 0.02 -0.01
8] -0.0	1 -0.01 -0.01	0.01 -0.01 0.04	0.39 0.51 0.56 -0.01 0.32 -0.09 0.02 0.01 0.01 0.01		0.05 0.1 -0.14 -0.08 0.3 0.02 -0.04 -0.04 0.1 0.06 -0.04 -0.07 -0.03 -0.03 -0.03 0.01 -0.01
9]		-0.07 -0.05	0.01 0.02 0.01 0.02 -0.01 -0.01 0.06 0.04 0.04 0.03 0.02 0.01 0.01		-0.01 -0.02 0.03 0.06 0.02 -0.06 -0.05 -0.05 0.34 -0.02 -0.05 0.06 0.92 0.06 0.05
0]	-0.08 0.0	.01 -0.16 -0.13 -0.13 -0.01	-0.03 -0.12 0.06 -0.02 0.03 -0.03 -0.01 0.41 0.26 0.26 0.17 0.09 0.08 0.07 0.02 0.02 0.01 0.01	-0.01	0.08 0.16 0.09 0.18 0.03 -0.34 0.01 -0.01 0.48 0.03 0.11 -0.04 -0.28 -0.04 0.1
1]	-0.03 -0	0.01 0.2 -0.07 -0.15 -0.01	0.01 0.04 -0.03 0.06 -0.04 -0.02 -0.47 -0.29 -0.29 -0.19 -0.1 -0.09 -0.07 -0.02 -0.02 -0.01 -0.01	0.01	0.07 0.15 0.08 0.16 0.02 -0.31 -0.01 -0.01 0.49 0.03 0.05 -0.03 -0.15 -0.04 0.06
2]	-0.07	-0.04 0.05 0.05	0.15 0.66 -0.48 0.33 -0.38 0.11 -0.06 0.09 0.06 0.06 0.04 0.02 0.02 0.01		-0.02 -0.03 0.01 0.04 0.04 -0.01 -0.01 -0.01 -0.03 0.05 0.02 0.05 -0.05 -0.03
3]	-0.01 0.78	-0.04 0.01 0.49 0.03	-0.01 -0.01 -0.02 -0.01 0.01		-0.03 -0.06 -0.01 -0.05 -0.05 0.13 -0.02 0.33 0.05 0.08 -0.01 -0.08 -0.01 -0.04 0.04
4]	-0.02 -0.52	-0.04 0.09 0.81 0.05	-0.01 -0.07 0.05 0.01 0.03 -0.02 -0.01 -0.03 -0.02 -0.01 -0.01 -0.01 -0.02		0.03 0.07 0.07 0.06 -0.12 -0.13 0.07 -0.01 -0.05 0.03 -0.01 -0.08 -0.03
5]	-0.34	0.01 -0.02 -0.08	-0.03 -0.01 -0.17 -0.02 0.02 0.08 -0.01 -0.01 -0.01	0.01 -0.01 -0.03 -0.01 -0.01	-0.14 -0.27 -0.15 -0.28 -0.02 0.59 -0.04 0.02 0.48 0.1 0.2 -0.12 0.02 0.03
6] /	0.01	0.07 0.91 -0.14 -0.01	-0.07 0.05 0.08 0.03 -0.03 -0.05 0.19 -0.04 -0.04 -0.01 -0.03 0.02 -0.09 -0.01 -0.01	0.01 0.02 -0.01	0.01 0.03 0.01 0.01 0.01 -0.03 -0.01 0.01 0.15 0.01 -0.02 0.01 0.01 -0.03 -0.23 -0.04
7]	-0.04	-0.01 -0.11 -0.02	0.04 -0.27 0.23 0.79 -0.07 -0.16 -0.35 -0.03 0.01 0.01 0.01 0.02	0.05 0.19 -0.03 -0.04 -0.07 -0.03 -0.03 -0.01 -0.01 -0.01	-0.02 -0.04 -0.05 -0.08 0.02 0.13 0.02 0.02 0.02 0.01 0.03 -0.02 0.06
8]	0.01	0.01	-0.01 0.06 -0.06 -0.17 0.02 0.03 0.07 0.01 0.01	0.26 0.87 -0.22 -0.18 -0.17 -0.09 -0.07 -0.02 -0.02 -0.02	0.01 0.02 0.02 0.02 -0.04 -0.01 -0.01 -0.01 -0.01



