

LFU: model dependence of the experimental measurements

Wolfgang Altmannshofer, Paula Alvarez Cartelle

Beyond the Flavour Anomalies
April 2021

Measuring LFU ratios

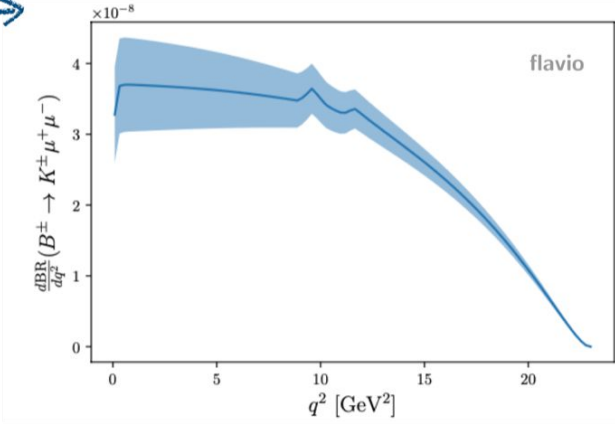
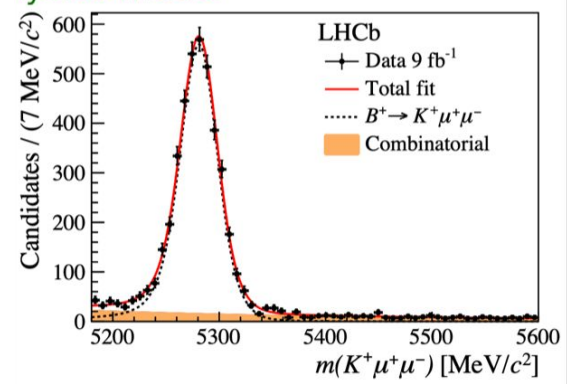
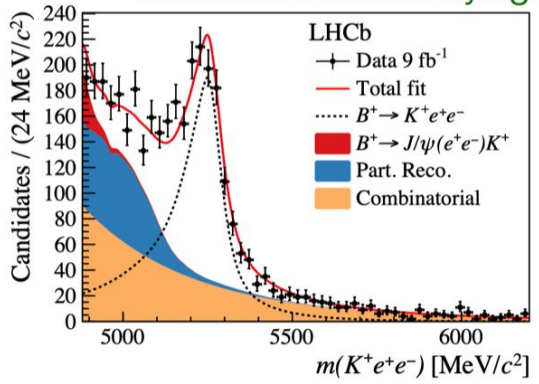
$$R_H = \left(\frac{1}{r_{J/\psi}} \right) \cdot \frac{N(B \rightarrow H^+ e^+ e^-)}{N(B \rightarrow H^+ \mu^+ \mu^-)} \cdot \frac{\epsilon(B \rightarrow H e^+ e^-)}{\epsilon(B \rightarrow H \mu^+ \mu^-)}$$

=1

Invariant mass shapes influenced by radiative corrections (QED) but otherwise not very correlated to underlying physics model

Efficiency calculation assumes SM differential decay rate (SM uncertainty included as syst)

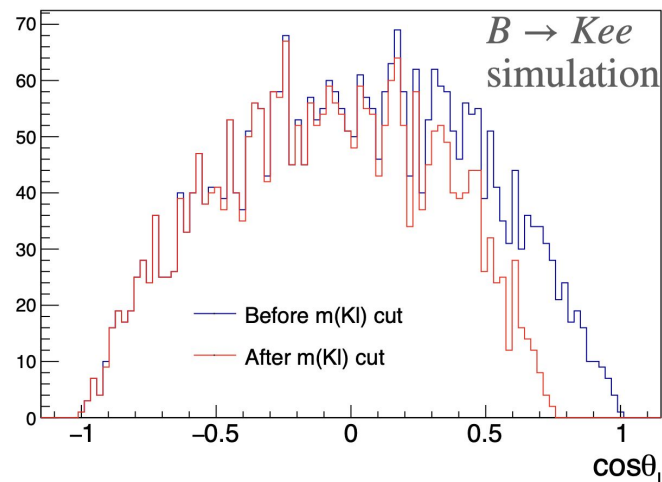
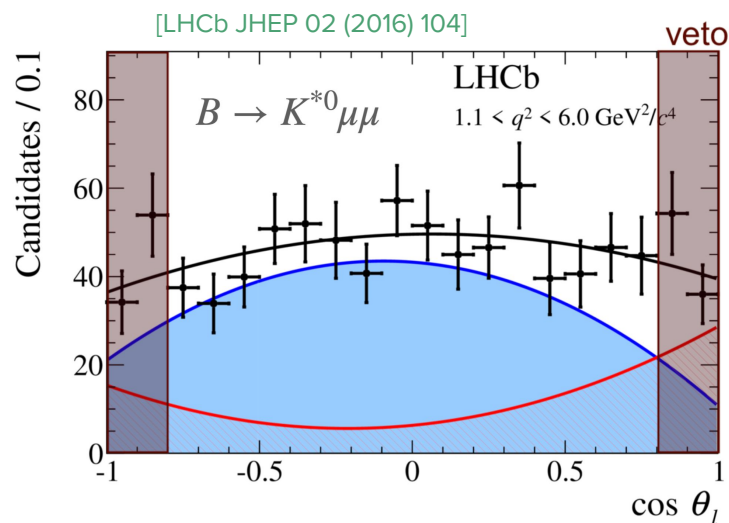
- q^2 spectrum
- angular distribution



