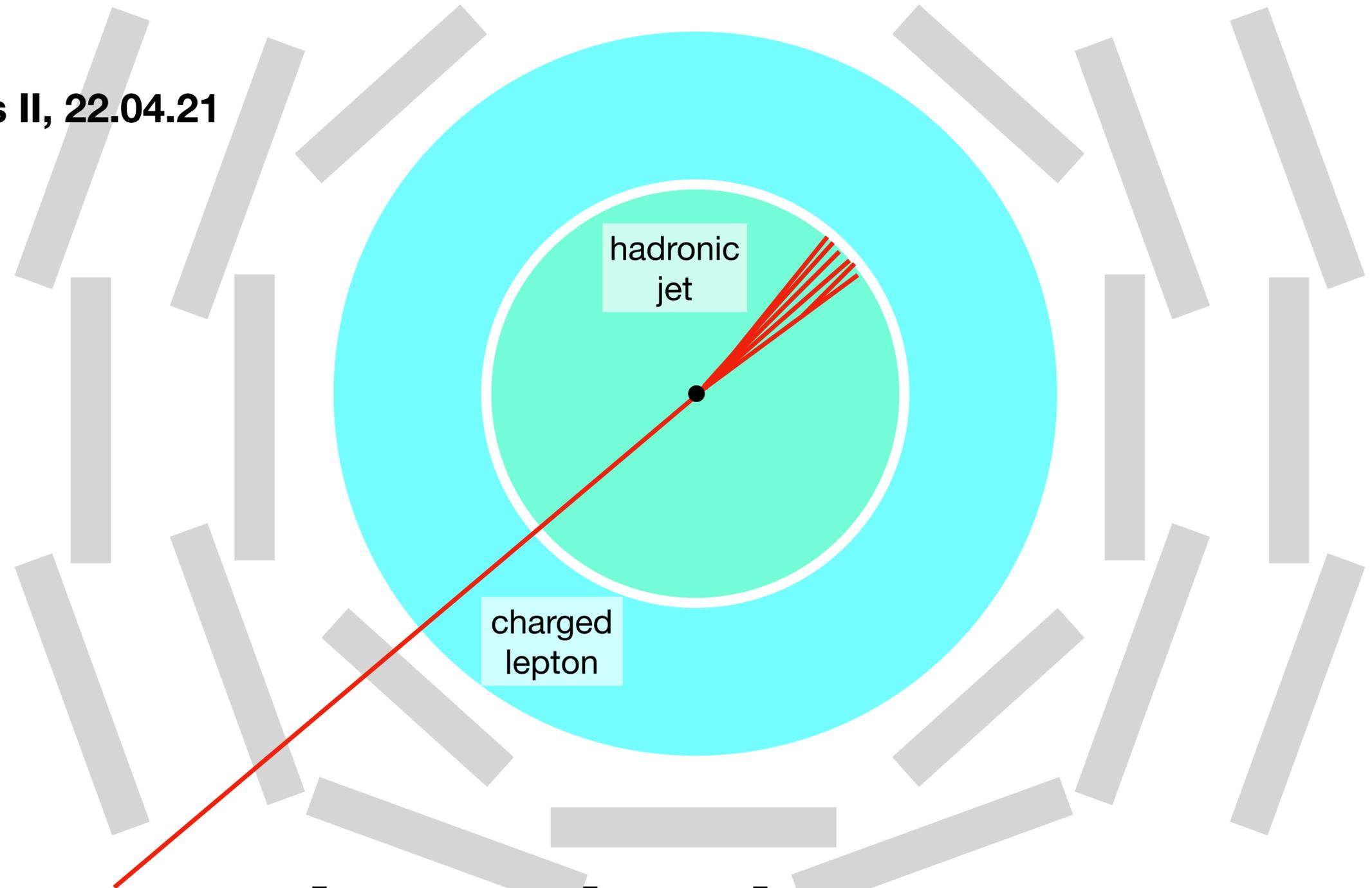


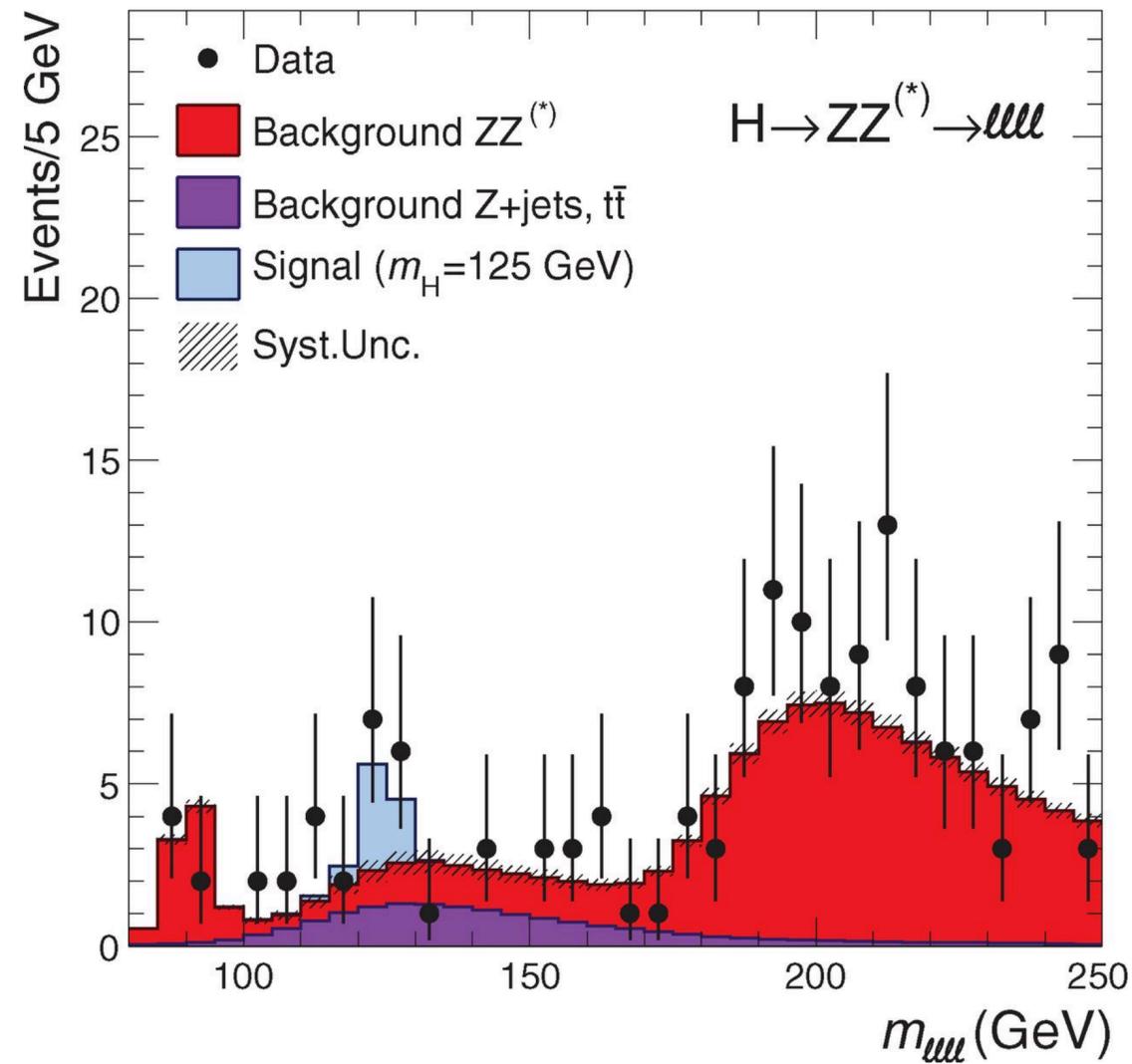
Uli Haisch, MPI Munich  
Beyond the Flavour Anomalies II, 22.04.21



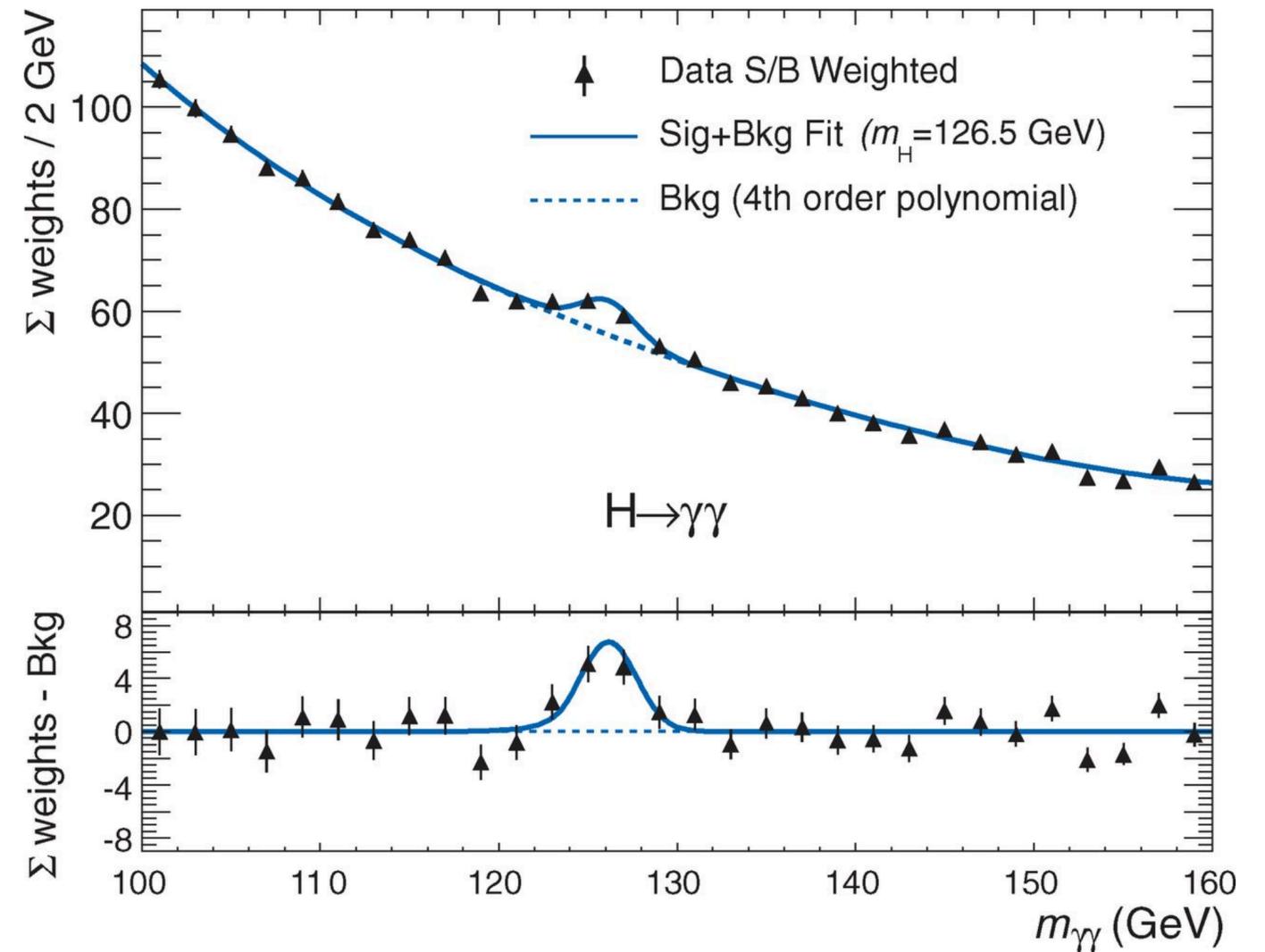
# Resonant leptoquark production

## LHC applications & constraints on B-physics anomalies

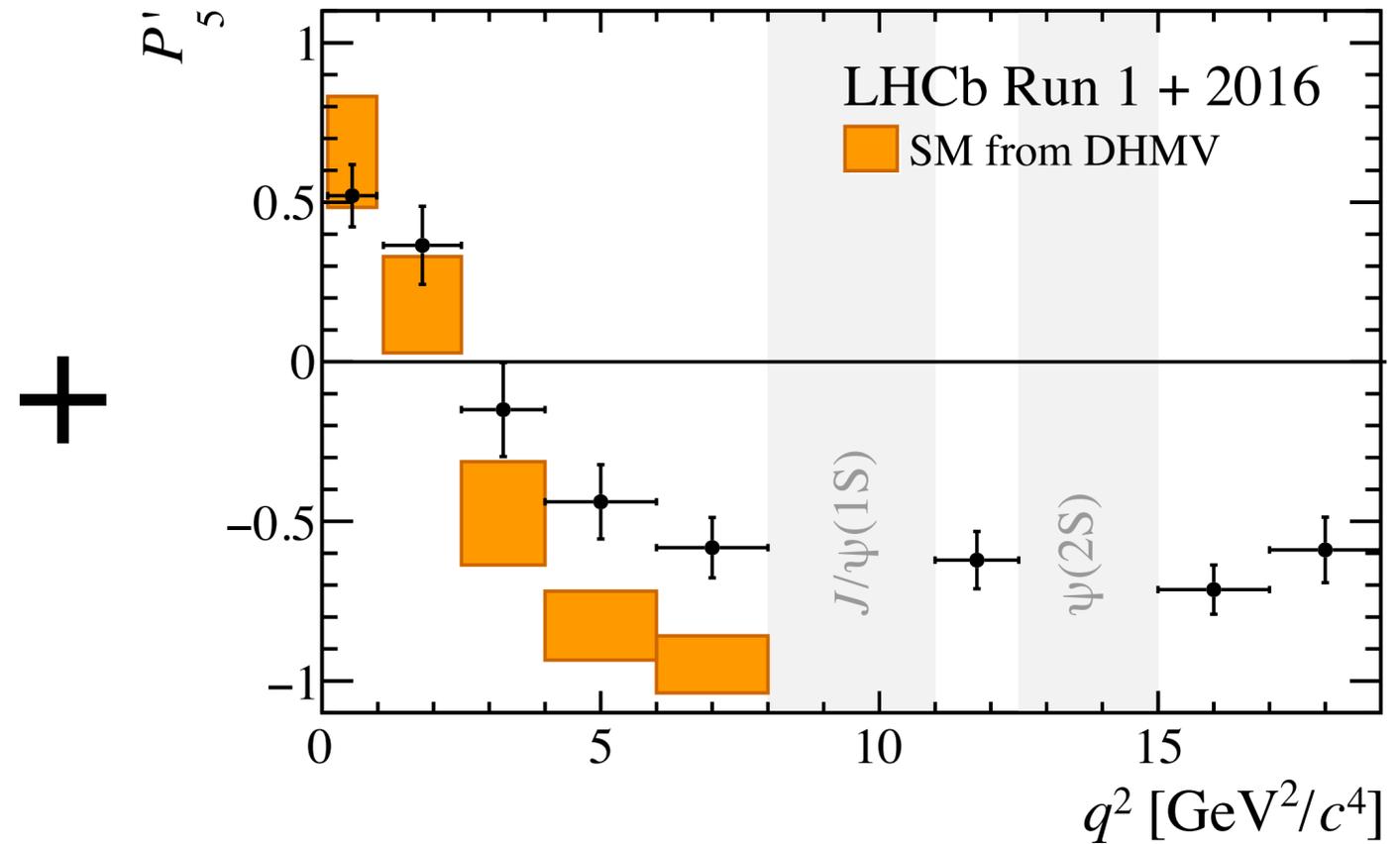
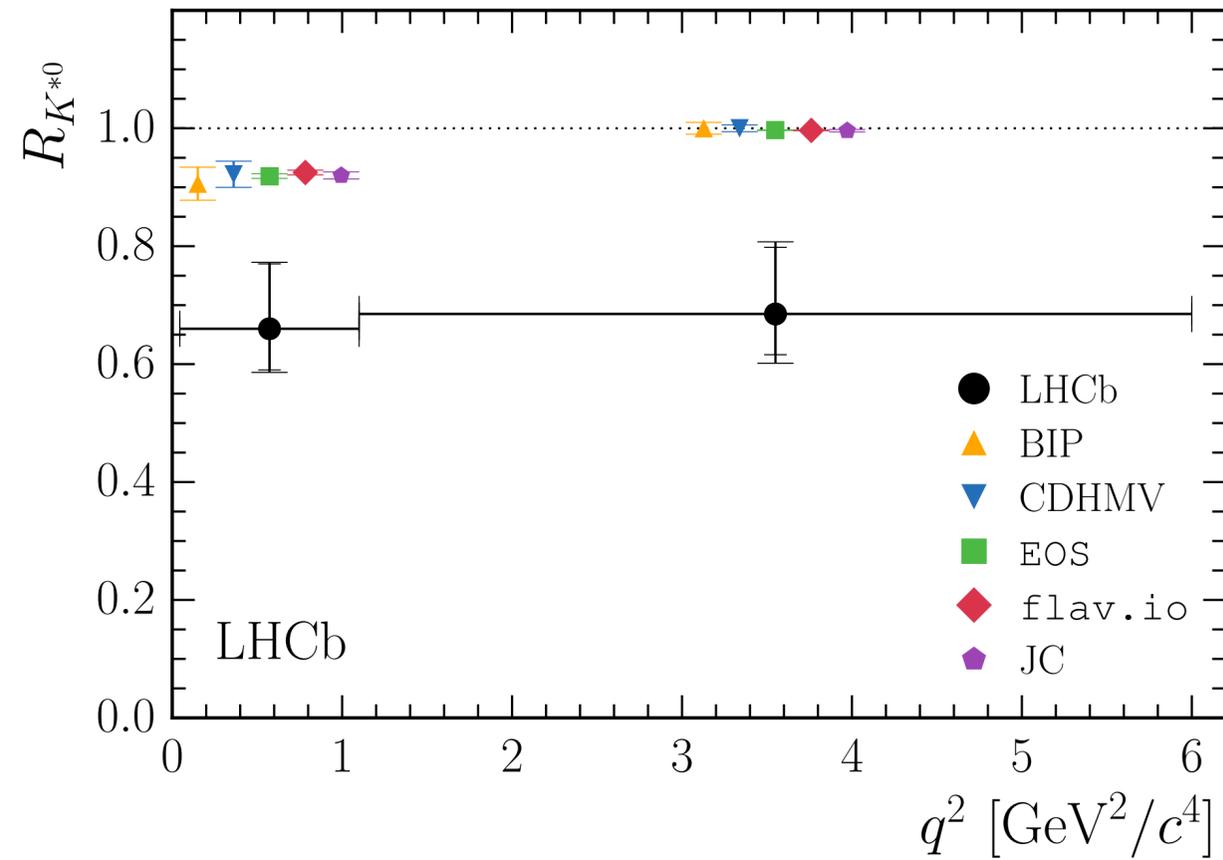
# 5 $\sigma$ : ATLAS Higgs discovery