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			-0.2
	0.01	0.01	
			-0.3
	0.01		-0.4
			-0.5
0.01			0.6
-0.03	-0.01		-0.0
-0.02			-0.7
	-0.02	-0.01	-0.8
0.05	0.01		
0.1			-0.9
-0.02	0.08	0.04	-1
Chyr	C/p	Cin	



CALINE OF A STATE OF Perform eigenvector decomposition and PCA to identify operator groups Small correlation between groups



LHC EFT WG: R. Balasubramanian et al. 42









[2105.00006]



ATLAS Global EFT fit

Input measurements:

- ATLAS Higgs boson data: A combined measurement of Higgs boson production and decay in exclusive kinematic regions of the production phase space, defined within the Simplified Template Cross-Section (STXS) framework [ATLAS-CONF-2021-053]
- ATLAS electroweak data: Differential cross-section measurements for diboson production and Z boson production via vector boson fusion (VBF) [ATL-PHYS-<u>PUB-2021-0221</u>.
- Electroweak precision data (EWPD): A combined measurements of electroweak precision observables (EWPO) on the Z resonance [arXiv:0509008] that were performed at LEP and SLC.

ATL-PHYS-PUB-2022-037]

Wilson	coefficient and operator	Affected process group				
		LEP/SLD EWPO	ATLAS Higgs	A eleo		
$c_{H\Box}$	$(H^\dagger H) \Box (H^\dagger H)$		\checkmark			
c_G	$f^{abc}G^{a u}_{\mu}G^{b ho}_{ u}G^{c\mu}_{ ho}$		\checkmark			
c_W	$\epsilon^{IJK}W^{I u}_{\mu}W^{J ho}_{ u}W^{K\mu}_{ ho}$		\checkmark			
c_{HD}	$\left(H^{\dagger}D_{\mu}H ight)^{*}\left(H^{\dagger}D_{\mu}H ight)$		\checkmark			
c_{HG}	$H^\dagger H G^A_{\mu u} G^{A\mu u}$		\checkmark			
C_{HB}	$H^\dagger H B_{\mu u} B^{\mu u}$		\checkmark			
c_{HW}	$H^\dagger H W^I_{\mu u} W^{I\mu u}$		\checkmark			
c_{HWB}	$H^\dagger au^I H W^I_{\mu u} B^{\mu u}$	\checkmark	\checkmark			
c_{eH}	$(H^{\dagger}H)(\bar{l}_{p}e_{r}H)$		\checkmark			
C_{uH}	$(H^{\dagger}H)(\bar{q}Y_{u}^{\dagger}u\widetilde{H})$		\checkmark			
c_{tH}	$(H^{\dagger}H)(ar{Q}\widetilde{H}t)$		\checkmark			
c _{bH}	$(H^\dagger H)(ar{Q}Hb)$		\checkmark			
$c_{Hl}^{(1)}$	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}H)(\bar{l}\gamma^{\mu}l)$	\checkmark	\checkmark			
$c_{Hl}^{(3)}$	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}^{I}H)(\bar{l}\tau^{I}\gamma^{\mu}l)$	\checkmark	\checkmark			
c_{He}	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}H)(\bar{e}\gamma^{\mu}e)$	\checkmark	\checkmark			
$c_{Hq}^{(1)}$	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}H)(\bar{q}\gamma^{\mu}q)$	\checkmark	\checkmark			
$c_{Hq}^{(3)}$	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}^{I}H)(\bar{q}\tau^{I}\gamma^{\mu}q)$	\checkmark	\checkmark			
c_{Hu}	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}H)(\bar{u}\gamma^{\mu}u)$	\checkmark	\checkmark			
C _{Hd}	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}H)(\bar{d}\gamma^{\mu}d)$	\checkmark	\checkmark			
$c_{HQ}^{(1)}$	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}H)(\bar{Q}\gamma^{\mu}Q)$	\checkmark	\checkmark			
$c_{HO}^{(3)}$	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}^{I}H)(\bar{Q}\tau^{I}\gamma^{\mu}Q)$	\checkmark	\checkmark			
c_{Hb}	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}H)(\bar{b}\gamma^{\mu}b)$	\checkmark				
c_{Ht}	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}H)(\bar{t}\gamma^{\mu}t)$	\checkmark	\checkmark			
c_{tG}	$(\bar{Q}\sigma^{\mu\nu}T^At)\widetilde{H}G^A_{\mu\nu}$		\checkmark			
c_{tW}	$(\bar{Q}\sigma^{\mu u}t)\tau^I\widetilde{H}W^I_{\mu u}$		\checkmark			
c_{tB}	$(\bar{Q}\sigma^{\mu\nu}t)\widetilde{H}B_{\mu\nu}$		\checkmark			
c _{ll}	$(\bar{l}\gamma_{\mu}l)(\bar{l}\gamma^{\mu}l)$	\checkmark				