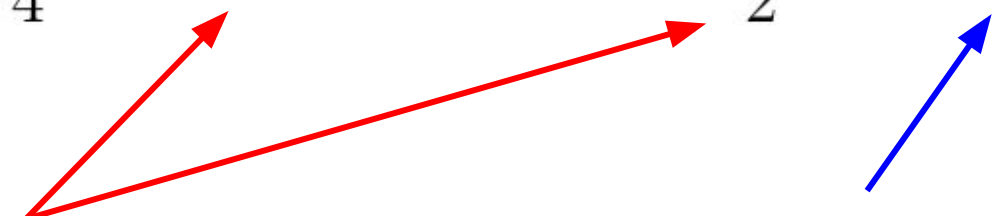
- For a null test of the SM, it does not matter
- But could this bias interpretation in terms of certain NP operators?

Signal selection

- The efficiency of the signal selection is not uniform
 - q^2 selection [1.1, 6] GeV^2/c^4
 - Cascade veto: To suppress background such as $B \rightarrow D^0 (\rightarrow K + \ell^- \nu) \ell^+ \nu$
 - $R(K)$: $m(K + \ell^-) > m(D^0)$
 - $R(K^*)$: $|\cos\theta_l| < 0.8$
 - More complicated effects induced by trigger, PID, ...
- How much/in what direction would R_K/R_{K^*} shift in a NP scenario?



The $B \rightarrow K\pi\pi$ Angular Distribution

$$\frac{1}{\Gamma} \frac{d\Gamma}{d \cos \theta_l} = \frac{3}{4}(1 - F_H)(1 - \cos^2 \theta_l) + \frac{1}{2}F_H + A_{FB} \cos \theta_l$$


“Flat term”

Forward-backward asymmetry

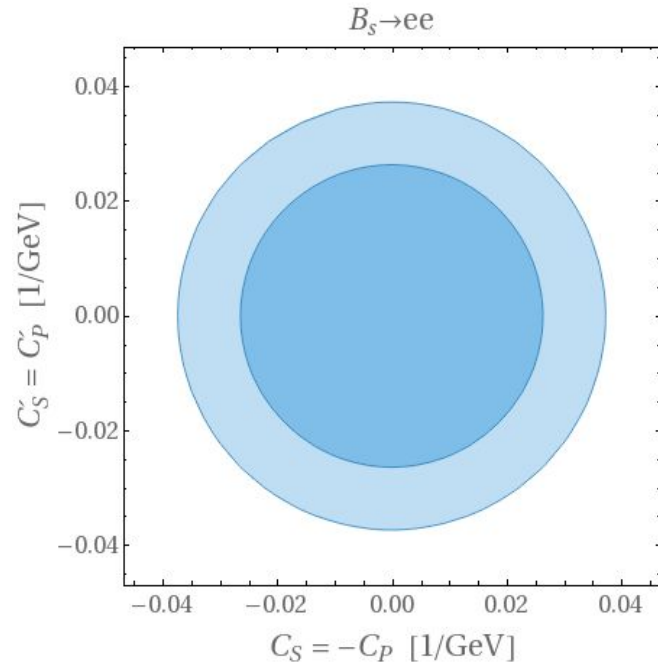
- $F_H \approx 0$ in the Standard Model
- Sensitive to (pseudo)scalar and tensor operators
- $A_{FB} \approx 0$ in the Standard Model
- Could become sizeable if both (pseudo)scalar and tensor operators are present at the same time

Scalar Operators in $B_s \rightarrow \ell\ell$ and $B \rightarrow K\ell\ell$

$$O_S^{(\ell)} = m_b(\bar{s}_{L(R)}b_{R(L)})(\bar{\mu}\mu), \quad O_P^{(\ell)} = m_b(\bar{s}_{L(R)}b_{R(L)})(\bar{\mu}\gamma_5\mu),$$