+

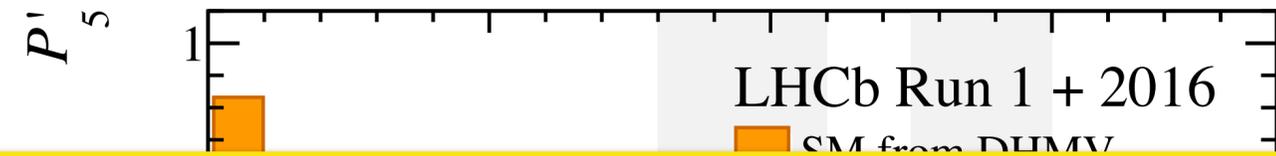
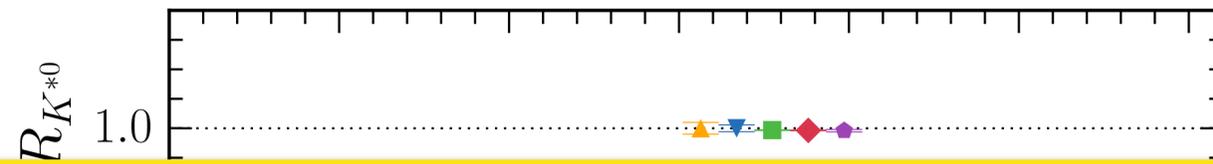


# 5 $\sigma$ : $b \rightarrow s$ anomalies



+ order 100 other observables

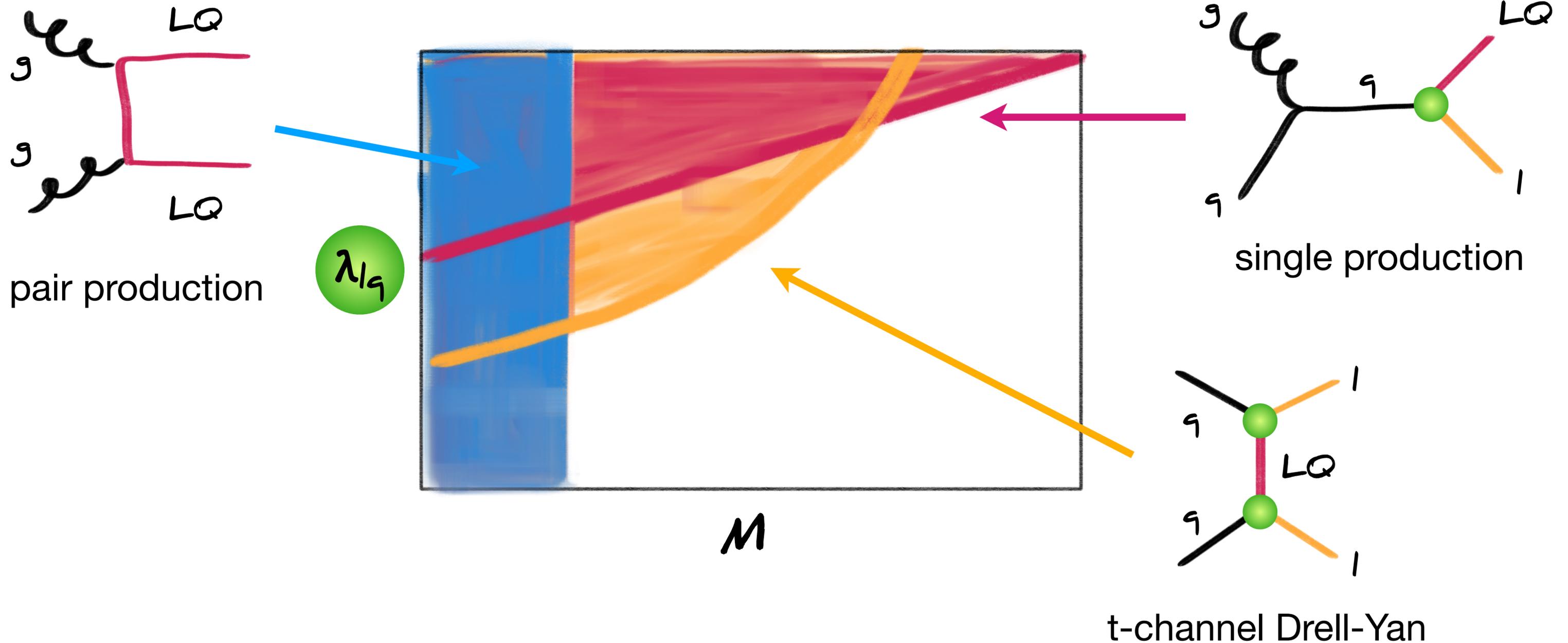
# 5 $\sigma$ : $b \rightarrow s$ anomalies



While  $b \rightarrow s$  anomalies may have similar significance than data that led to Higgs discovery, there are at least two important differences. First, Higgs has been discovered by two independent experiments & second Higgs has been detected by observing a resonance in two different final states. Case of flavour anomalies would be significantly stronger IMHO, if ATLAS/CMS would also see hints of them, in best-case scenario by finding a bump in a high- $p_T$  search

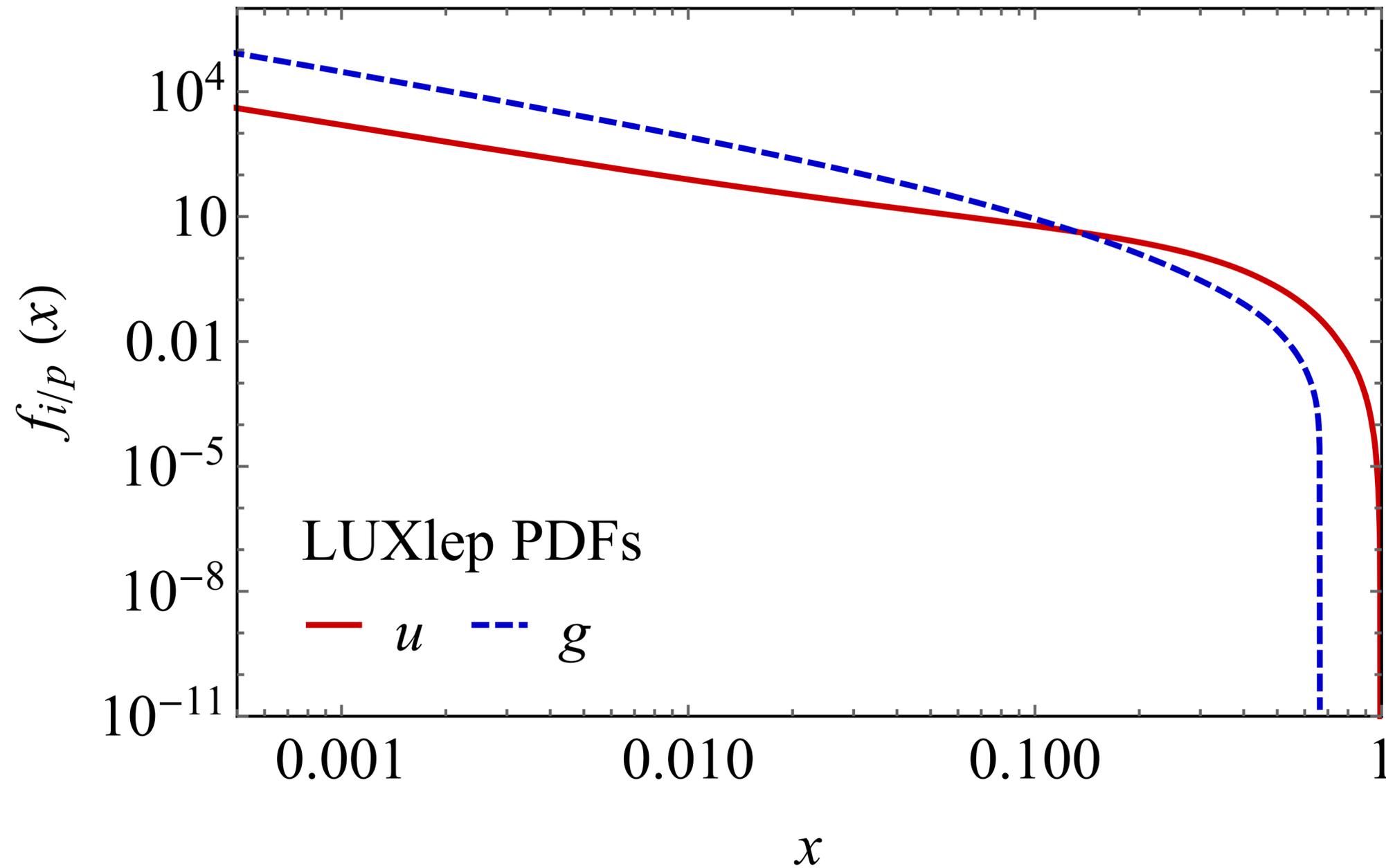
+ order 100 other observables

# Leptoquark (LQ) search strategies @ the LHC

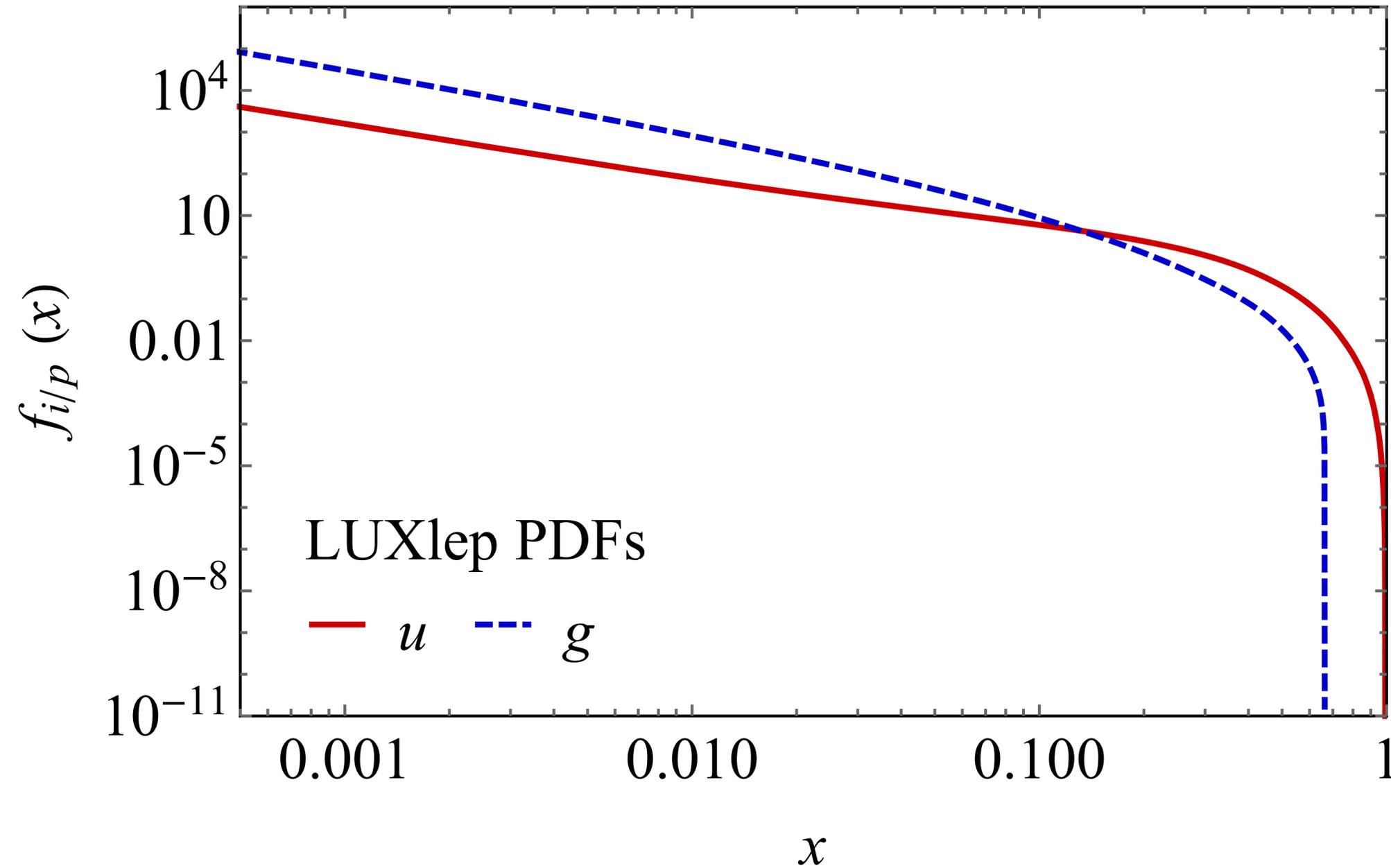


[sketch adopted from Dorsner & Greljo, JHEP 05 (2018) 126]

