

Soft Anomalous Dimension in QCD

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Outline

- Motivation for studying the Soft Anomalous Dimension
- Introduction: Set-up and Definition
- Wilson Lines and Lie Algebra
- Color and Kinematics in diagrams



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- Only fully connected diagrams present (Gardi, Smilie, White 2013)
- Bootstrapped at 3 loops for massless partons (Almelid, Duhr, Gardi Mcleod, White 2017)
- High energy (Regge) limit of the SAD=soft limit of the BFKL resummation (Del Duca, Duhr, Gardi, Magnea, White 2011)
- Important for resummation of large logs in n -jet process to N^3LL accuracy, possible check for higgs production calculations

Current Status of Soft Anomalous Dimension

Massless Partons

- Ansatz at four loops in QCD (Becher, Neubert 2019)
- Exact calculation and bootstrapping at three loops (Almelid, Duhr Gardi 2016, + McLeod, White 2017))

Other

- Production of single top quark at three loops (Kidonakis 2019)
- Massive particle 2 loops (Mitov 2009)
- Exact Soft Anomalous Dimension in Super Yang-Mills (Giombi, Komatsu 2020)

Introduction: Set-up 1/2

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$$\frac{i(\not{p}+\not{k})}{(p+k)^2+i0} \mathcal{M}$$

$$2p \cdot k = 2p_0 k_0 (1 - \cos \theta_{pk})$$

Introduction: Set-up 1/2

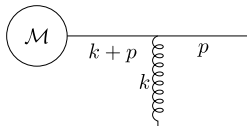
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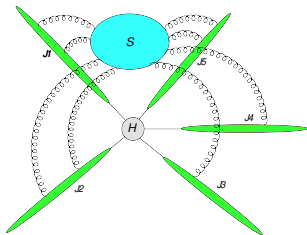
Soft: $k_0 = 0$

Collinear: $\cos \theta_{pk} = 1$



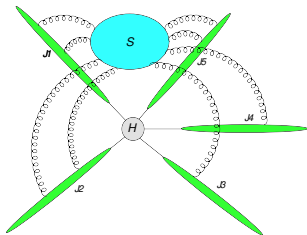
Introduction: Set-up 2/2

- Hard/Soft/Jet: Amplitude factorisation

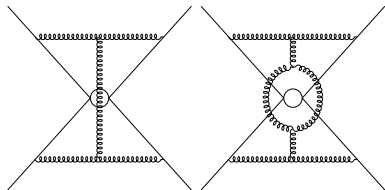


Introduction: Set-up 2/2

- Hard/Soft/Jet: Amplitude factorisation



- Loops



Soft Anomalous dimension(SAD) Definition

Amplitude Factorises with μ_f as the scale

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Renormalisation group equation for the scale μ_f

$$\frac{d}{d \ln \mu_f} Z_n(\{p_i\}, \{\mathbf{T}_i\}, \epsilon, \alpha_s(\mu_f)) = -Z_n \Gamma_n(\{p_i\}, \{\mathbf{T}_i\}, \mu_f, \alpha_s(\mu_f)) \quad (2)$$

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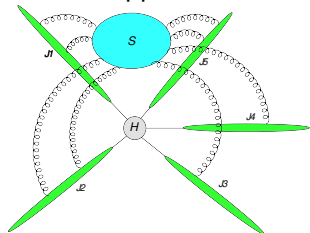
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$$Z_n = \mathcal{P} \exp \left\{ -\frac{1}{2} \int_0^{\mu_f^2} \frac{d\lambda^2}{\lambda^2} \Gamma_n(\{p_i\}, \{\mathbf{T}_i\}, \lambda, \alpha_s(\lambda^2)) \right\}, \quad (3)$$

Wilson Lines

Non-abelian exponentiation: Gardi, Smilie, White 1304.7040

Eikonal Approximation



$$\frac{i(\not{p} + \not{k})}{(p + k)^2 + i0} \mathcal{M} \xrightarrow{k \ll p} \frac{p^\mu}{p \cdot k + i0} \mathcal{M} = \frac{\beta^\mu}{\beta \cdot k + i0} \quad (4)$$

Wilson Line representation

$$\Phi_\beta(0, \infty) = \mathcal{P} \left(\exp \left[i \int_C dx^\mu \beta_i A_\mu^a(\beta_i x) T^a \right] \right) \quad (5)$$

$SU(N_c)$ Lie Algebra 1/2

1. Commutator

$$\mathbf{T}_1^a \mathbf{T}_1^b - \mathbf{T}_1^b \mathbf{T}_1^a = [\mathbf{T}_1^a, \mathbf{T}_1^b] = i f^{abc} \mathbf{T}_1^c \quad (7)$$