- Scalar operators are strongly constrained by $\text{BR}(B_s \rightarrow \ell\ell)$, if SMEFT relations are imposed
[\[Alonso et al. 1407.7044\]](#)

$$C_S = -C_P \quad C'_S = C'_P$$



Scalar Operators in $B_s \rightarrow \ell\ell$ and $B \rightarrow K\ell\ell$

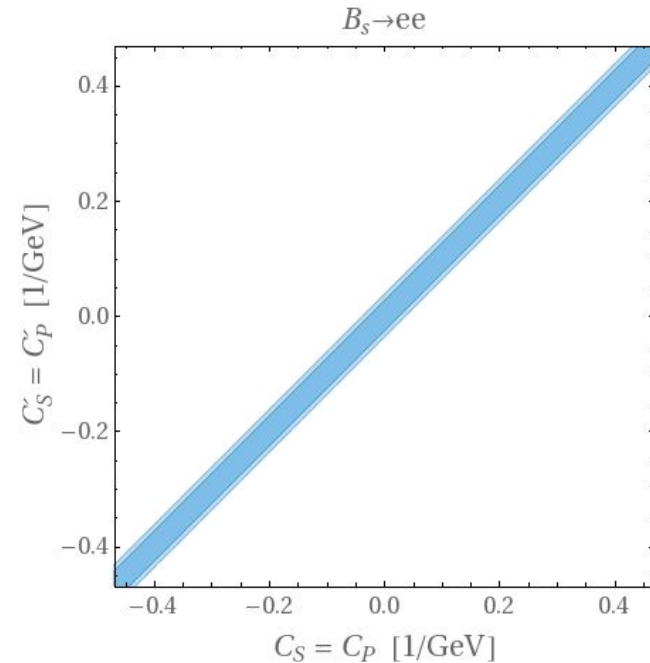
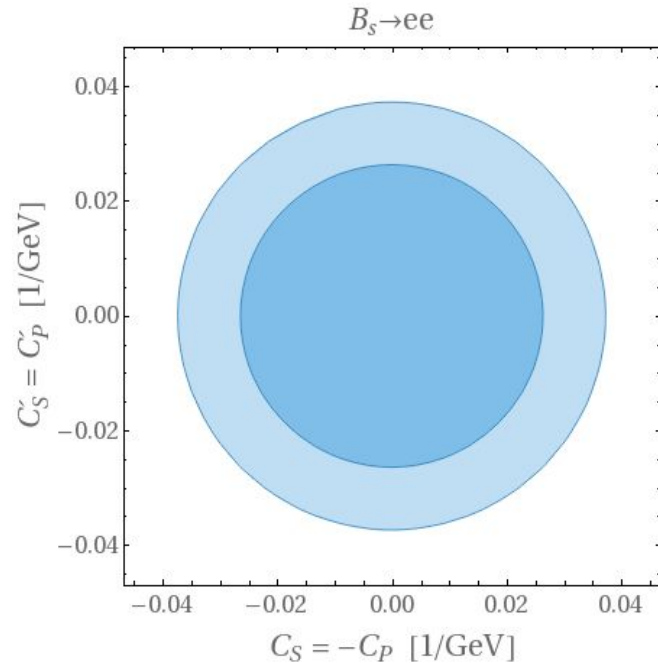
$$O_S^{(\ell)} = m_b(\bar{s}_{L(R)}b_{R(L)})(\bar{\mu}\mu), \quad O_P^{(\ell)} = m_b(\bar{s}_{L(R)}b_{R(L)})(\bar{\mu}\gamma_5\mu),$$

- Scalar operators are strongly constrained by $\text{BR}(B_s \rightarrow \ell\ell)$, if SMEFT relations are imposed
[Alonso et al. 1407.7044]

$$C_S = -C_P \quad C'_S = C'_P$$

- Beyond SMEFT, there are “blind directions” in Wilson coefficient space that are not probed by the purely leptonic decays

[Becirevic et al 1205.5811]



Modification of the B→Kll Angular Distribution and effect on R(K)

$$\frac{1}{\Gamma} \frac{d\Gamma}{d \cos \theta_l} = \frac{3}{4}(1 - F_H)(1 - \cos^2 \theta_l) + \frac{1}{2}F_H + A_{FB} \cos \theta_l$$

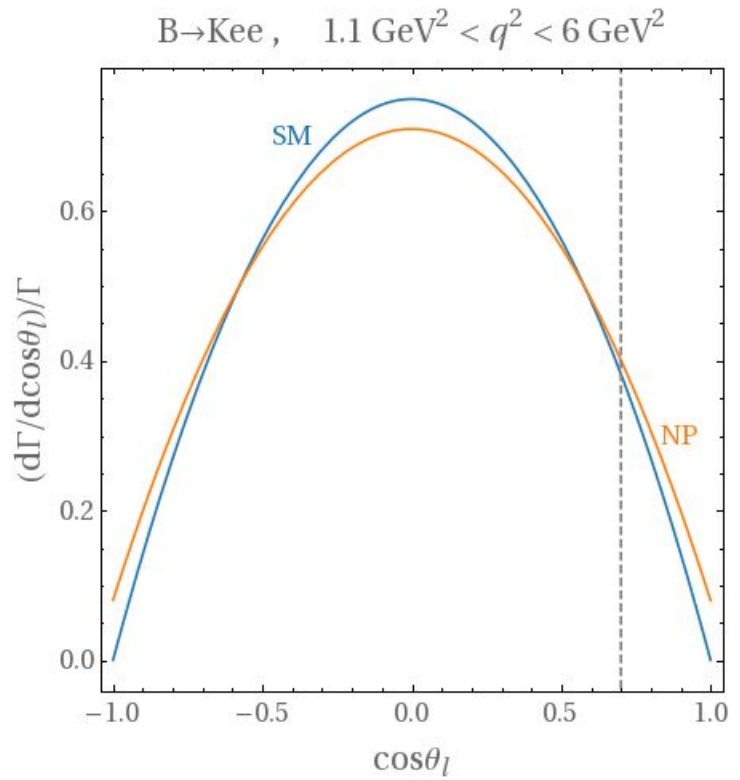
- Example benchmark point with R(K)= 0.85