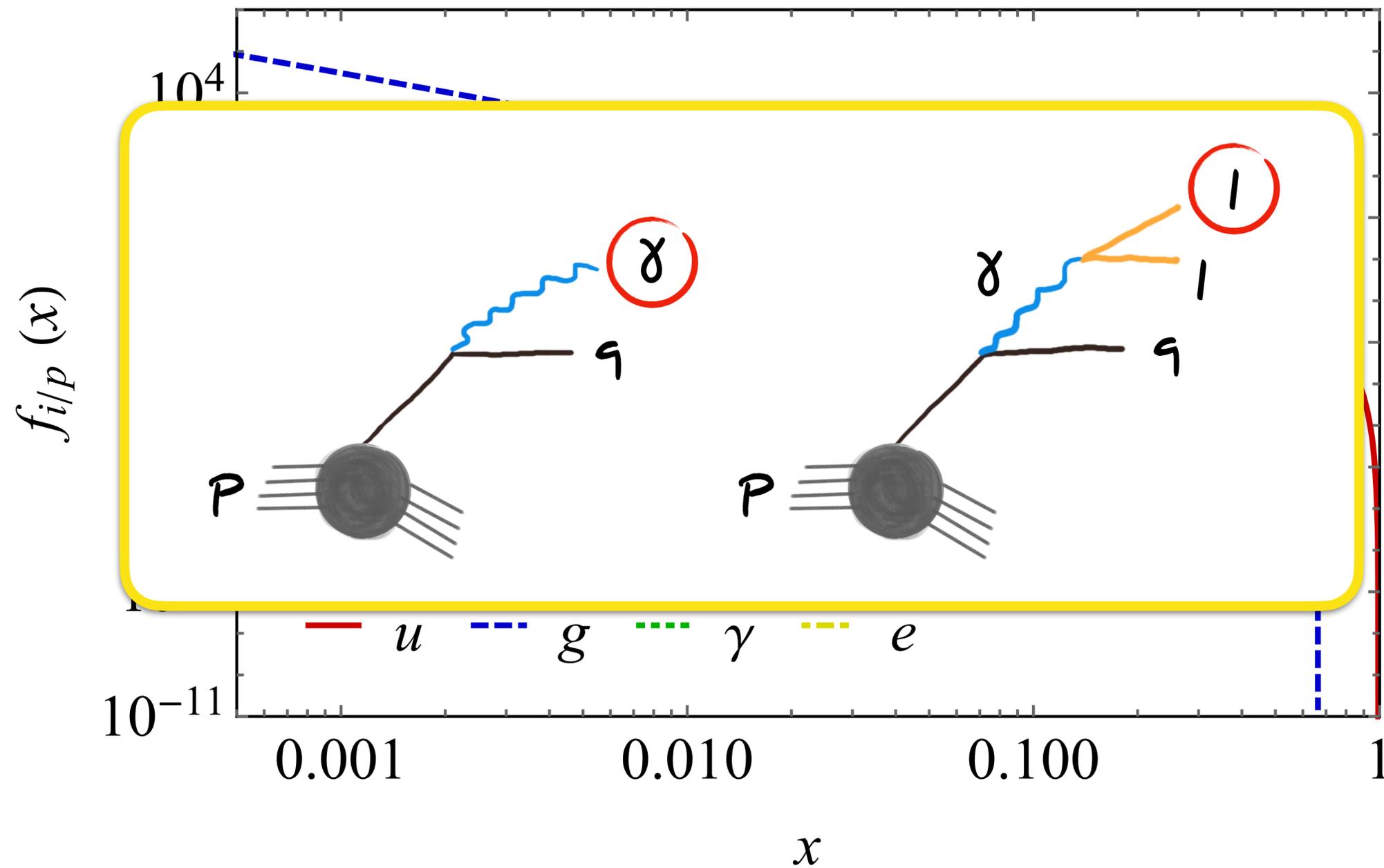
# But @ LHC no resonant LQ production ...



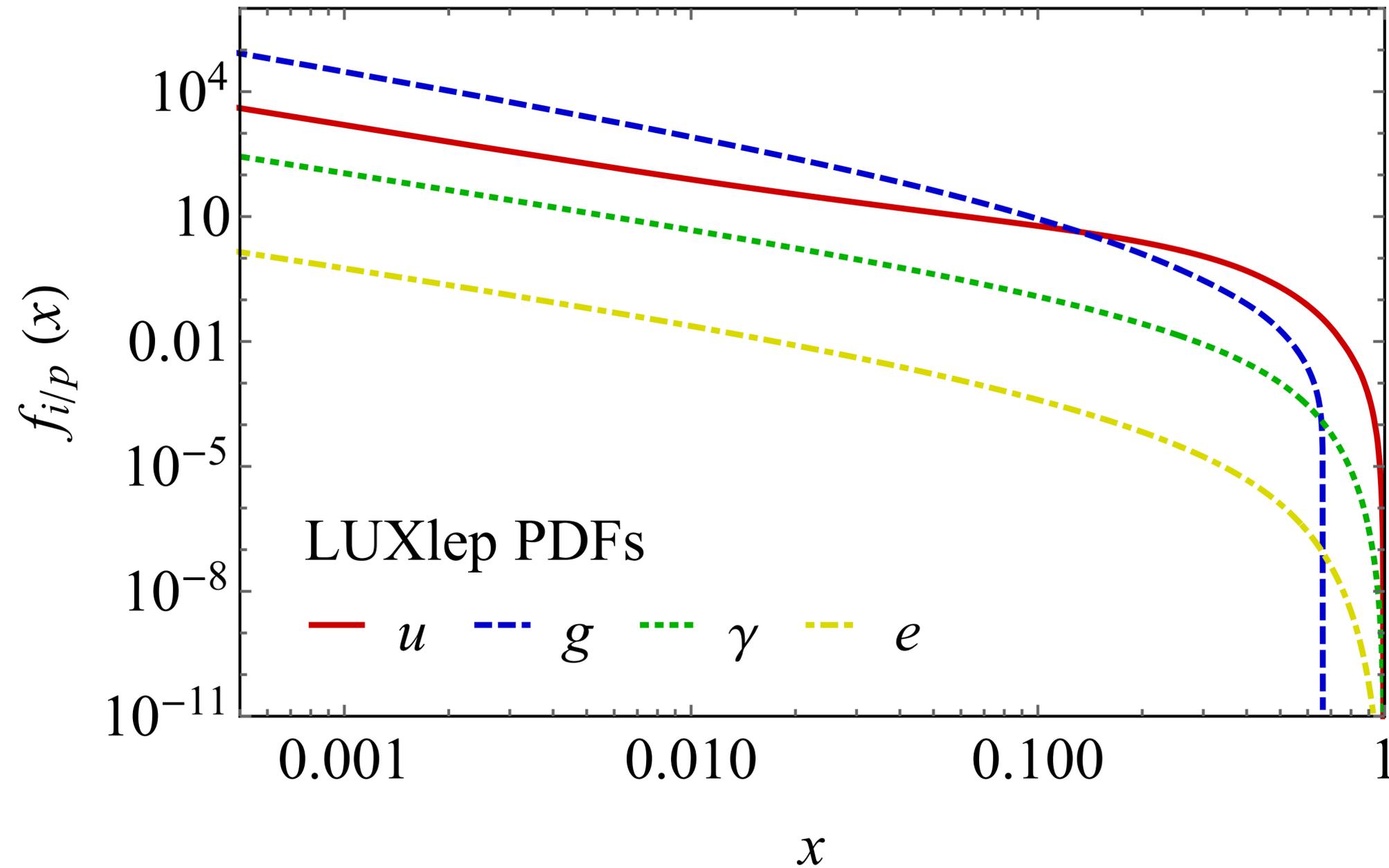
... since a proton consists of quarks & gluons



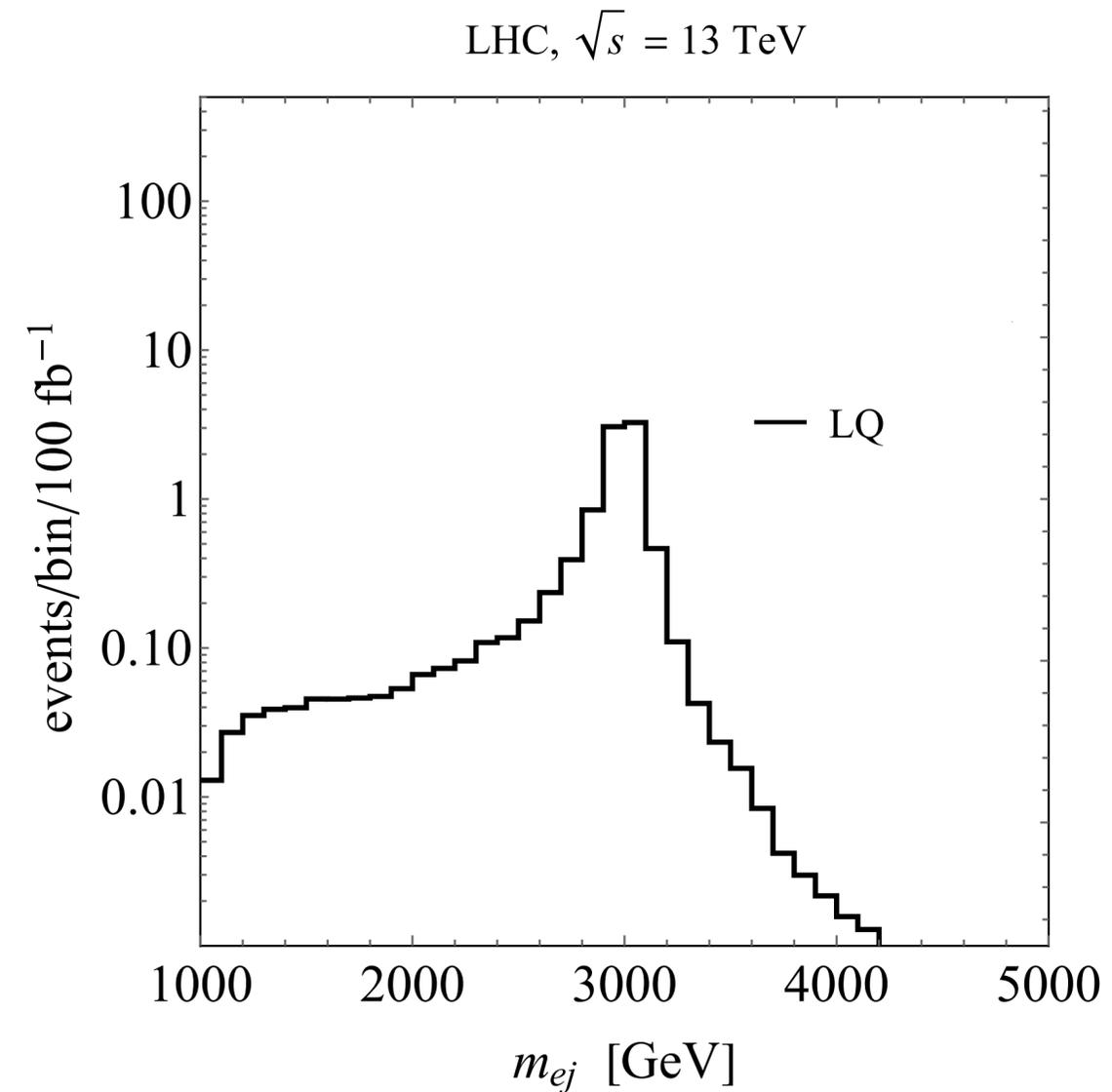
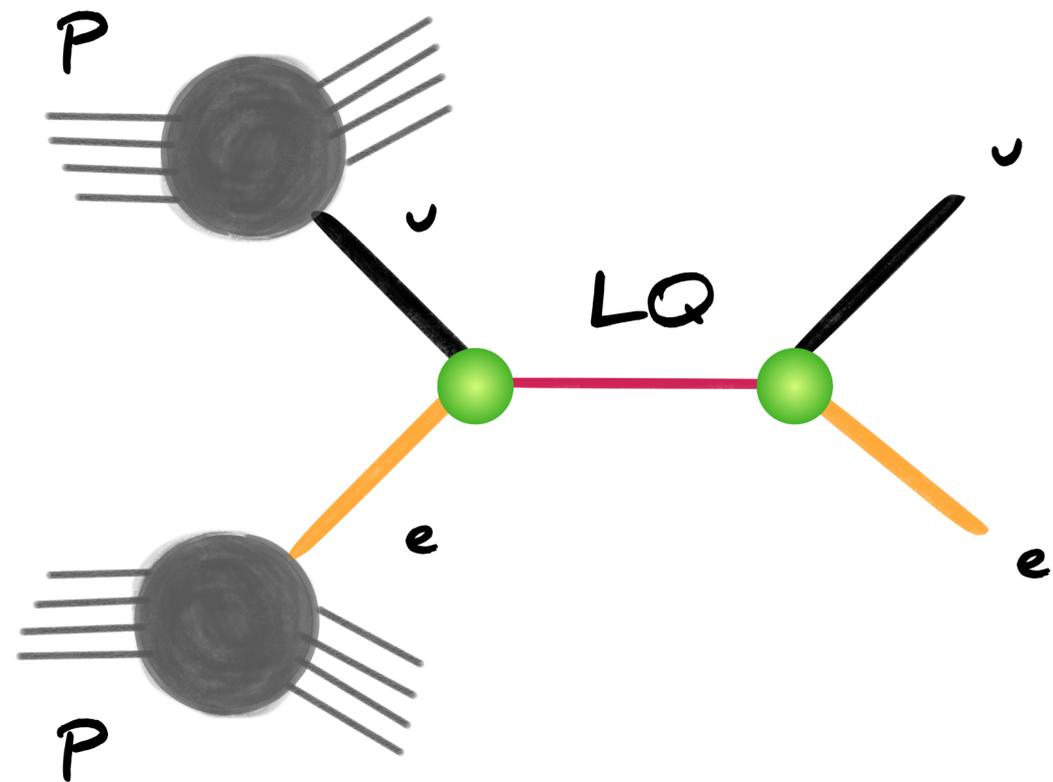
# QFT to the rescue!



