

Young Theorists Forum 20 16th December 2020

NEVER HAVE I EVER used a constant to fix my problem

Bruno Valeixo Bento

Partly based on work in collaboration with Susha Parameswaran, Ivonne Zavala, Dibya Chakraborty

The Cosmological Constant

Albert
Hey guys! Here's gravity!

$$G_{\mu\nu} = 8\pi G T_{\mu\nu}$$

As you see, the war treated me kindly enough, in spite of the heavy gunfire, to
allow me to get away from it all and take this walk in the land of your ideas.

$$ds^2 = -\left(1 - \frac{2GM}{r}\right)dt^2 + \left(1 - \frac{2GM}{r}\right)^{-1}dr^2 + r^2(d\theta^2 + \sin^2\theta \, d\phi^2)$$
Albert
Ups...There was a mistake!
 $G_{\mu\nu} + \Lambda g_{\mu\nu} = 8\pi G T_{\mu\nu}$
That term is necessary only for the purpose of making
possible a quasi-static distribution of matter, as required by
the fact of the small velocities of the stars.



[Chapter 1]

Karl

As you see, the war treated me kindly enough, in spite of the heavy gunfire, to allow me to get away from it all and take this walk in the land of your ideas.



$$ds^{2} = -\left(1 - \frac{2GM}{r}\right)dt^{2} + \left(1 - \frac{2GM}{r}\right)^{-1} + r^{2}(d\theta^{2} + \sin^{2}\theta \, d\phi^{2})$$

Albert

Ups...There was a mistake! No static Universe otherwise...

$$G_{\mu\nu} + \Lambda g_{\mu\nu} = 8\pi G T_{\mu\nu}$$

Alexander

I found nice expanding solutions with the first equation...



 $G_{\mu\nu} = 8\pi G T_{\mu\nu}$

The Cosmological teamstant gh, in spite of Chapter 1

$$ds^{2} = -\left(1 - \frac{2GM}{r}\right)dt^{2} + \left(1 - \frac{2GM}{r}\right)^{-1} + r^{2}(d\theta^{2} + \sin^{2}\theta \, d\phi^{2})$$

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Karl

I found nice expanding solutions with the first equation...



Georges

Totally agree with this! ^



The Cosmological Constant

[Chapter 1]

FLRW background curvature

$$ds^{2} = -dt^{2} + a^{2}(t) \left(\frac{dr^{2}}{1 - kr^{2}} + r^{2}(d\theta^{2} + \sin^{2}\theta \, d\varphi^{2}) \right)$$

Einstein Equations ⇒ Friedmann Equations

acceleration requires

$$\omega_{eff} = \frac{P_{tot}}{\rho_{tot}} < -\frac{1}{3}$$

$$H^{2} = \left(\frac{\dot{a}}{a}\right)^{2} = \frac{\rho_{tot}}{3M_{Pl}^{2}} - \frac{k}{a^{2}} \qquad \frac{\ddot{a}}{a} = -\frac{1}{6M_{Pl}^{2}}(\rho_{tot} + 3P_{tot})$$





The Cosmological Constant [Chapter 1]



Georges

Totally agree with this! ^



Edwin

Sorry bro... Those galaxies are definitely running away! Red as a tomato...



$$R_{\mu\nu} - \frac{1}{2} R g_{\mu\nu} = T_{\mu\nu}$$

The Cosmological Constant



The Fake Return or the subtleties of Inflation



The Cosmological Constant The Fake Return or the subtleties of Inflation

A <u>temporary</u> cosmological constant in the early universe would help with...

- 1. Horizon problem
- 2. Flatness problem
- 3. Magnetic monopoles problem

⇒ INFLATION

<u>Period</u> of exponential expansion in the early universe

[Chapter 2]

Any vacuum energy will look like a cosmological constant!



The Cosmological Constant The Fake Return or the subtleties of Inflation

[Chapter 2]

Inflation has its own <u>specific needs</u>:

- Enough expansion ($N \sim 60$)
- To reproduce power spectrum, e.g.
 - $n_s = 0.9626 \pm 0.0057$
 - *r* < 0.056
- Not the Λ

- To end -
- with reheating



The anisotropies of the CMB as observed by Planck. Image by ESA and Planck collaboration

> Planck collaboration [arxiv:1807.06211]



[1998] Houston, we have a problem... The Universe is accelerating! [Saul Perlmutter, Brian P. Schmidt, Adam G. Riess]

- 1. Bring back Λ (Λ CDM)
- 2. Have a positive energy vacuum (de Sitter)
- 3. Have a slowly rolling scalar field (Quintessence)
- 4. ???