$$C_S=C_P=C'_S=C'_P= 0.35/\text{GeV}$$

(equality of the Wilson coefficients has to hold at the 10% level to escape the bound from B_s→ee)

- B→Kee efficiency is slightly reduced (~2%) in the new physics example => true R(K) is ~2% lower
 [note: using a cut cos(θ)_l<0.7]

(The new physics scenario is contrived, but I don't think it is excluded)



Modification of the $B \rightarrow K\ell\ell$ Angular Distribution and effect on $R(K)$

$$\frac{1}{\Gamma} \frac{d\Gamma}{d \cos \theta_\ell} = \frac{3}{4}(1 - F_H)(1 - \cos^2 \theta_\ell) + \frac{1}{2}F_H + A_{FB} \cos \theta_\ell$$

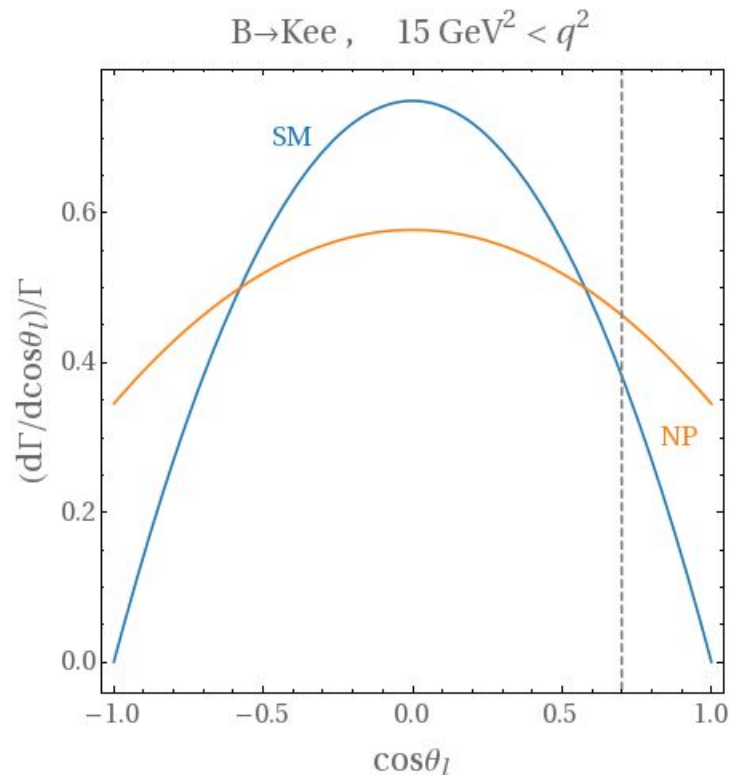
- Example benchmark point with $R(K) = 0.85$

$$C_S = C_P = C'_S = C'_P = 0.35/\text{GeV}$$

(equality of the Wilson coefficients has to hold at the 10% level to escape the bound from $B_s \rightarrow ee$)

- $B \rightarrow K\ell\ell$ efficiency is slightly reduced ($\sim 2\%$) in the new physics example \Rightarrow true $R(K)$ is $\sim 2\%$ lower
[note: using a cut $\cos(\theta_\ell) < 0.7$]
- The effect of scalar operators is more pronounced at high q^2 ($\sim 7\%$ in this example)

(The new physics scenario is contrived, but I don't think it is excluded)



The $B \rightarrow K^* \ell \bar{\ell}$ Angular Distribution

$$\frac{1}{\Gamma} \frac{d\Gamma}{d \cos \theta_\mu} = \frac{3}{4} F_L (1 - \cos^2 \theta_\mu) + \frac{3}{8} (1 - F_L) (1 + \cos^2 \theta_\mu) + A_{FB} \cos \theta_\mu$$

- The $\cos(\theta_\mu)$ distribution depends on:
 1. the longitudinal polarization fraction of the K^*
 2. the forward backward asymmetry
- F_L and A_{FB} can be modified by the usual semi-leptonic operators

$$O_9^{bs\ell\ell} = (\bar{s} \gamma_\mu P_L b) (\bar{\ell} \gamma^\mu \ell),$$

$$O_{10}^{bs\ell\ell} = (\bar{s} \gamma_\mu P_L b) (\bar{\ell} \gamma^\mu \gamma_5 \ell),$$

$$O_9^{'bs\ell\ell} = (\bar{s} \gamma_\mu P_R b) (\bar{\ell} \gamma^\mu \ell),$$

$$O_{10}^{'bs\ell\ell} = (\bar{s} \gamma_\mu P_R b) (\bar{\ell} \gamma^\mu \gamma_5 \ell),$$