# Proton has a little bit of photons & leptons!

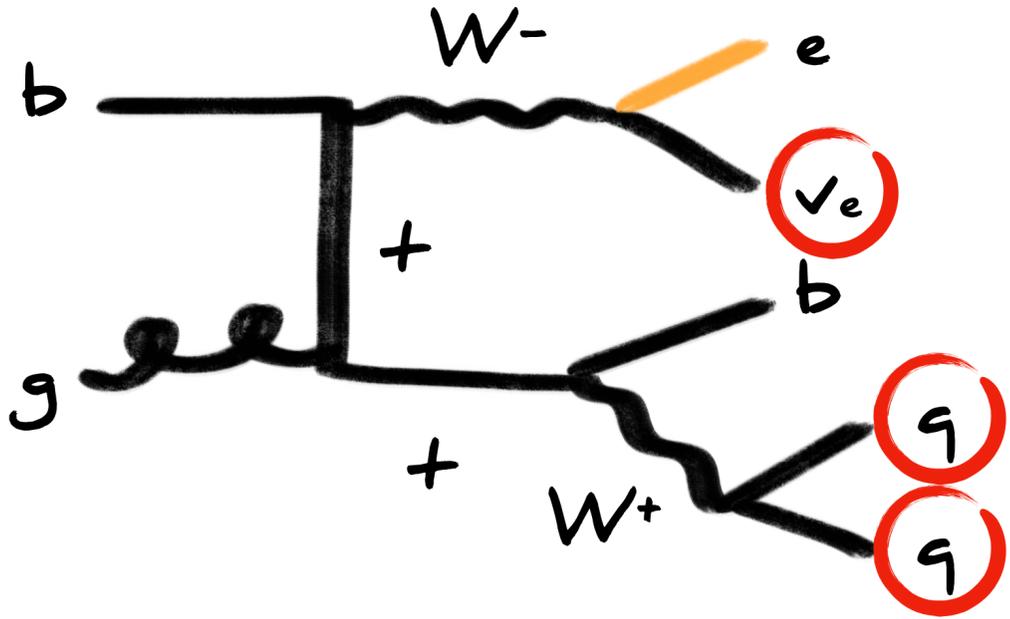


# Resonant LQ production @ the LHC



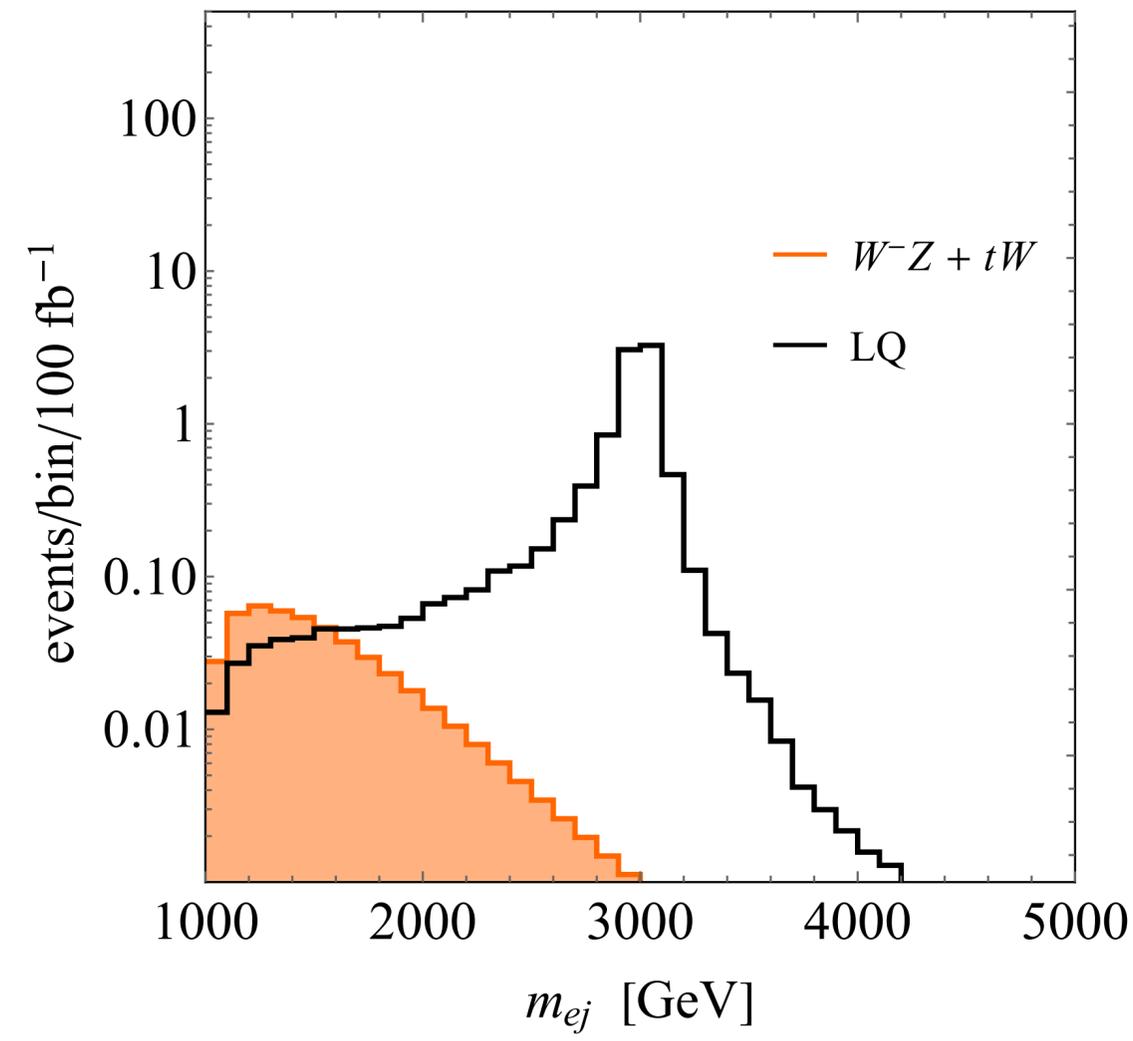
At 13 TeV LHC, 9 events per 100 fb<sup>-1</sup> for minimal scalar LQ of  $M = 3$  TeV &  $\lambda_{eu} = 1$

# Resonant LQ production @ the LHC



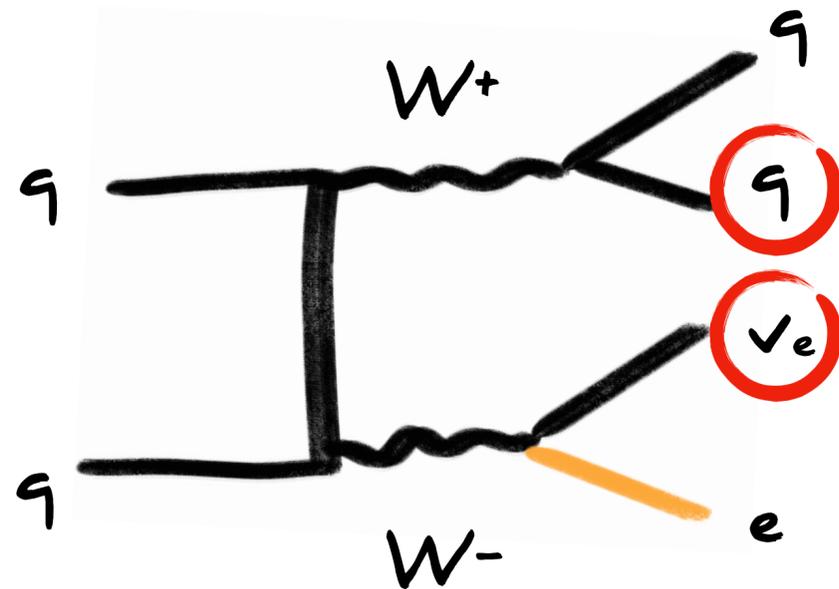
Suppressed by  $E_{T,miss}$  requirement & jet veto

LHC,  $\sqrt{s} = 13$  TeV

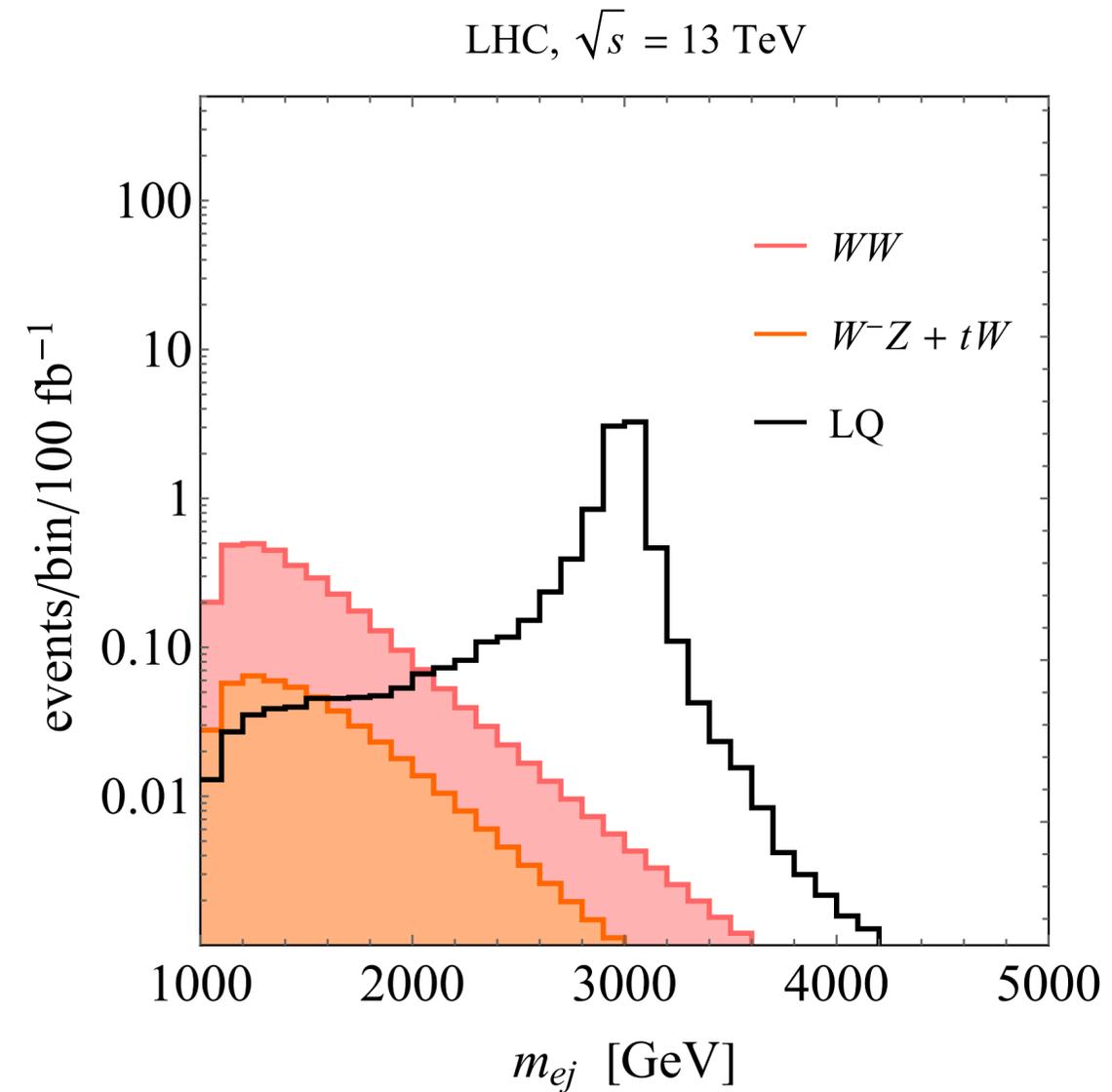


[Buonocore, UH, Nason, Tramontano & Zanderighi, PRL 125 (2020) 23]

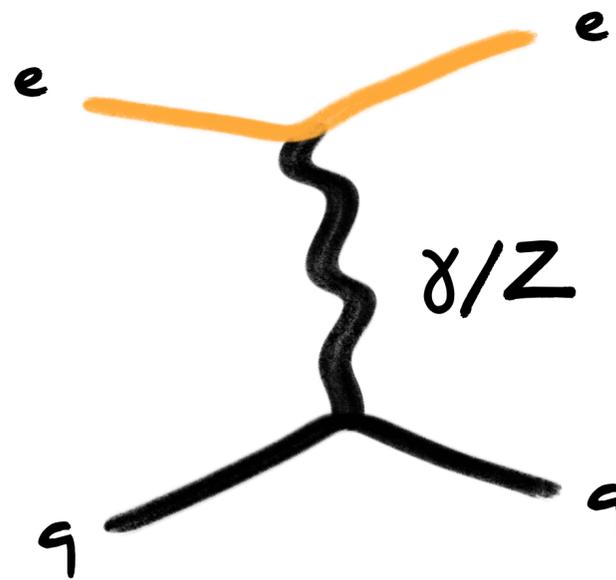
# Resonant LQ production @ the LHC



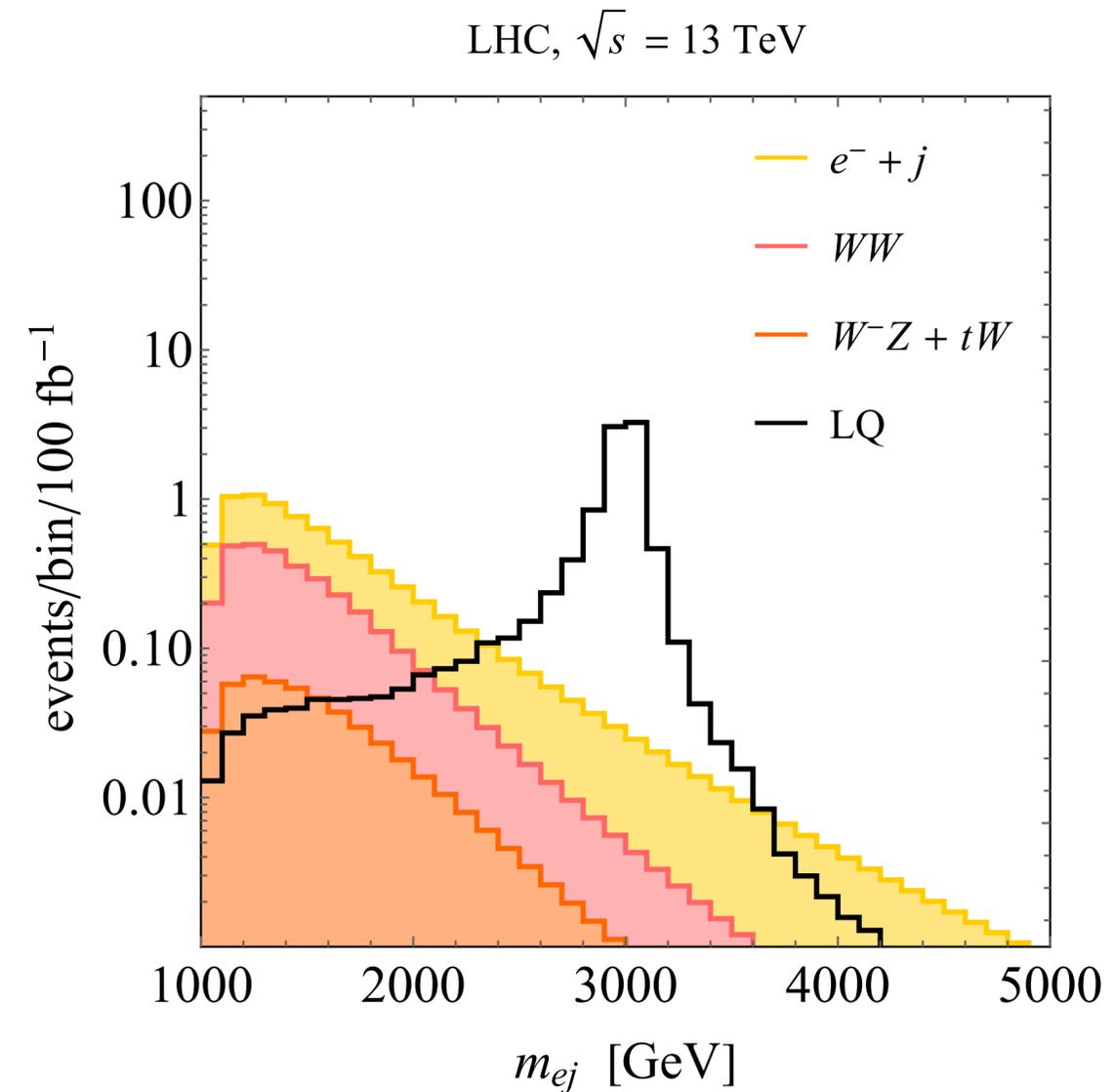
Suppressed by  $E_{T,miss}$  requirement & jet veto