High pT Top+X







	$t\bar{t}W(j)$	$t\bar{t}WW$	$t\bar{t}Z(j)$	$t\bar{t}\gamma(j)$	$tar{t}\gamma\gamma$	$t\bar{t}\gamma Z$	$t\bar{t}ZZ$	VBF
$t W \to t W$	1	1						1
$t Z \to t Z$			1				1	1
$tZ \to t\gamma$			1	1		1		1
$t\gamma \to t\gamma$				1	1			1

Table 6: The set of two-top $2 \rightarrow 2$ scattering amplitudes without Higgs bosons considered in this work mapped to the collider processes in which they are embedded.

A summary of the maximal energy growths obtained in our helicity amplitude computations, taken from Tables 13–16, is shown in Table 7. A clear favourite emerges in $tW \rightarrow tW$ scattering, which displays maximal and interfering energy growth for all current operators. It has equal or better energy growth for all other operators apart from $\mathcal{O}_{\varphi B}$. In contrast, the other two amplitudes show at most linear growth in all cases barring the dipole operators, which have a tendency to grow maximally everywhere.

	$\mathcal{O}_{arphi D}$	$\mathcal{O}_{arphi^{\Box}}$	$\mathcal{O}_{arphi B}$	$\mathcal{O}_{arphi W}$	$\mathcal{O}_{arphi WB}$	\mathcal{O}_W	${\cal O}_{tarphi}$	${\cal O}_{tB}$	\mathcal{O}_{tW}	$\mathcal{O}^{(1)}_{arphi Q}$	$\mathcal{O}^{(3)}_{arphi Q}$	$\mathcal{O}_{arphi t}$
$t W \to t W$	E	E	_	E	E	E^2	E	E	E^2	E^2	E^2	E^2
$t Z \to t Z$	E	E	E	E	E	_	E	E^2	E^2	E	E	E
$t Z \to t \gamma$	_	_		E	E	_	_	E^2	E^2	_	_	_
$t\gamma \to t\gamma$	_	_	E	E	E	_	_	E	E	_	_	_

[1904.05637]



High pT Top+X



Figure 42: Same as Figure 40 for $\mathcal{O}_{\varphi t}$

[1904.05637]



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High pT Top+X

	$\mathcal{O}_{arphi D}$	$\mathcal{O}_{arphi^{\Box}}$	${\cal O}_{arphi B}$	$\mathcal{O}_{arphi W}$	$\mathcal{O}_{arphi WB}$	\mathcal{O}_W	${\cal O}_{tarphi}$	${\cal O}_{tB}$	\mathcal{O}_{tW}	$\mathcal{O}^{(1)}_{arphi Q}$	${\cal O}^{(3)}_{arphi Q}$	$\mathcal{O}_{arphi t}$	$\mathcal{O}_{arphi tb}$
$b W \rightarrow t Z$		_	_	_	E	E^2	_	E^2	E^2	E	E^2	E	E^2
$b W \rightarrow t \gamma$	-	_	_	_	E	E^2	_	E^2	E^2	_	_	_	_
$b W \rightarrow t h$	_	_	_	E	_	_		_	E^2	_	E^2	_	E^2

able 5: Maximal energy growths induced by each operator on the set of single top attering amplitudes considered. '-' denotes either no contribution or no energy growth nd the red entries denote the fact that the interference between the SMEFT and the SM nplitudes also grows with energy.

	tWj	tZj	$t\gamma j$	tWZ	$tW\gamma$	thj	thW
$b W \to t Z$	1	1		1			
$b W ightarrow t \gamma$	1		1		1		
$b W \rightarrow t h$						√	1







[1904.05637]





EFT | Extending to differential Higgs

Simplified Template Cross section (STXS)

- Going beyond kappa framework
- Compromise for channels without stats for fiducial crosssection.