Modification of the $B \rightarrow K^* \ell \ell$ Angular Distribution and effect on $R(K^*)$

$$\frac{1}{\Gamma} \frac{d\Gamma}{d \cos \theta_\mu} = \frac{3}{4} F_L (1 - \cos^2 \theta_\mu) + \frac{3}{8} (1 - F_L) (1 + \cos^2 \theta_\mu) + A_{FB} \cos \theta_\mu$$

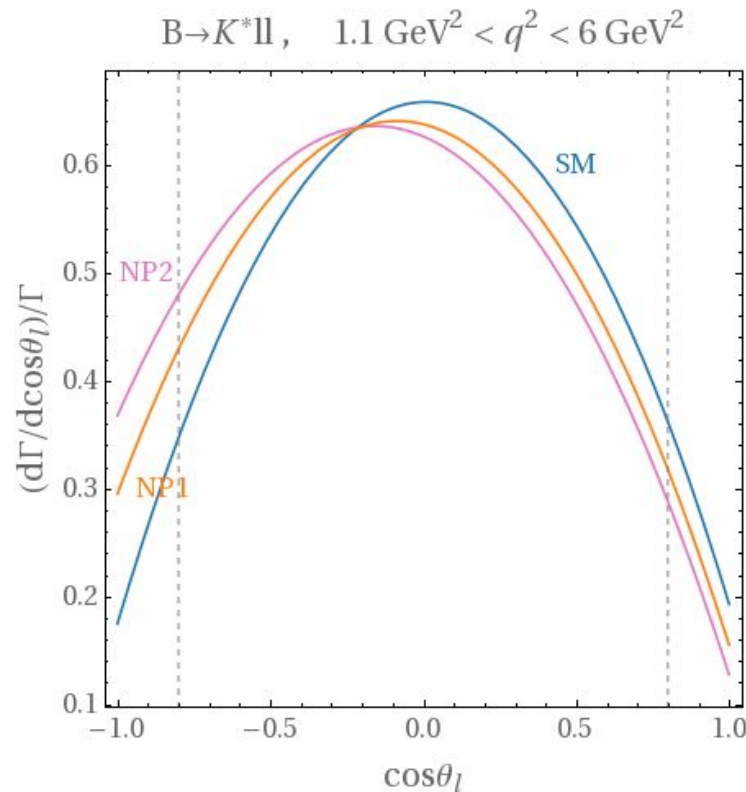
- Two new physics examples

1) $C_9^{bs\mu\mu} = -1.0$ $\rightarrow R(K^*) = 0.85$
2) $C_9^{bsee} = C_{10}^{bsee} = -1.5$ $\rightarrow R(K^*) = 0.66$

- At low q^2 : $1.1 \text{ GeV}^2 < q^2 < 6 \text{ GeV}^2$

In new physics **example 1**, the $B \rightarrow K^* \mu \mu$ efficiency is $\sim 1\%$ smaller \Rightarrow **true $R(K^*)$ is $\sim 1\%$ higher**

In new physics **example 2**, the $B \rightarrow K^* e e$ efficiency is $\sim 2\%$ smaller \Rightarrow **true $R(K^*)$ is $\sim 2\%$ lower**



Modification of the $B \rightarrow K^* \ell \ell$ Angular Distribution and effect on $R(K^*)$

$$\frac{1}{\Gamma} \frac{d\Gamma}{d \cos \theta_\mu} = \frac{3}{4} F_L (1 - \cos^2 \theta_\mu) + \frac{3}{8} (1 - F_L) (1 + \cos^2 \theta_\mu) + A_{FB} \cos \theta_\mu$$