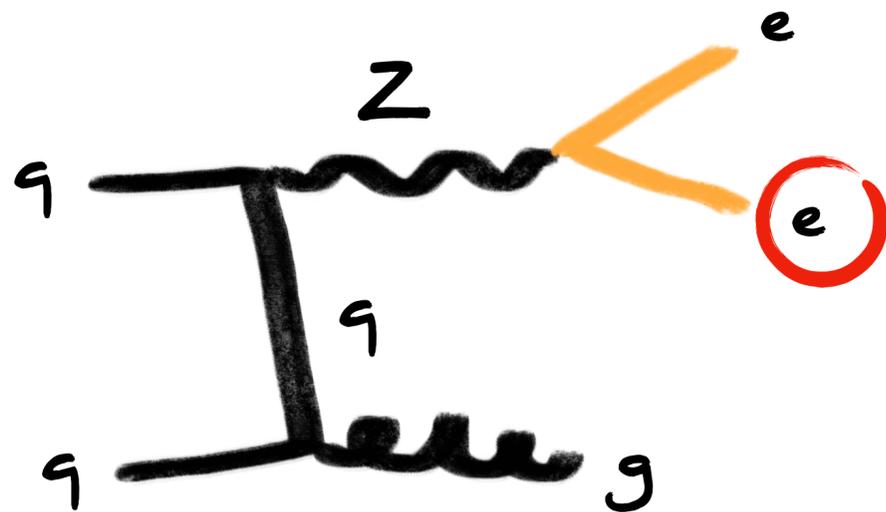
# Resonant LQ production @ the LHC



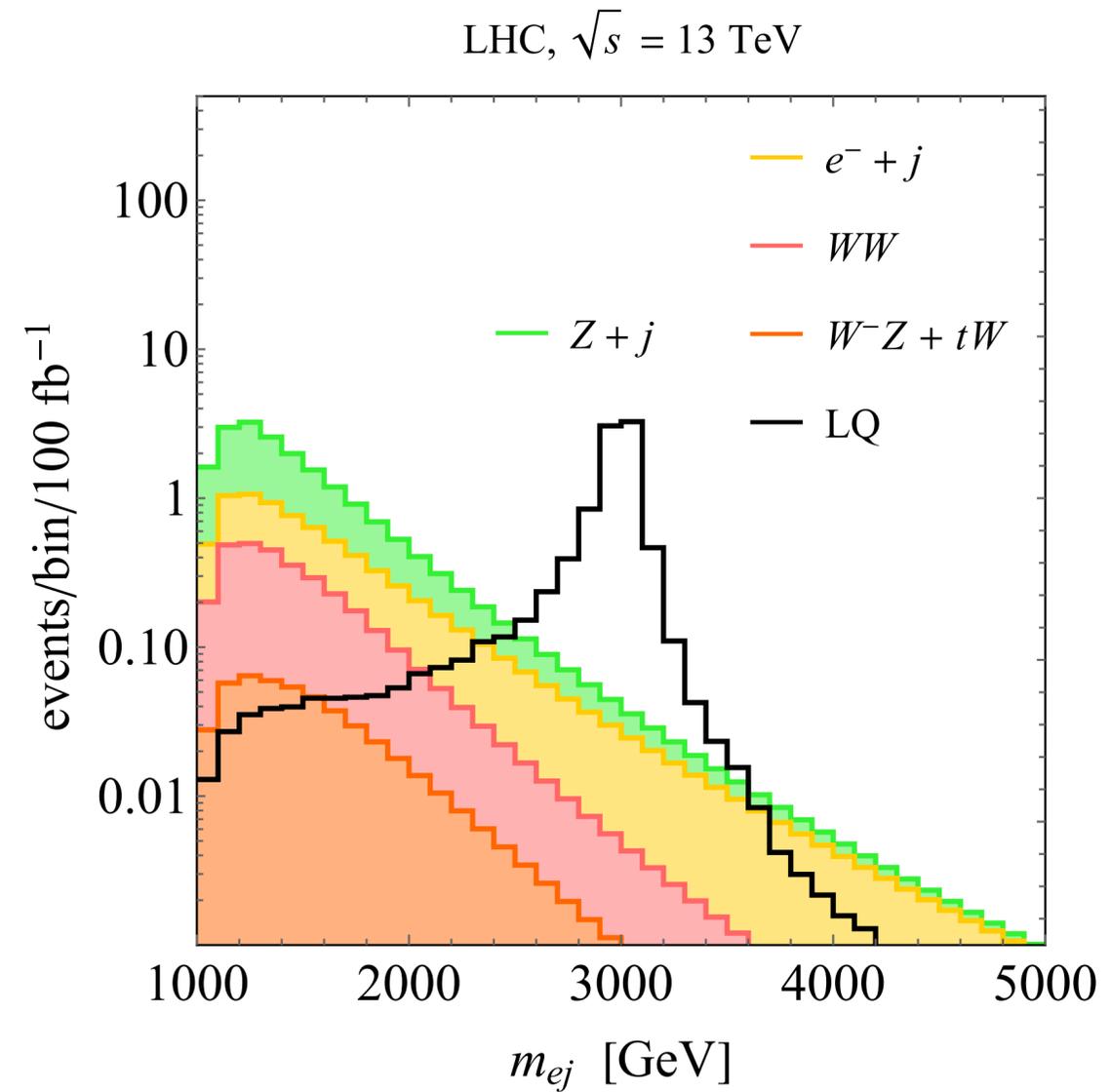
Irreducible background particularly relevant  
@ high invariant lepton-jet mass



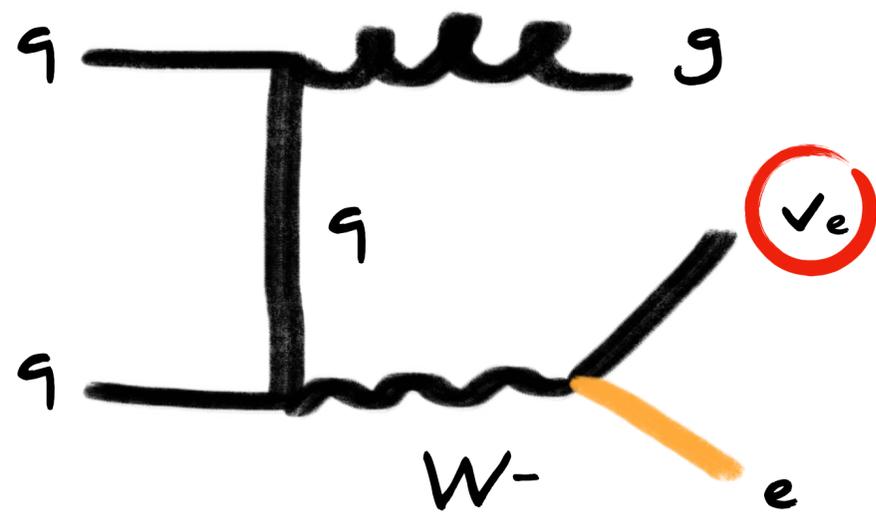
# Resonant LQ production @ the LHC



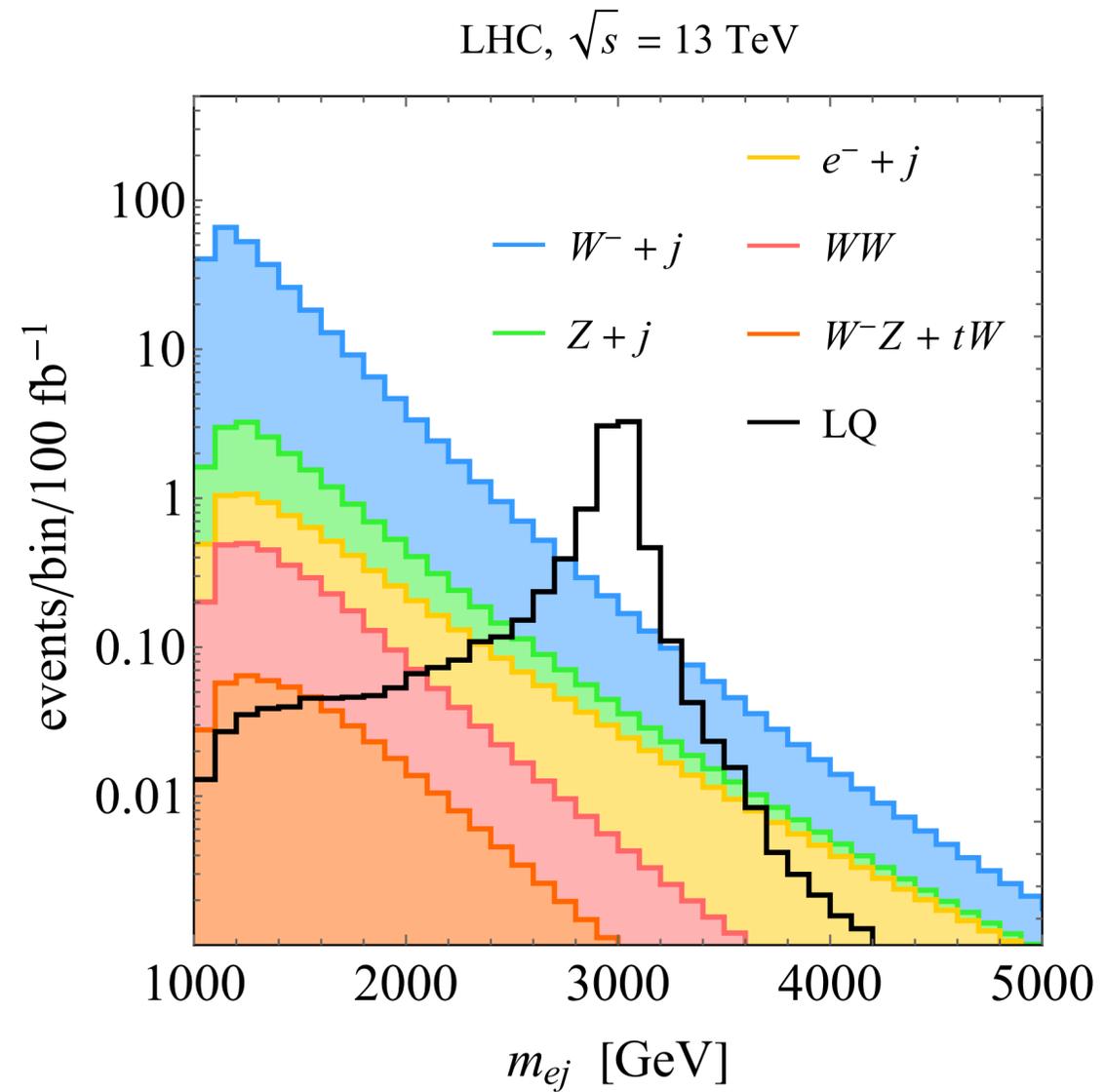
Suppressed by lepton veto



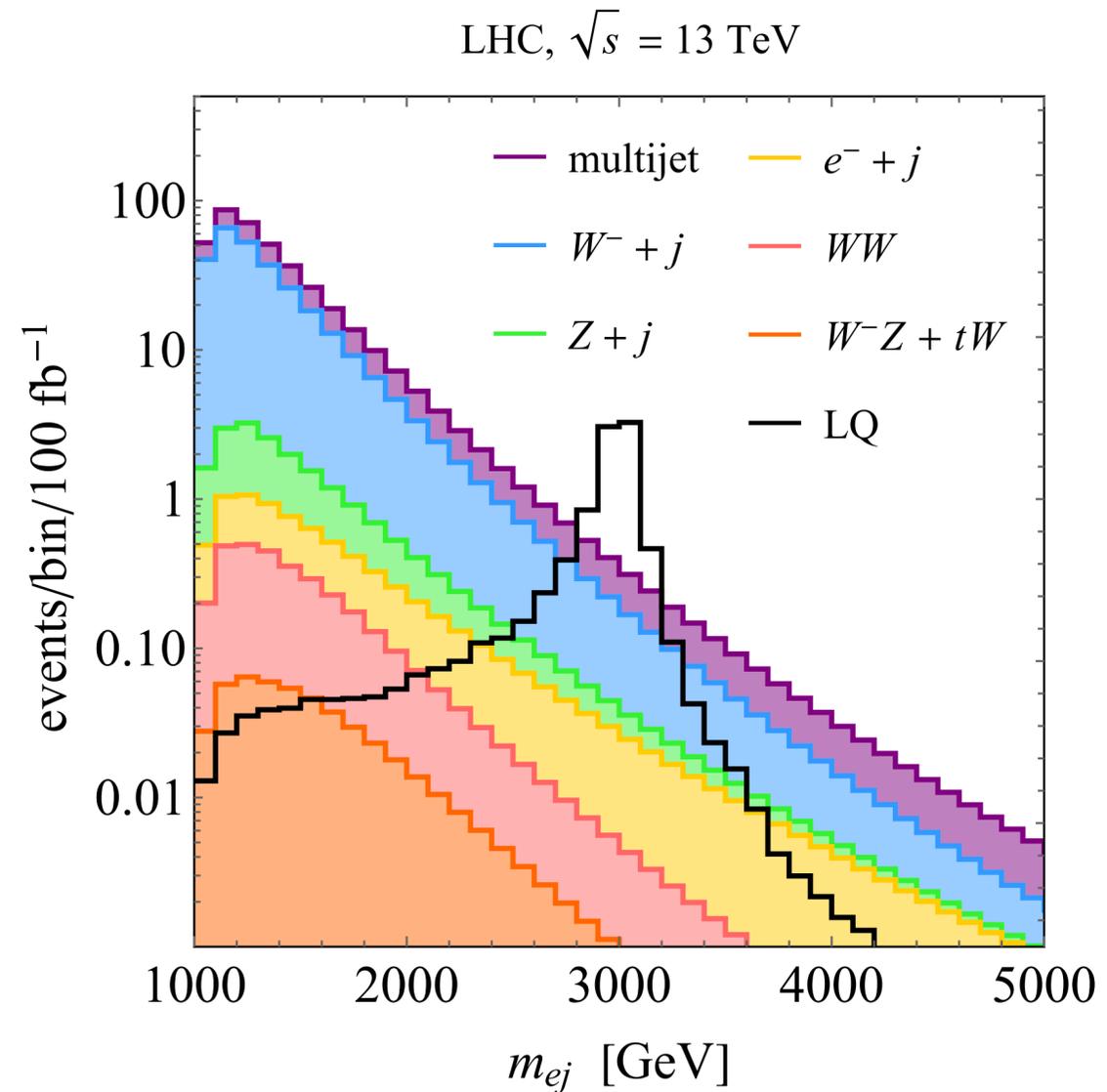
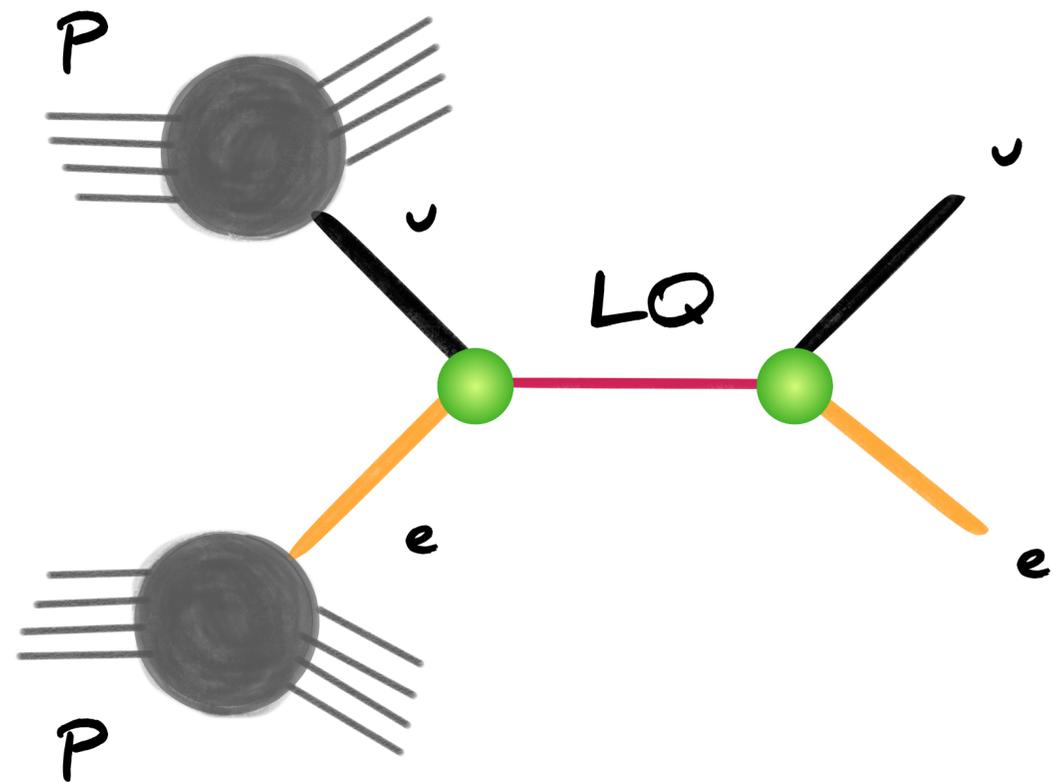
# Resonant LQ production @ the LHC