Key features:

- Regions defined inclusively in Higgs decay & kinematics
- Specific to Higgs prod. mode, topology and kinematics

Enables:

- Combination of multiple Higgs decays to extract STXS bins in prod. mode/event topology
- Permits use of MVA techniques.



EFT | In practise... Higgs ATLAS Higgs combination











EFT Higgs: USE Higgs: ATLAS Higgs combination

- EFT parameter sensitivity
- Same parameter decomposition as before





EFT | Higgs

- ATLAS Higgs combination
 - **EFT** parameter sensitivity
 - Same parameter decomposition as before

Results

- All parameters consistent with SM expectation
- Somewhat tighter tighter constraints from lin+quad fit
 - Implies non-negligible influence of these terms.





I Higgs EFT STXS also lends itself to wider interpretation



[arXiv:2012.02779]





EFT | Higgs STXS also lends itself to wider interpretation

Significant impact of STXS measurements in global fits

More on this later...







[arXiv:2012.02779]



EFT Flavour Assumptions

Starting-point proposal

more **Top-philic** (extending the 'universal' scenario) restrictive new physics couples dominantly to bosons+tops e.g. realized in composite Higgs scenarios • $\mathcal{O}(30)$ CP-even + $\mathcal{O}(10)$ CP-odd d.o.f. not radiatively stable baseline $SU(2)_{u,q}^2 \times SU(3)_{d,l,e}^3$ · basically MFV with all breakings neglected apart from y_t • $\mathcal{O}(100)$ d.o.f. $[\mathcal{O}(180) \text{ for } SU(2)^3_{a,u,d} \times SU(3)^2_{l,e}]$ · reasonable approx. for dim \leq 4 too radiatively stable then · massless b (5F scheme), no $h \rightarrow b \bar{b}$ or $\mu^+ \mu^$ less LFU-violating restrictive • separating e, μ, τ (e.g. $U(1)_{e}^{3} \times U(1)_{I}^{3}$, $[U(1)_{I+e}]^{3}$, $U(2)^{5}$, etc.) \cdot not needed, a priori, given the limited interplay with B anomalies (only in high-mass Drell-Yan) • $\mathcal{O}(15)$ Warsaw operators with leptons

[LHC EFT WG Meeting]



EFT | Top+X





- SMEFT describes ~all BSM models characterised by a scale Λ ($\gg v(E)$)
- Additions to the SM built from dim>4 operators based on SM fields
- Schematically the Lagrangian is:

$$\mathscr{L}_{SMEFT} = \mathscr{L}_{SM} + \mathscr{L}_{SM}$$

$y_{5} + \mathcal{L}_{6} + \mathcal{L}_{7} + \mathcal{L}_{8} + \dots$



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$$\mathscr{L}_{SMEFT} = \mathscr{L}_{SM} + \mathscr{L}_{SM}$$

$$\int \mathsf{SM} (\dim 4)$$

$z_{5} + Z_{6} + Z_{7} + Z_{8} + \dots$



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1 lepton n

umber violating parameter

$'_{5} + \mathscr{L}_{6} + \mathscr{L}_{7} + \mathscr{L}_{8} + \dots$



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$$\mathscr{L}_{SMEFT} = \mathscr{L}_{SM} + \mathscr{L}_5 + \mathscr{L}_6 + \mathscr{L}_7 + \mathscr{L}_8 + \dots$$

~100-1000 parameters depending on flavour structure



- SMEFT describes ~all BSM models characterised by a scale Λ ($\gg v(E)$)
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- Schematically the Lagrangian is:

$$\begin{aligned} \mathscr{L}_{SMEFT} &= \mathscr{L}_{SM} + \mathscr{L}_5 + \mathscr{L}_6 + \mathscr{L}_7 + \mathscr{L}_8 + \dots \\ & & & & & \\ & & & & \\ & & & & 30 \text{ operators} \\ & & & & & \text{violating lepton/} \\ & & & & & & \\ & & & & & & \text{baryon number} \end{aligned}$$