[Chapter 3]

[1998] Houston, we have a problem... The Universe is accelerating! [Saul Perlmutter, Brian P. Schmidt, Adam G. Riess]



Antonio Padilla [arxiv:1502.05296]



[1998] Houston, we have a problem... The Universe is accelerating! [Saul Perlmutter, Brian P. Schmidt, Adam G. Riess]

Use
$$\Lambda$$
 to cance
 $R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} + (\Lambda^{\text{vac}} - \Lambda)g_{\mu\nu} = \tilde{T}_{\mu\nu}$

The Cosmological Constant Problem

The problem is **not** that the constant needs to be extremely small, it's that it needs to *stay extremely small at all orders (radiative instability)*

Antonio Padilla [arxiv:1502.05296]

A tempting guess is to use (un)broken symmetries...

What about UV complete theories?

Really hard to get de Sitter in supergravity and String Theory

- SUSY breaking scale is at least 1TeV;
- Not compatible with SUSY;
- KKLT and LVS

$$V = e^{K} \left[G^{I\bar{J}}(D_{I}W)(D_{\bar{J}}\overline{W}) - 3|W|^{2} \right]$$

with S. Parameswaran, I. Zavala, D. Chakraborty



[Chapter 3]



De Sitter Space and the Swampland

Georges Obied, Hirosi Ooguri, Lev Spodyneiko, Cumrun Vafa

[arxiv:1806.08362]

Distance and de Sitter Conjectures on the Swampland

Hirosi Ooguri, Eran Palti, Gary Shiu, Cumrun Vafa

[arxiv:1810.05506]

Refined de Sitter Conjecture: The scalar potential $V(\phi)$ for scalar fields in a low energy EFT of any consistent Quantum Gravity must satisfy either:

$$\frac{|\nabla V|}{V} \ge \frac{c}{M_p^2} \qquad \text{or} \qquad \frac{\min(\nabla_i \nabla_j V)}{V} \le -\frac{c'}{M_p^2} \qquad (1)$$

for some universal constants $c, c' \sim \mathcal{O}(1)$, where $\min(\nabla_i \nabla_j V)$ stands for the minimum eigenvalue of the Hessian matrix in an orthonormal frame.

This forbids de Sitter vacua.



[Chapter 4]

- Plenty of scalar fields in string compactifications;
- Vast landscape might help find flat enough potentials;
- Could be a maximum (hilltop quintessence).



The Cosmological Constant Endgame or the Rise of Quintessence

Dark Energy in String Theory

Bruno Valeixo Bento, Dibya Chakraborty, Susha L. Parameswaran, Ivonne Zavala

[arxiv:2005.10168]

$$K = -n \log \left(\Phi + \overline{\Phi} \right)$$
 Bulk/fibre moduli $W = W_0 + A e^{-a\Phi}$

$$K = k_0 + \frac{(\Phi + \bar{\Phi})^{2n}}{k_1}$$

Blow-up moduli

 $W = W_0 + A\Phi^p$

[Chapter 4]

 $K = k_0 + \frac{|\Phi|^{2n}}{k_1}$

Deformation moduli

Generalizing the analysis with all the possible combinations, what do we get?

The Cosmological Constant Endgame or the Rise of Quintessence

de Sitter
$$V > 0$$

Slow-roll $\epsilon_V = \frac{M_{Pl}^2}{2} g^{\phi\phi} \left(\frac{V'}{V}\right)^2 < 1$

Only for a local modulus, with a perturbative leading W

$$K = k_0 + \frac{|\Phi|^{2n}}{k_1} \qquad K = k_0 + \frac{(\Phi + \bar{\Phi})^{2n}}{k_1}$$
$$W = W_0 + A\Phi^p \qquad \qquad p = n$$