Suppressed by  $E_{T,miss}$  requirement

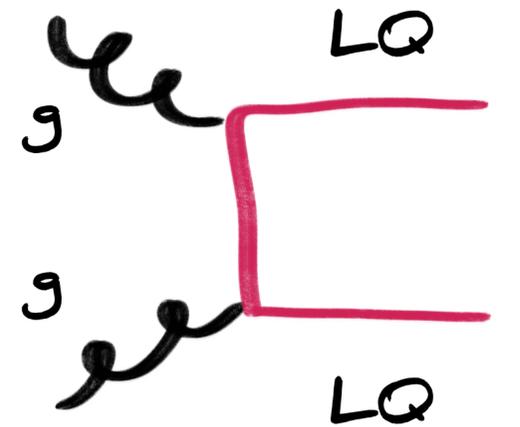


# Resonant LQ production @ the LHC

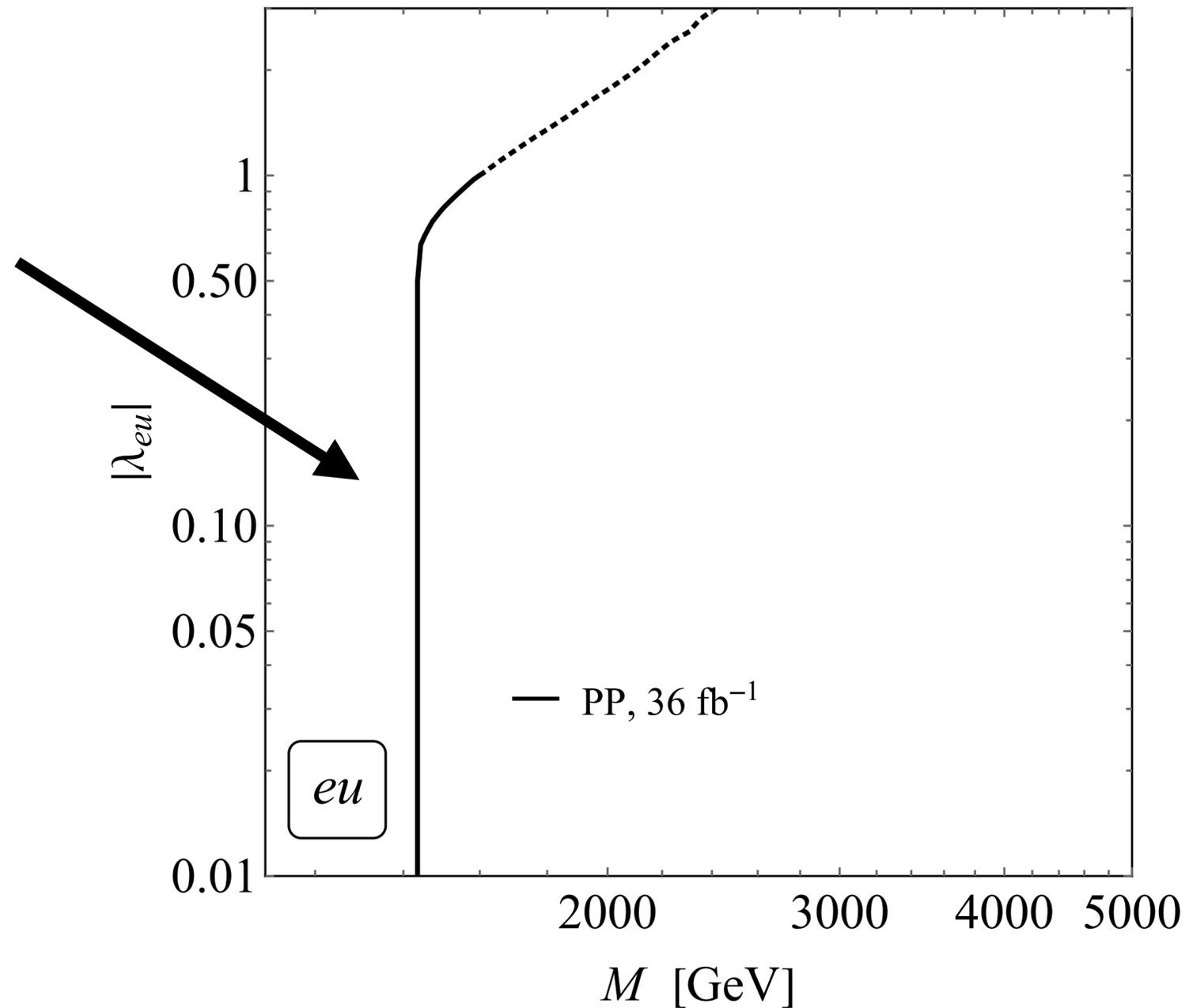


Sum over backgrounds is a steeply falling distribution, while signal exhibits a narrow peak

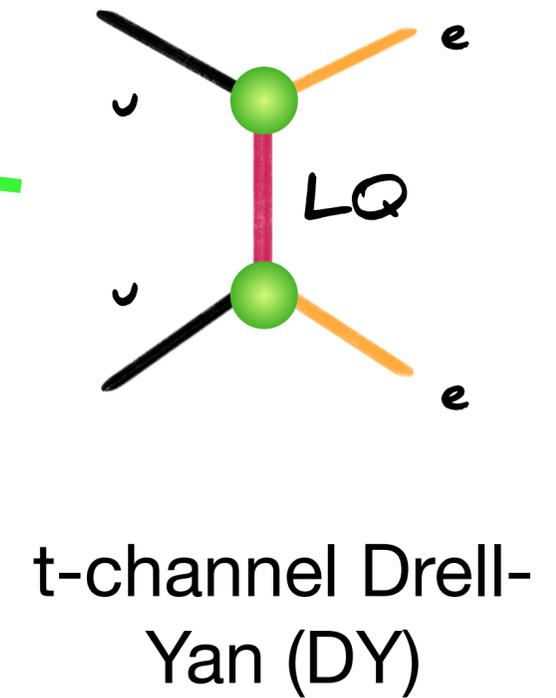
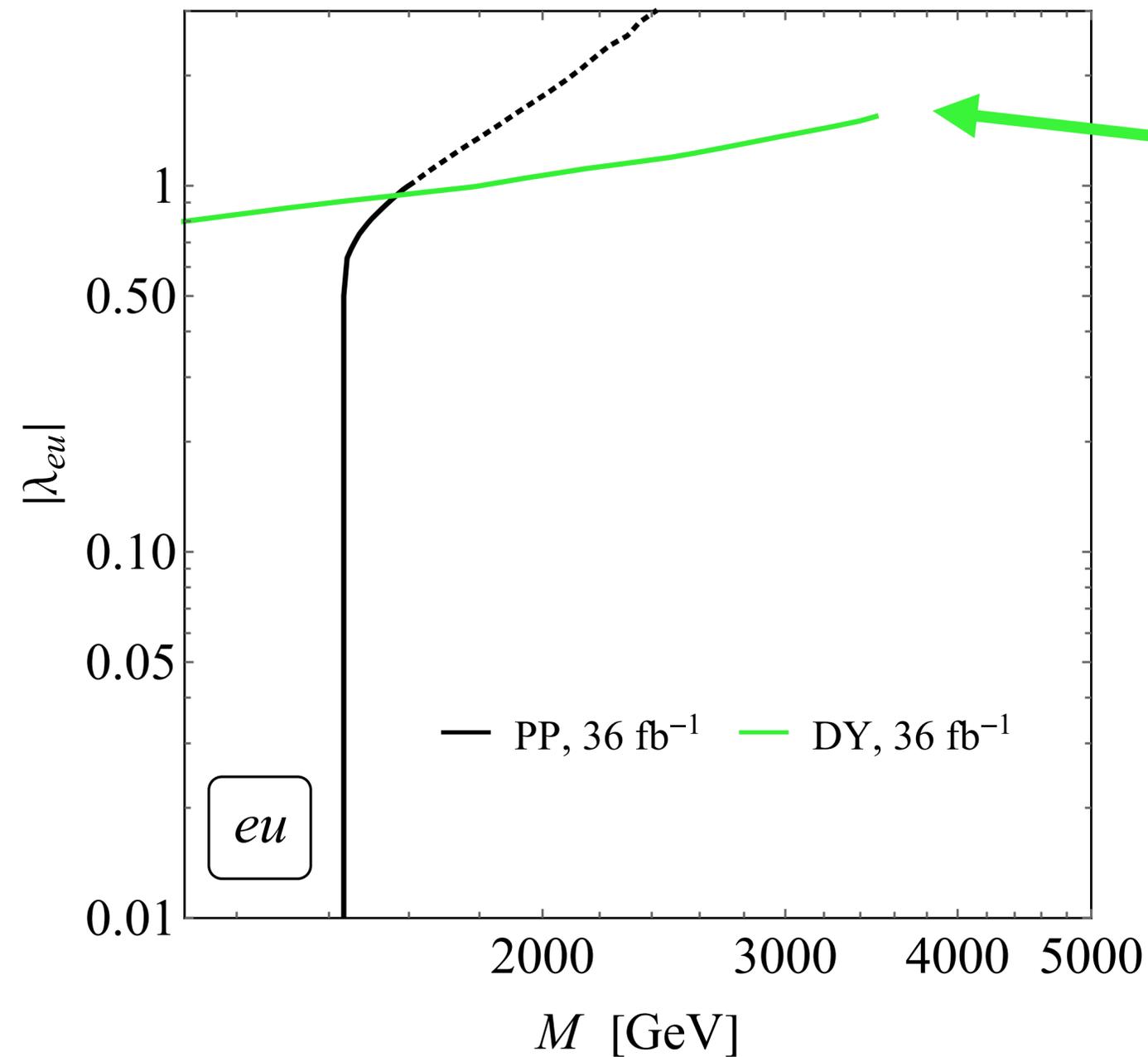
# LHC limits on 1<sup>st</sup> & 2<sup>nd</sup> generation LQs



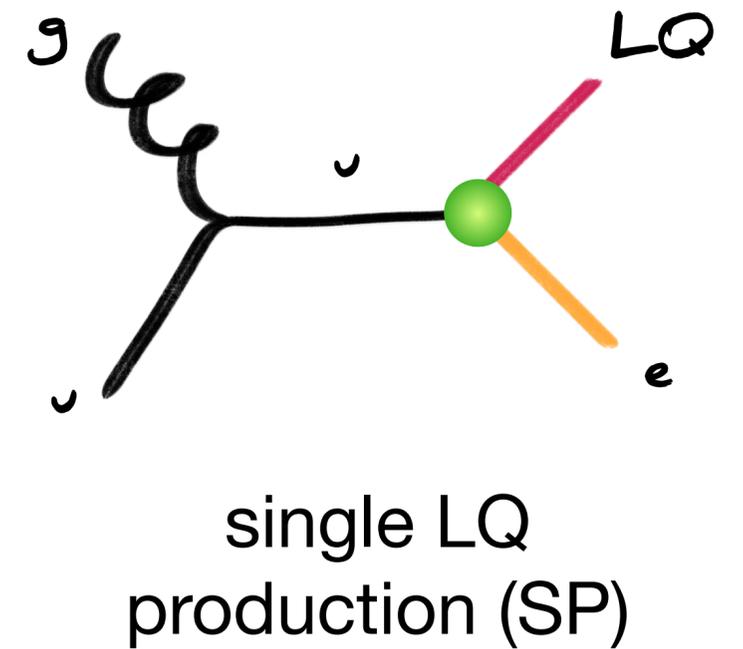
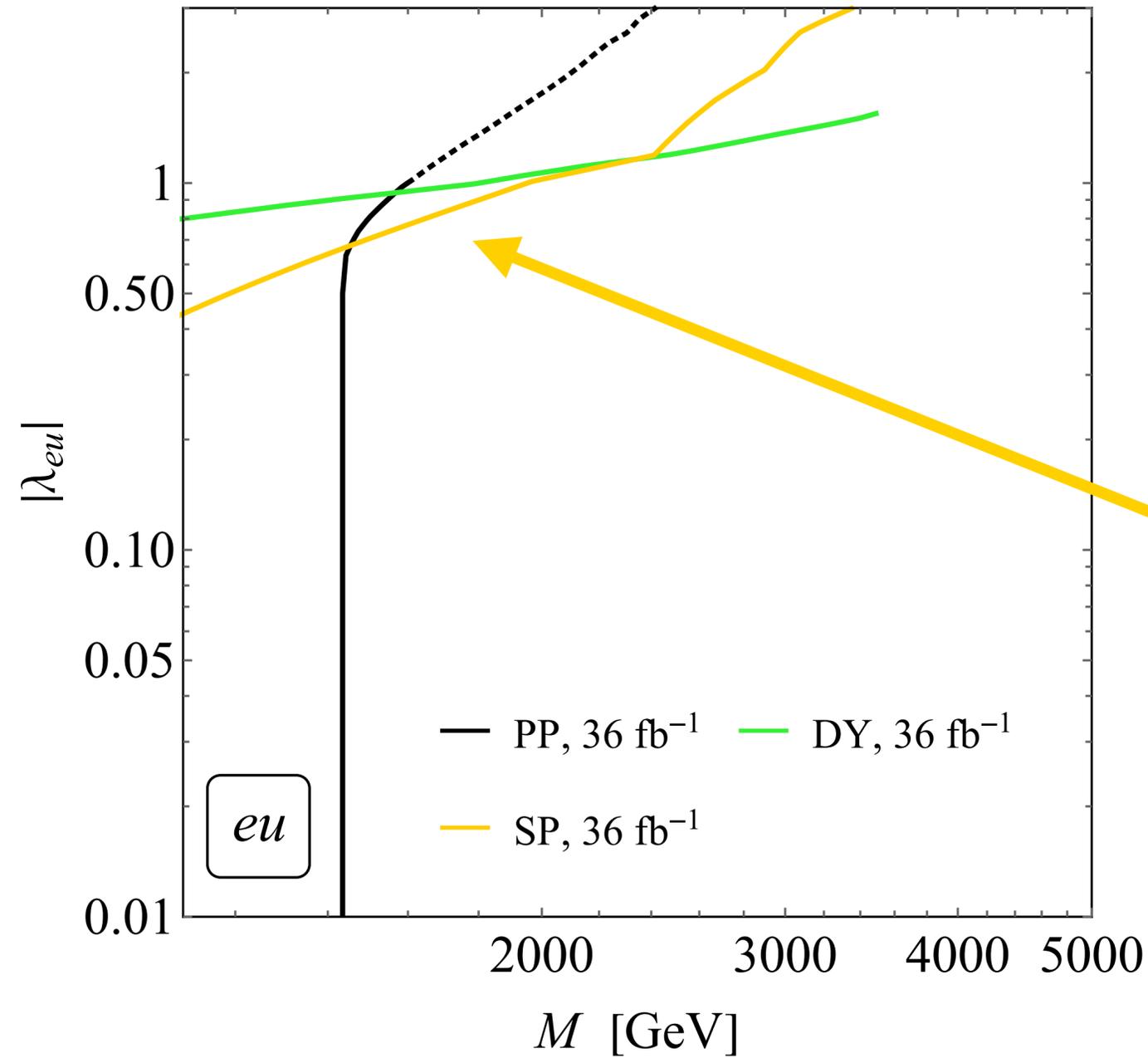
LQ pair  
production (PP)