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- Additions to the SM built from dim>4 operators based on SM fields
- Schematically the Lagrangian is:

$$\mathscr{L}_{SMEFT} = \mathscr{L}_{SM} + \mathscr{L}_{SM}$$

 ${}^{2}_{5} + \mathscr{L}_{6} + \mathscr{L}_{7} + (\mathscr{L}_{8}) + \dots$ ~1000 operators for N_f=1 scenario



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$$\mathscr{L}_{SMEFT} = \mathscr{L}_{SM} + \mathscr{L}_{SM}$$

$y_{5} + \mathcal{L}_{6} + \mathcal{L}_{7} + \mathcal{L}_{8} + \dots$



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- Additions to the SM built from dim>4 operators based on SM fields
- Schematically the Lagrangian is:

$$\mathscr{L}_{SMEFT} = \mathscr{L}_{SM} + \frac{\mathscr{L}}{\Lambda}$$





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In practise (for simplicity) we often only consider the dimension 6 operators





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- In practise (for simplicity) we often only consider the dimension 6 operators
- In terms of operators and coefficients

$$\frac{\mathscr{L}_6}{\Lambda^2} = \sum_{i}^{I}$$







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- Additions to the SM built from dim>4 operators based on SM fields
- Schematically the Lagrangian is:

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- In practise (for simplicity) we often only consider the dimension 6 operators
- In terms of operators and coefficients
 - Scale of new physics
 - Wilson coefficients
 - Dimension 6 operators







the corresponding observables are:

$$\mathscr{A} = \mathscr{A}_{SM} + \mathscr{A}_6$$





the corresponding observables are:

$$\mathcal{A} = \mathcal{A}_{SM} + \mathcal{A}_6 \quad \rightarrow$$

 $\frac{\mathscr{L}_6}{\Lambda^2} = \sum_{i} \frac{C_i \mathscr{O}_i^{a=\upsilon}}{\Lambda^2}$

$|\mathscr{A}_{SM}|^2 + 2Re\mathscr{A}_{SM}\mathscr{A}_6^{\dagger} + |\mathscr{A}_6|^2$



the corresponding observables are:

$$\mathcal{A} = \mathcal{A}_{SM} + \mathcal{A}_6 \quad \rightarrow$$

 $\frac{\mathscr{L}_6}{\Lambda^2} = \sum_{i=1}^{\infty} \frac{C_i \mathcal{O}_i^{d=0}}{\Lambda^2}$

 $|\mathscr{A}_{SM}|^2 + 2Re\mathscr{A}_{SM}\mathscr{A}_6^{\dagger} + |\mathscr{A}_6|^2$

"Interference terms"



the corresponding observables are:

$$\mathcal{A} = \mathcal{A}_{SM} + \mathcal{A}_6 \quad \rightarrow$$

 $\left|\mathscr{A}_{SM}\right|^{2} + 2Re\mathscr{A}_{SM}\mathscr{A}_{6}^{\dagger} + \left|\mathscr{A}_{6}\right|^{2}$ "Interference terms" "Quadratic terms"




EFT | Basics Given these Lagrangian terms

the corresponding observables are:

$$\mathscr{A} = \mathscr{A}_{SM} + \mathscr{A}_6 \quad \rightarrow$$

The C_i are what we try to constrain in measurements.

 $\left|\mathscr{A}_{SM}\right|^{2} + 2Re\mathscr{A}_{SM}\mathscr{A}_{6}^{\dagger} + \left|\mathscr{A}_{6}\right|^{2}$ "Interference terms" "Quadratic terms"

Josh McFayden | Top2022 | 6/9/2022





EFT | Basics Given these Lagrangian terms

the corresponding observables are:

$$\mathscr{A} = \mathscr{A}_{SM} + \mathscr{A}_6 \quad \rightarrow$$

The C_i are what we try to constrain in measurements.

• (For a given Λ , conventionally set to 1 TeV)

 $\left|\mathscr{A}_{SM}\right|^{2} + 2Re\mathscr{A}_{SM}\mathscr{A}_{6}^{\dagger} + \left|\mathscr{A}_{6}\right|^{2}$ "Interference terms" "Quadratic terms"

Josh McFayden | Top2022 | 6/9/2022