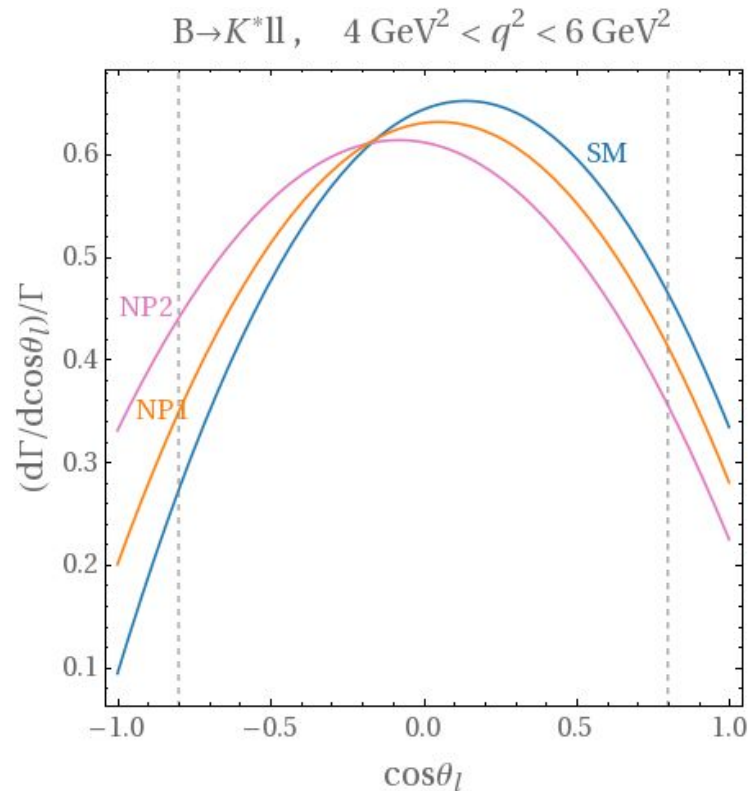
- Two new physics examples

1) $C_9^{bs\mu\mu} = -1.0$ $\rightarrow R(K^*) = 0.85$
2) $C_9^{bsee} = C_{10}^{bsee} = -1.5$ $\rightarrow R(K^*) = 0.66$

- At low q^2 : $4 \text{ GeV}^2 < q^2 < 6 \text{ GeV}^2$

In new physics **example 1**, the $B \rightarrow K^* \mu \mu$ efficiency is $\sim 1\%$ smaller \Rightarrow true $R(K^*)$ is $\sim 1\%$ higher

In new physics **example 2**, the $B \rightarrow K^* e e$ efficiency is $\sim 2\%$ smaller \Rightarrow true $R(K^*)$ is $\sim 2\%$ lower



Modification of the $B \rightarrow K^* \ell \ell$ Angular Distribution and effect on $R(K^*)$

$$\frac{1}{\Gamma} \frac{d\Gamma}{d \cos \theta_\mu} = \frac{3}{4} F_L (1 - \cos^2 \theta_\mu) + \frac{3}{8} (1 - F_L) (1 + \cos^2 \theta_\mu) + A_{FB} \cos \theta_\mu$$

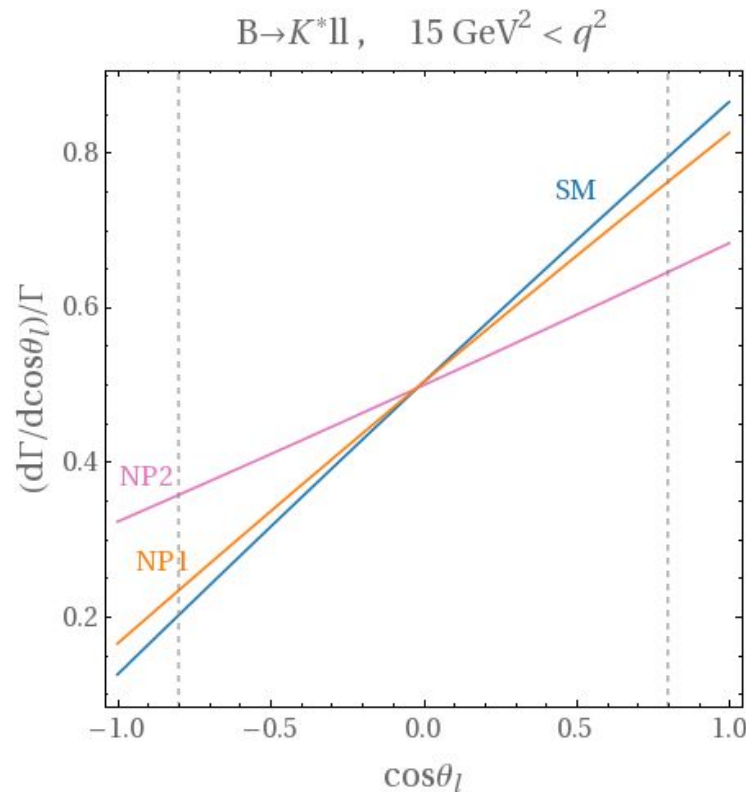
- Two new physics examples

1) $C_9^{bs\mu\mu} = -1.0$ $\rightarrow R(K^*) = 0.85$
2) $C_9^{bsee} = C_{10}^{bsee} = -1.5$ $\rightarrow R(K^*) = 0.66$

- At high q^2 : $15 \text{ GeV}^2 < q^2$

In new physics **example 1**, the $B \rightarrow K^* \mu \mu$ efficiency is \sim unchanged \Rightarrow **true $R(K^*)$ is \sim unchanged**

In new physics **example 2**, the $B \rightarrow K^* e e$ efficiency is \sim unchanged \Rightarrow **true $R(K^*)$ is \sim unchanged**



Light New Physics?

- Can light new physics that is not covered by the effective Hamiltonian formalism affect the measurements of $R(K)$ and $R(K^*)$? (dark photons, axions, light Z' bosons, ...)