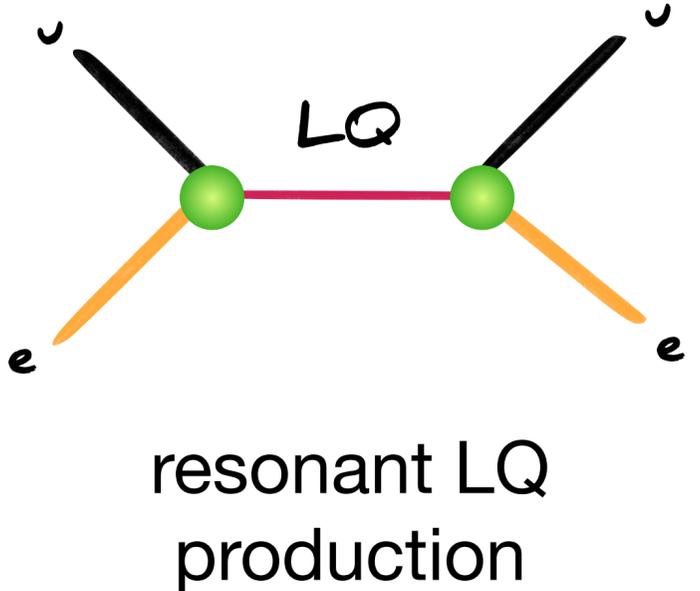
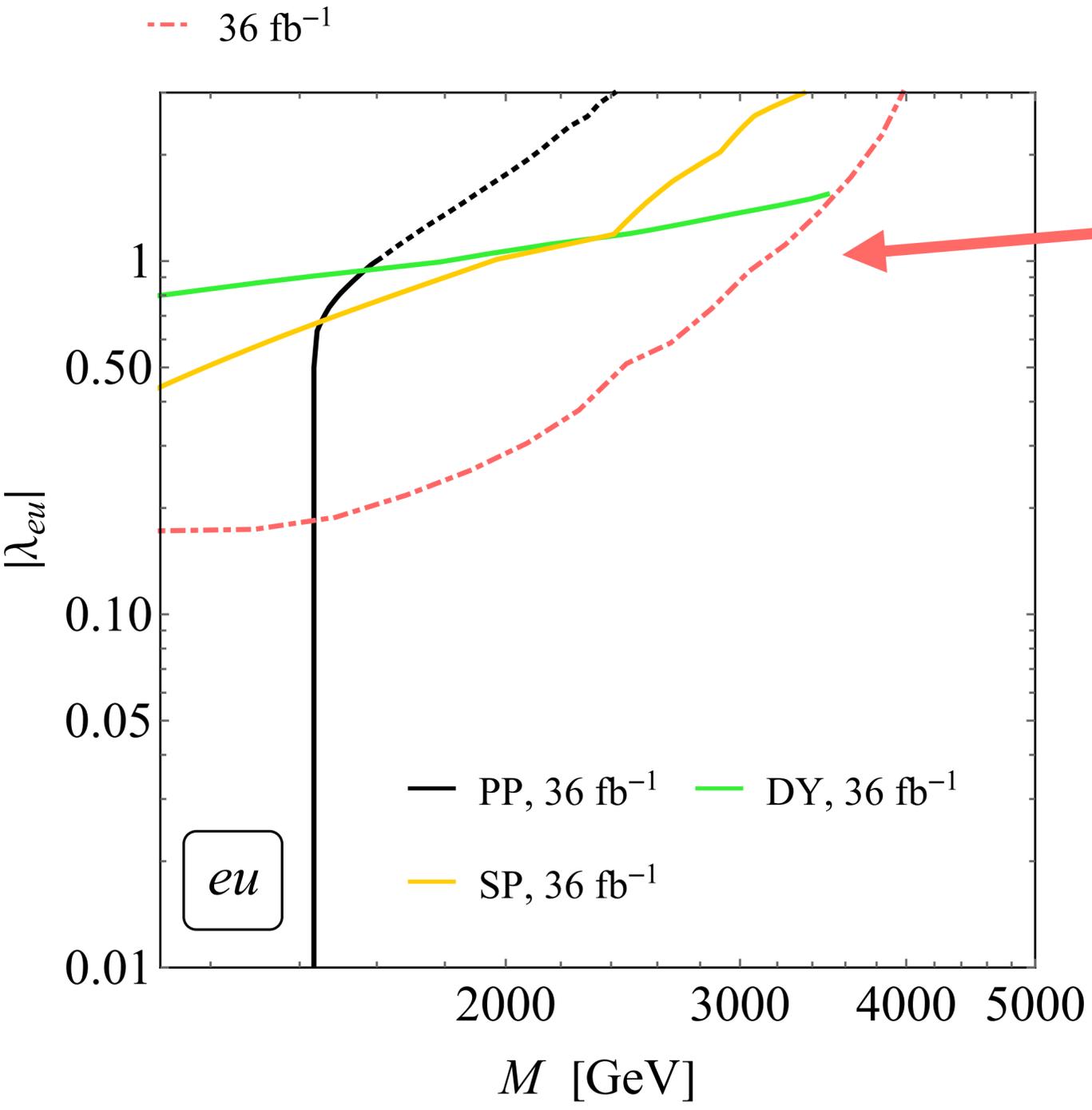
# LHC limits on 1<sup>st</sup> & 2<sup>nd</sup> generation LQs



# LHC limits on 1<sup>st</sup> & 2<sup>nd</sup> generation LQs

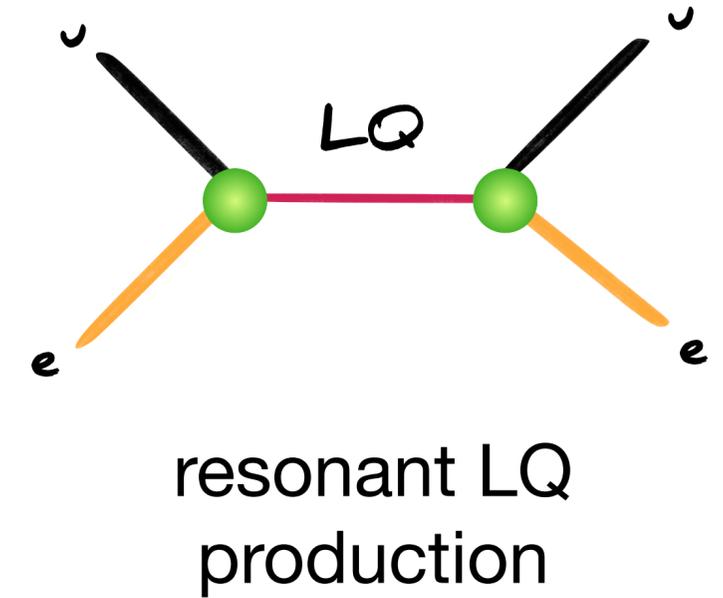
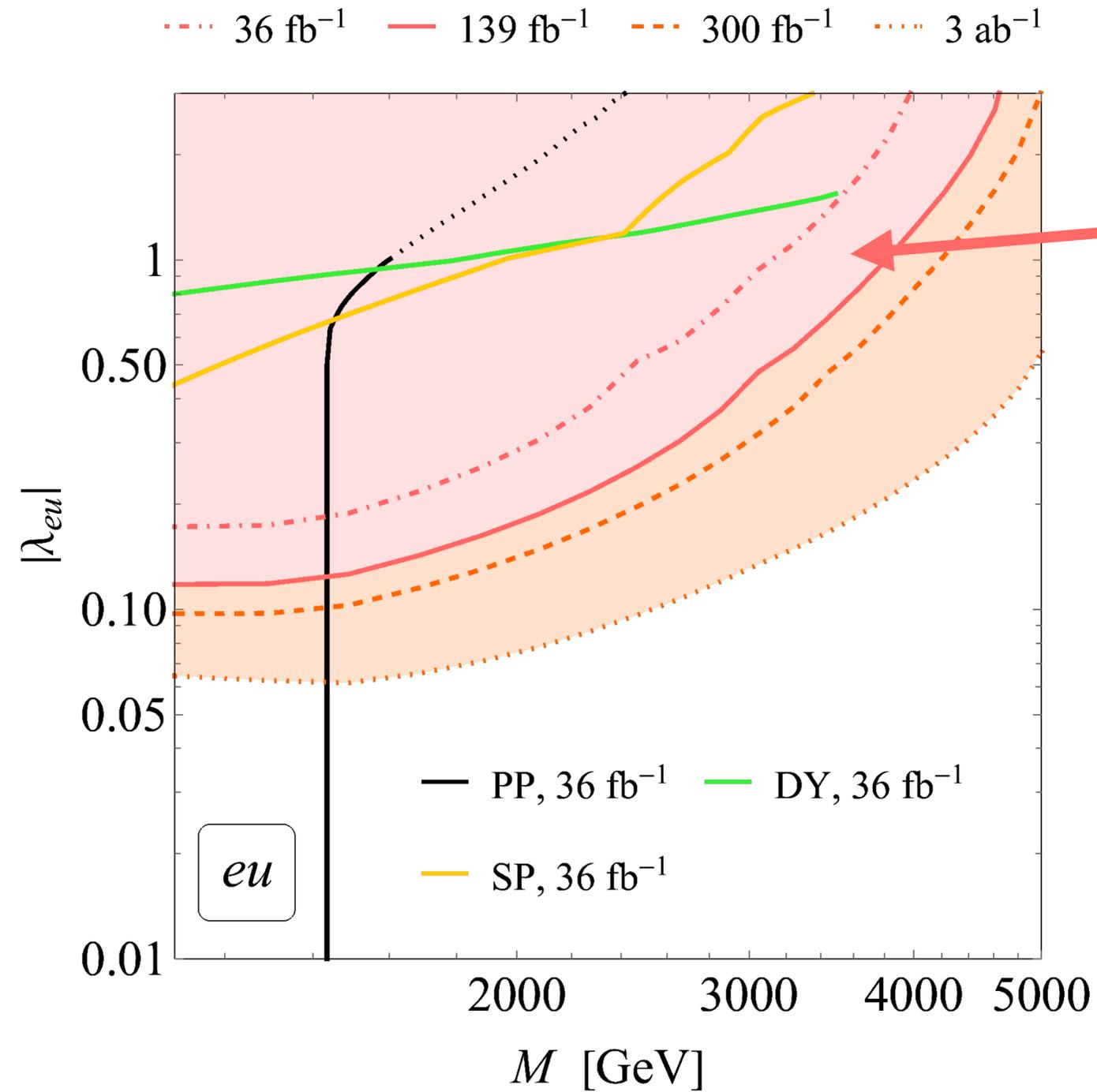


# LHC limits on 1<sup>st</sup> & 2<sup>nd</sup> generation LQs

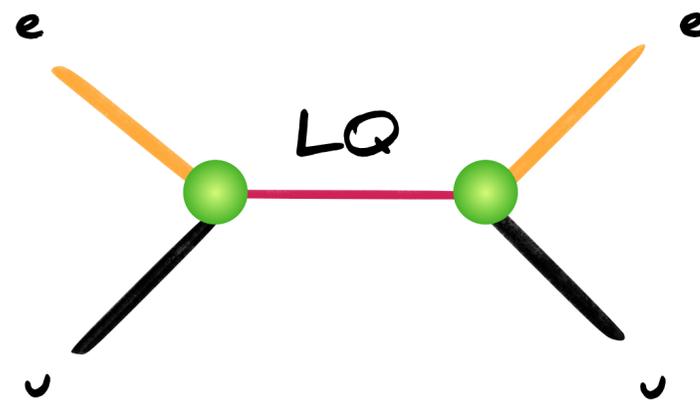


[Buonocore, UH, Nason, Tramontano & Zanderighi, PRL 125 (2020) 23]

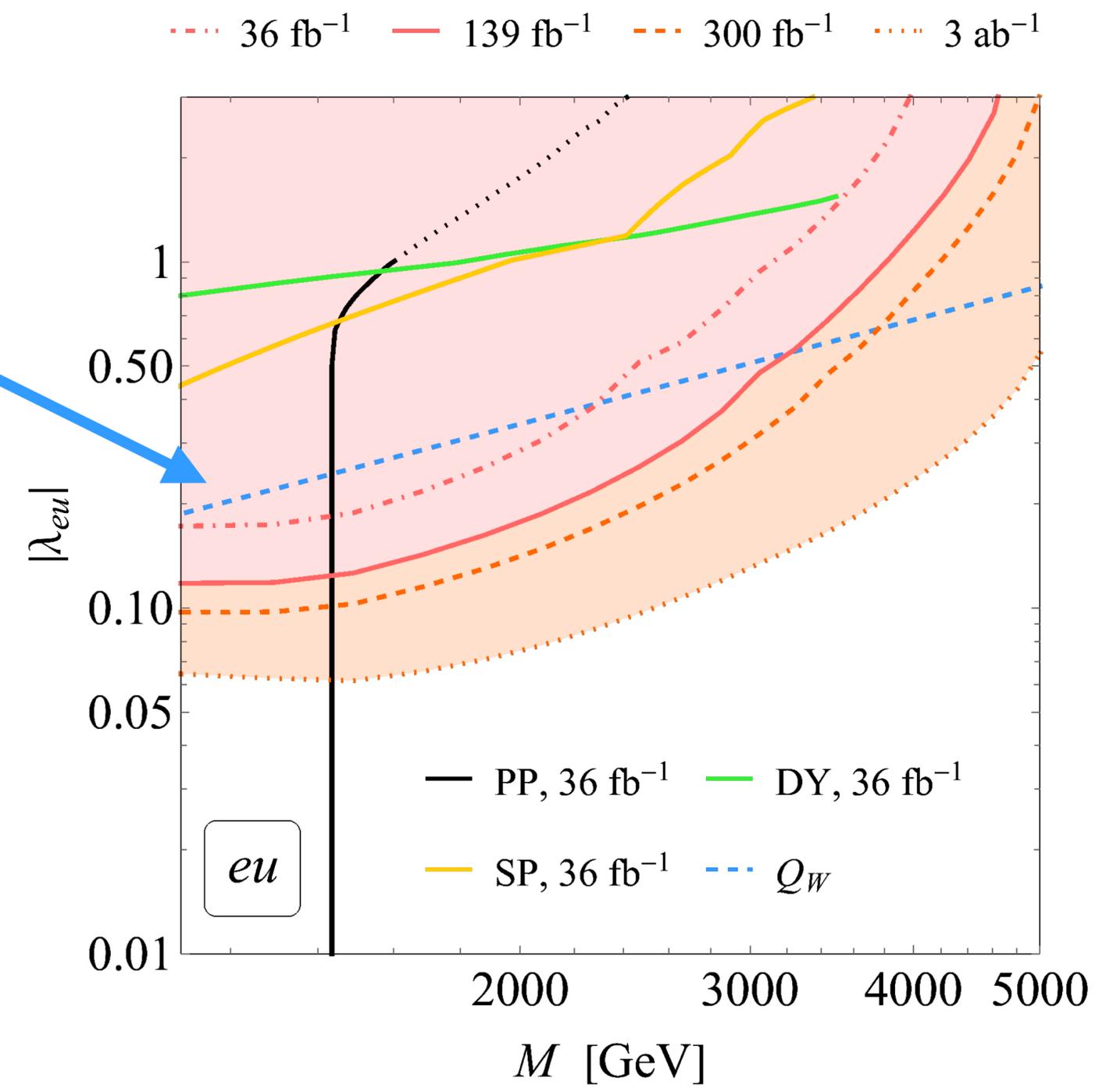
# LHC limits on 1<sup>st</sup> & 2<sup>nd</sup> generation LQs