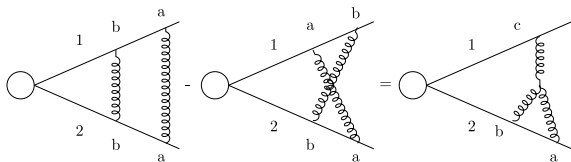
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$SU(N_c)$ Lie Algebra 2/2

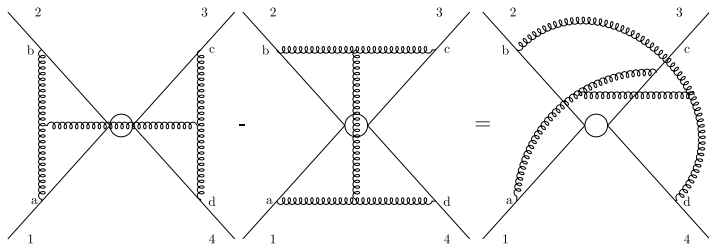
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Kinematics and Weight

- weight: a Multiple PolyLogarithm of weight n is

$$G(a_1, a_2, \dots, a_n; z) = \int_0^z \frac{dt}{t - a_1} G(a_2, \dots, a_n; t), \quad (11)$$

where G is a MPL and a_i are indices that make up a length n vector.

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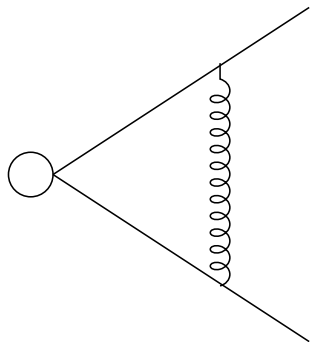
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- For 4 loops, Γ_n contains weight 7 functions.

Dipole



$$\mathbf{T}_i \cdot \mathbf{T}_j = (\mathbf{T}_i^a)_{km} (\mathbf{T}_j^a)_{pq}$$

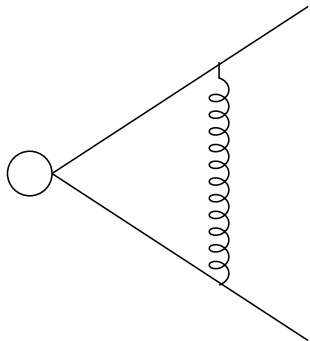
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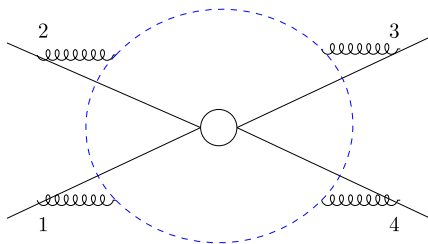
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$$\log\left(\frac{-s_{ij}}{\mu^2}\right)$$

New at four loops

- $\frac{1}{6} \sum_{\sigma} \text{Tr}(\mathbf{T}_R^a \mathbf{T}_R^{\sigma(b)} \mathbf{T}_R^{\sigma(c)} \mathbf{T}_R^{\sigma(d)}) \mathbf{T}_1^a \mathbf{T}_2^b \mathbf{T}_3^c \mathbf{T}_4^d$



- Kinematics: Completely symmetric
- Two independent cross ratios
- $G^R(\beta_{1243}, \beta_{1342})$

Conclusion and outlook

- Provided motivation for studying the soft anomalous dimension
- Looked at diagrammatic representation of the color structures
- Commented on some symmetries and the kinematics
- Find the form of the kinematic functions at four loops using the limits

References

- 0901.1091
- 1008.0098
- 1304.7040
- 1908.11379
- 1911.10174
- <https://wmresources.org/august-2019-the-long-and-winding-road/>

Kinematic Functions: Single-Valued Harmonic Polylogarithms

- Harmonic Polylogarithm

$$H(a_1, a_2, \dots, a_n; z) = \int_0^z \frac{dt}{t - a_1} H(a_2, \dots, a_n; t), \quad (12)$$

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- Single-valued means there are no branch cuts e.g

$$\begin{aligned} \mathcal{L}_0 &= \ln(z) + \ln(\bar{z}) = H(0, z) + H(0, \bar{z}) = G(0, z) + G(0, \bar{z}) \\ \mathcal{L}_1 &= \ln(1 - z) + \ln(1 - \bar{z}) = H(1, z) + H(1, \bar{z}) \\ &= -(G(1, z) + G(1, \bar{z})) \end{aligned} \quad (13)$$

- The kinematic functions of Γ_s are made up weight 5 SVHPLs

Casimirs

Quadratic Casimir: $C_2(i)\mathbb{I} = (\mathbf{T}_i^a)_{km}(\mathbf{T}_i^a)_{mq}$