[Sala, Straub 1704.06188; Datta et al. 1705.08423; WA, Baker, Gori, Harnik, Pospelov, Thamm 1711.07494]

- Typically one would expect a prominent bump in the q^2 -spectrum.
- No new physics resonance is seen in $B \rightarrow K^* X$, $X \rightarrow \mu\mu$ and $B \rightarrow K X$, $X \rightarrow \mu\mu$
- What about electrons?

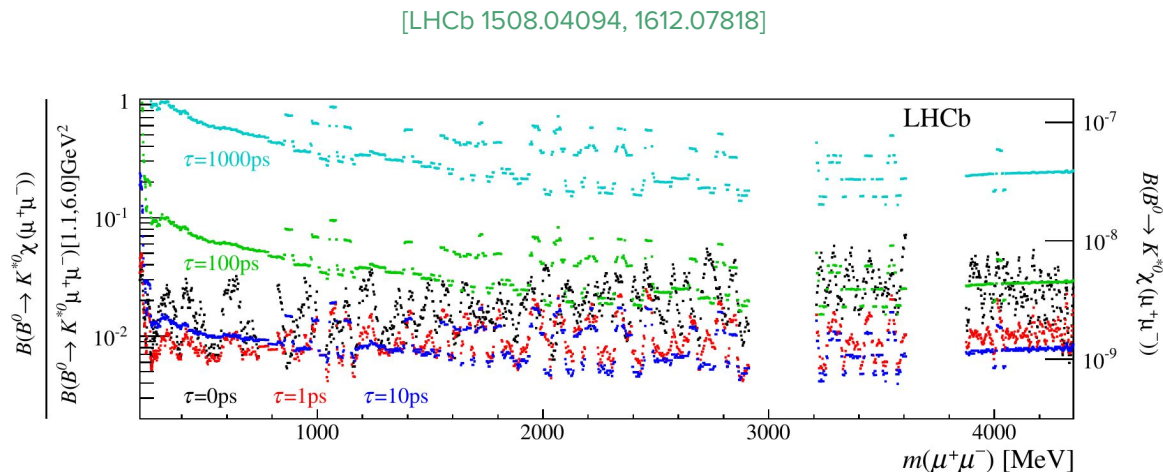


Figure 7: Upper limits at 95% CL for (left axis) $\mathcal{B}(B^0 \rightarrow K^{*0} \chi(\mu^+ \mu^-)) / \mathcal{B}(B^0 \rightarrow K^{*0} \mu^+ \mu^-)$, with $B^0 \rightarrow K^{*0} \mu^+ \mu^-$ in $1.1 < m^2(\mu^+ \mu^-) < 6.0 \text{ GeV}^2$, and (right axis) $\mathcal{B}(B^0 \rightarrow K^{*0} \chi(\mu^+ \mu^-))$. Same as Fig. 4 in the Letter but including the $\tau = 0$ and 1 ps limits.

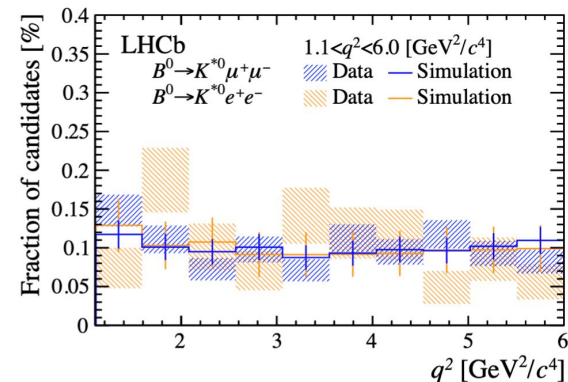
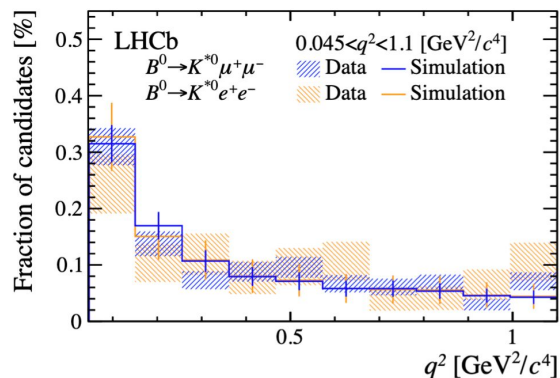
Light New Physics?

- Can light new physics that is not covered by the effective Hamiltonian formalism affect the measurements of $R(K)$ and $R(K^*)$? (dark photons, axions, light Z' bosons, ...)

[Sala, Straub 1704.06188; Datta et al. 1705.08423; WA, Baker, Gori, Harnik, Pospelov, Thamm 1711.07494]

- Typically one would expect a prominent bump in the q^2 -spectrum.
- No new physics resonance is seen in $B \rightarrow K^* X$, $X \rightarrow \mu\mu$ and $B \rightarrow K X$, $X \rightarrow \mu\mu$
- What about electrons?