			$K = k_0 + \frac{ \Phi ^{2n}}{k_1}$,	$W = W_0 + A\Phi^p$					
$V = \frac{W_0^2}{M_{-x}^2} \frac{e^{k_0 + y}}{n^2 y} ((px + n(1 + x)y)^2 - 3n^2(1 + x)^2 y)$									
$\mathcal{E}_{V} = \frac{(p^{3}x^{2} + 3n^{2}px(1+x)(y-1)y + n^{3}(1+x)^{2}(y-2)y^{2} + np^{2}x(y+x(3y-1)))^{2}}{n^{2}y(p^{2}x^{2} + 2npx(1+x)y + n^{2}(1+x)^{2}(y-3)y)^{2}}$									
Parameters			$V \rightarrow$	$\epsilon_V \rightarrow$	$V > 0$ $\varepsilon_V < 1$				
$p \neq n$	$x \gg 1$		$\frac{W_0^2}{M_{pl}^2} e^{k_0} \frac{p^2}{n^2} \frac{x^2}{y} > 0$	$\frac{(p-n)^2}{n^2}\frac{1}{y} > 1$	No-go				
	<i>x</i> ≪ 1	$x \gg y$	$\frac{W_0^2}{M_{pl}^2} \frac{p^2 e^{k_0}}{n^2} \left(\frac{x^2}{y} - 3\frac{n^2}{p^2}\right)$	$\frac{(n-p)^2}{n^2 \left(\frac{x^2}{y} - 3\frac{n^2}{p^2}\right)^2} \frac{x^2}{y} \left(\frac{x}{y}\right)^2 > 1$	No-go				
		$x \ll y$	$-\frac{3W_0^2 e^{k_0}}{M_{pl}^2} < 0$	$\frac{4y}{9} < 1$	No-go				
Parameters			$V \rightarrow$	$\varepsilon_V ightarrow$	$V > 0$ $\varepsilon_V < 1$				
	$x \gg 1$		$rac{W_0^2}{M_{pl}^2}e^{k_0}rac{x^2}{y}>0$	$(1+xy)^2 \frac{y}{x^2} < 1$	Yes				
p = n	<i>x</i> ≪ 1	$x^2 \gg y$	$rac{W_0^2}{M_{pl}^2}e^{k_0}rac{x^2}{y}>0$	$\frac{4y}{x^2} < 1$	Yes				
		$x^2 \ll y$	$-\frac{3W_0^2 e^{k_0}}{M_{pl}^2} < 0$	$\frac{4y}{9}\left(1+\frac{x}{y}\right)^2 < 1$	No-go				

			$K = k_0 + \frac{(\Phi + \bar{\Phi})^{2n}}{k_1}$,	$W = W_0 + A\Phi^p$					
$V = \frac{W_0^2}{M_{\pi^2}^2} \frac{e^{k_0 + y}}{n(2n-1)y} (2p^2x^2 + 4npx(1+x)y + n(1+x)^2(3+2n(y-3))y)$									
$\mathcal{E}_{V} = \frac{2(2p^{3}x^{2} + 2n^{3}(1+x)^{2}(y-2)y^{2} + 3n^{2}(1+x)y(2px(y-1) + (1+x)y) + npx(3(1+x)y + 2p(y+x(3y-1))))^{2}}{n(2n-1)y(2p^{2}x^{2} + 4npx(1+x)y + n(1+x)^{2}(3+2n(y-3))y)^{2}}$									
Parameters			$V \rightarrow$	$\epsilon_V ightarrow$	$V > 0$ $\varepsilon_V < 1$				
$p \neq n$	$x \gg 1$		$\frac{\frac{W_0^2}{M_{pl}^2}}{\frac{2p^2e^{k_0}}{n(2n-1)}}\frac{x^2}{y} > 0$	$\frac{2(n-p)^2}{n(2n-1)}\frac{1}{y} > 1$	No-go				
	<i>x</i> ≪ 1	$x \gg y$	$\frac{W_0^2}{M_{pl}^2} \frac{2p^2 e^{k_0}}{n(2n-1)} \left(\frac{x^2}{y} - \frac{3n(2n-1)}{2p^2}\right)$	$\frac{2(n-p)^2}{n(2n-1)\left(\frac{x^2}{y} - \frac{3n(2n-1)}{2p^2}\right)^2} \frac{x^2}{y} \left(\frac{x}{y}\right)^2 > 1$	No-go				
		$x \ll y$	$-rac{3W_0^2 e^{k_0}}{M_{pl}^2} < 0$	$\frac{2n(4n-3)^2}{9(2n-1)^3}y < 1$	No-go				
Parameters			$V \rightarrow$	$arepsilon_V ightarrow$	$V > 0$ $\varepsilon_V < 1$				
	$x \gg 1$		$rac{W_0^2}{M_{ hol}^2}e^{k_0}rac{2n}{2n-1}rac{x^2}{y}>0$	$\frac{9y}{2n(2n-1)} < 1$	Yes				
p = n	<i>x</i> ≪ 1	$x^2 \gg y$	$\frac{\frac{W_0^2}{M_{pl}^2}}{M_{pl}^2}e^{k_0}\frac{2n}{2n-1}\frac{x^2}{y} > 0$	$\frac{(4n-3)^2}{2n(2n-1)}\frac{y}{x^2} < 1$	Yes				
		$x^2 \ll y$	$-\frac{3W_0^2 e^{k_0}}{M_{pl}^2} < 0$	$\frac{2n}{9}\frac{(4n-3)^2}{(2n-1)^3}\left(1+\frac{x}{y}\right)^2 y$	No-go				





with S. Parameswaran, I. Zavala, D. Chakraborty

To be continued...

[Chapter 4]

What's the take away?

O The cosmological constant problem is **not its size**

- O [Inflation] Scalar fields are natural solutions for accelerated expansion
- O It's hard to get de Sitter in String Theory
- O Quintessence is a very natural alternative