[LHCb, JHEP 08 (2017) 055]



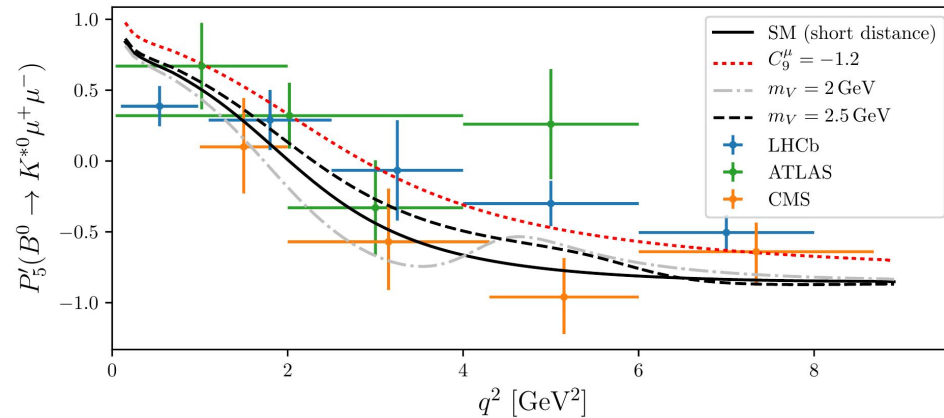
Light New Physics and the Angular Distribution

The constraints from the q^2 Spectrum can be avoided in several ways:

← increasing "exoticness"

- Hide the light new physics particle behind QCD resonances.
- Put the new particle close to thresholds.
- Give the new particle a very large width.
- Introduce many resonances with mass splitting of the order of the invariant mass resolution.
- Consider an "unparticle continuum".

[Sala, Straub 1704.06188]



Could light new physics lead to an "exotic" distortion of the angular distribution and affect the $R(K)$ or $R(K^*)$ efficiencies in a peculiar way?

Some other thoughts...

- C/C' interplay for different helicity amplitudes?

$$C + C' : K, K_{\perp}^*, \dots$$

$$C - C' : K_0(1430), K_{0,\parallel}^*, \dots$$

[Hiller, Schmaltz, JHEP 02 (2015) 055]

- Impact of S-wave on B->K*0|| - not separated in R_{K*}

- measured in B⁰->Kπμμ to be ~10%

[LHCb, JHEP 04 (2017) 142]

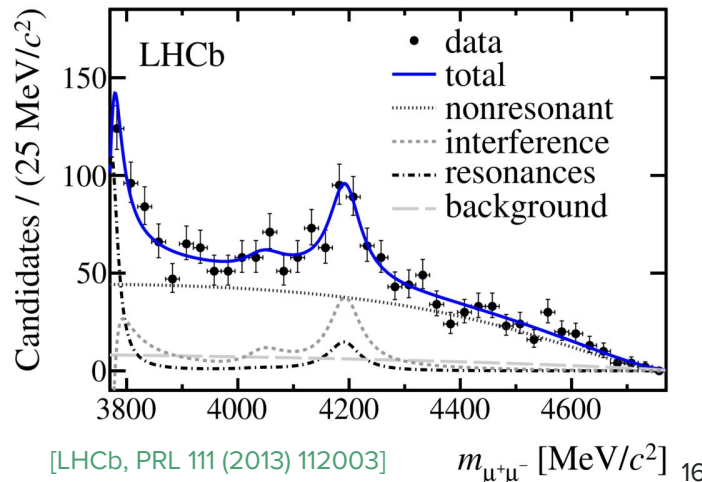
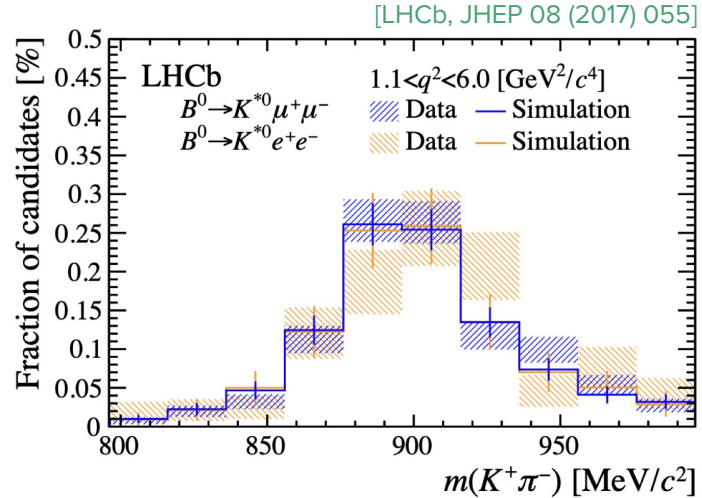
- higher waves?

[No significant D-wave in B⁰->Kπμμ at high m(Kπ)]

[LHCb, JHEP 04 (2017) 142]

- For LFU at high q²: What is the impact of resonances?

- How to include these bins in global fits?



[LHCb, PRL 111 (2013) 112003]

$m_{\mu^+\mu^-}$ [MeV/c²] 16