# LHC limits on 1<sup>st</sup> & 2<sup>nd</sup> generation LQs

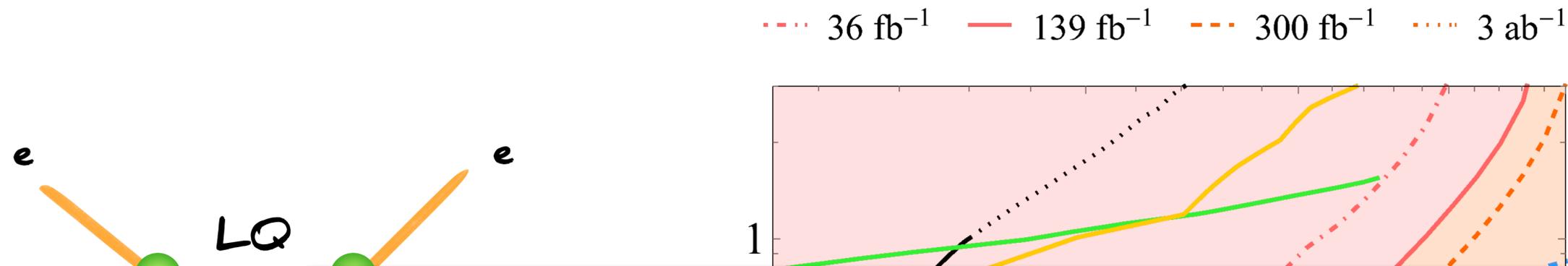


weak charge measurements ( $Q_W$ )



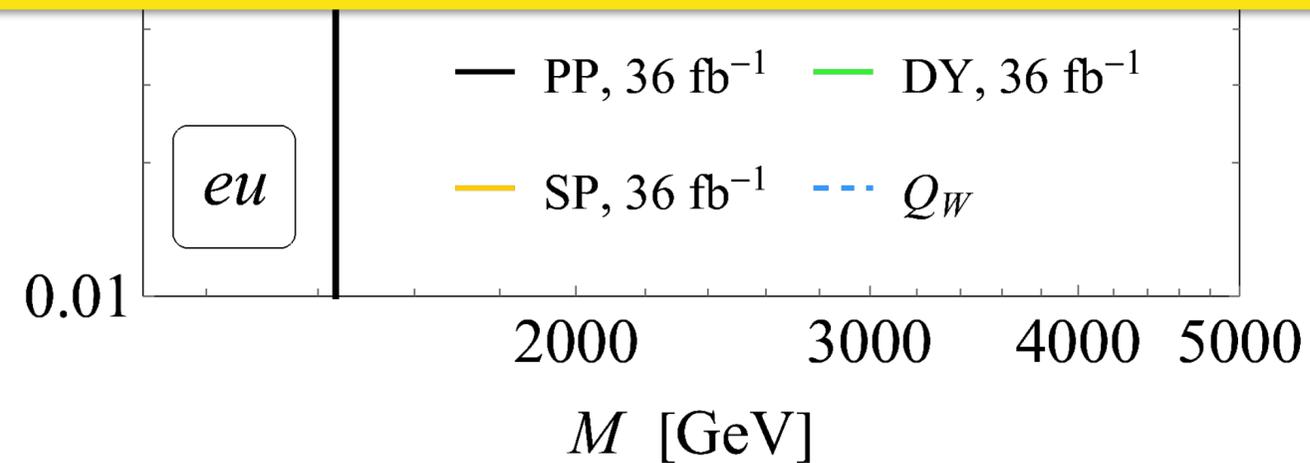
[Buonocore, UH, Nason, Tramontano & Zanderighi, PRL 125 (2020) 23]

# LHC limits on 1<sup>st</sup> & 2<sup>nd</sup> generation LQs



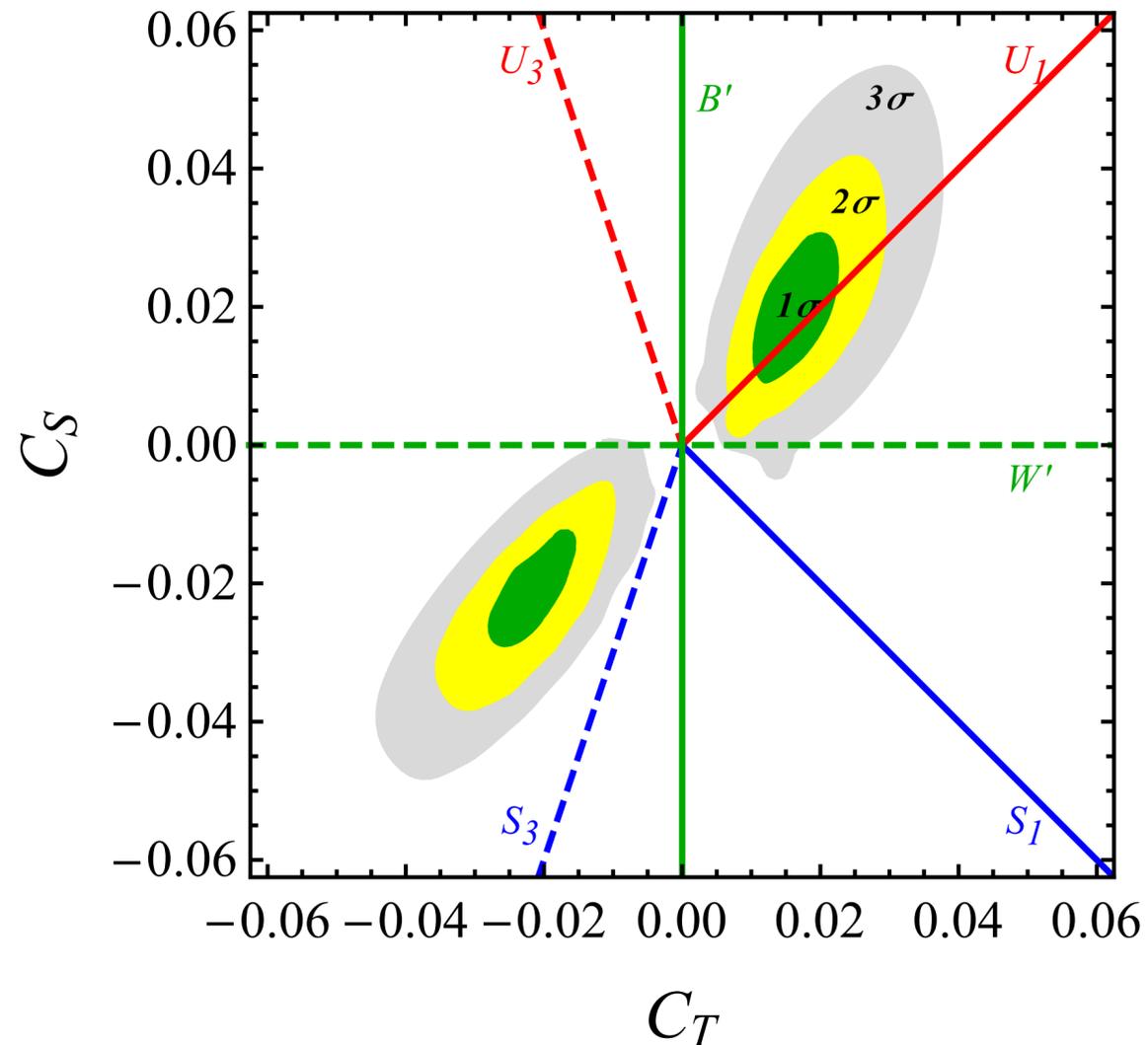
Given discovery reach of resonant LQ signature, dedicated searches for final states with a light lepton & a light-flavour jet should be added to exotics search canon of ATLAS & CMS

weak  
measurement



# Simplified models for B anomalies

$$\lambda_{ij}^q \lambda_{\alpha\beta}^l \left( C_T (\bar{Q}_L^i \gamma_\mu \sigma^a Q_L^j) (\bar{L}_L^\alpha \gamma^\mu \sigma^a L_L^\beta) + C_S (\bar{Q}_L^i \gamma_\mu Q_L^j) (\bar{L}_L^\alpha \gamma^\mu L_L^\beta) \right)$$

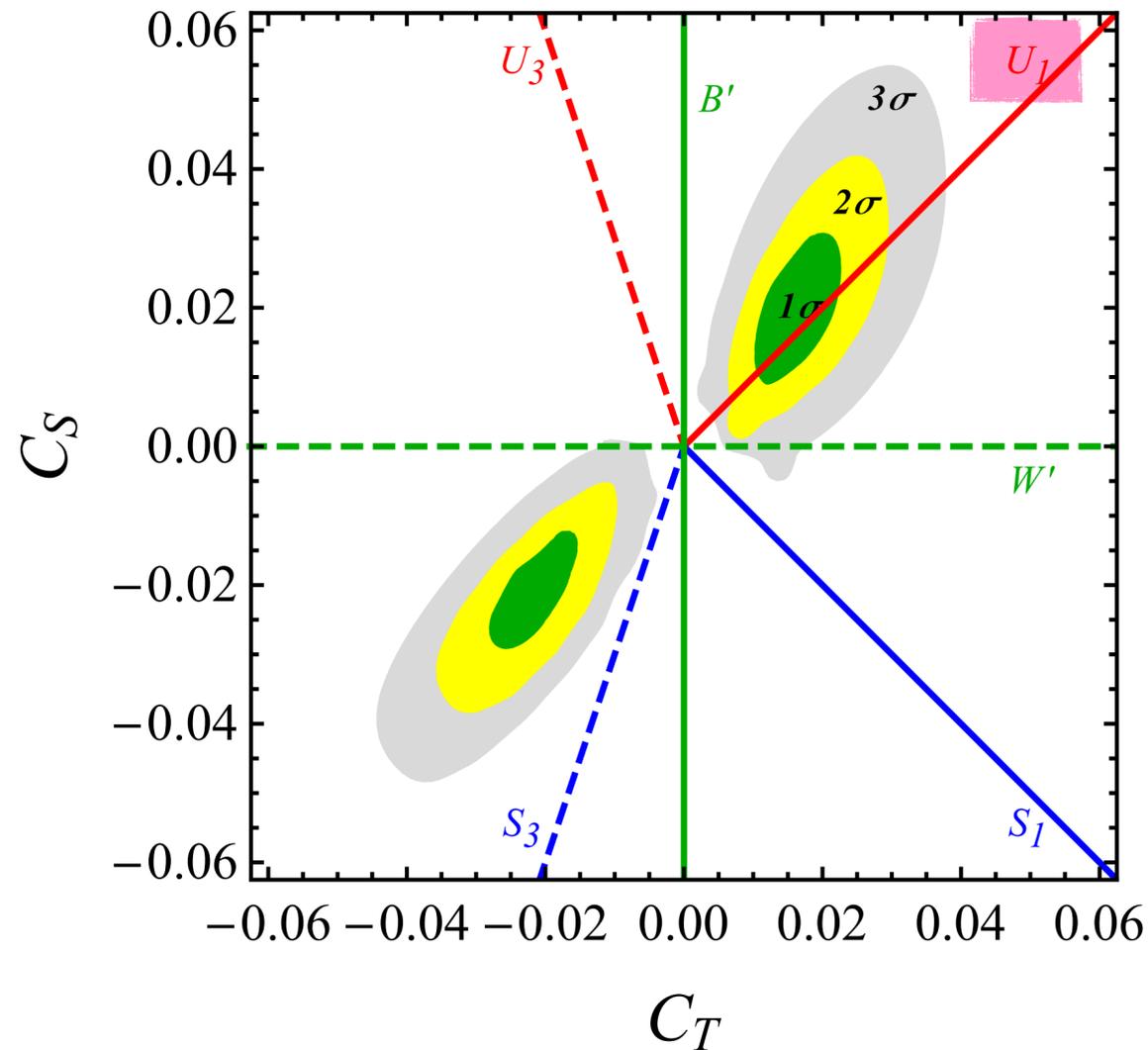


Model	Mediator	$b \rightarrow s$	$b \rightarrow c$
Colorless vectors	$B' = (1, 1, 0)$	✓	✗
	$W' = (1, 3, 0)$	✗	✓
Scalar leptoquarks	$S_1 = (\bar{3}, 1, 1/3)$	✗	✓
	$S_3 = (\bar{3}, 3, 1/3)$	✓	✗
Vector leptoquarks	$U_1 = (3, 1, 2/3)$	✓	✓
	$U_3 = (3, 3, 2/3)$	✓	✗

$b \rightarrow s$  ( $b \rightarrow c$ ) anomalies alone can be accommodated by several simple single-mediator models

# Simplified models for B anomalies

$$\lambda_{ij}^q \lambda_{\alpha\beta}^l \left( C_T (\bar{Q}_L^i \gamma_\mu \sigma^a Q_L^j) (\bar{L}_L^\alpha \gamma^\mu \sigma^a L_L^\beta) + C_S (\bar{Q}_L^i \gamma_\mu Q_L^j) (\bar{L}_L^\alpha \gamma^\mu L_L^\beta) \right)$$



Model	Mediator	$b \rightarrow s$	$b \rightarrow c$
Colorless vectors	$B' = (1, 1, 0)$	✓	✗
	$W' = (1, 3, 0)$	✗	✓
Scalar leptoquarks	$S_1 = (\bar{3}, 1, 1/3)$	✗	✓
	$S_3 = (\bar{3}, 3, 1/3)$	✓	✗
Vector leptoquarks	$U_1 = (3, 1, 2/3)$	✓	✓
	$U_3 = (3, 3, 2/3)$	✓	✗

$U_1$  singlet vector LQ is the only single-mediator model that can explain both sets of anomalies

# Singlet vector LQ models for B anomalies

$$\mathcal{L} \supset \frac{g_U}{\sqrt{2}} \left[ \beta_L^{ij} \bar{Q}_L^{i,a} \gamma_\mu L_L^j + \beta_R^{ij} \bar{d}_R^{i,a} \gamma_\mu \ell_R^j \right] U^{\mu,a} + \text{h.c.}, \quad |\beta_L^{22}| \lesssim |\beta_L^{32}| \ll |\beta_L^{23}| \lesssim |\beta_L^{33}| = \mathcal{O}(1)$$

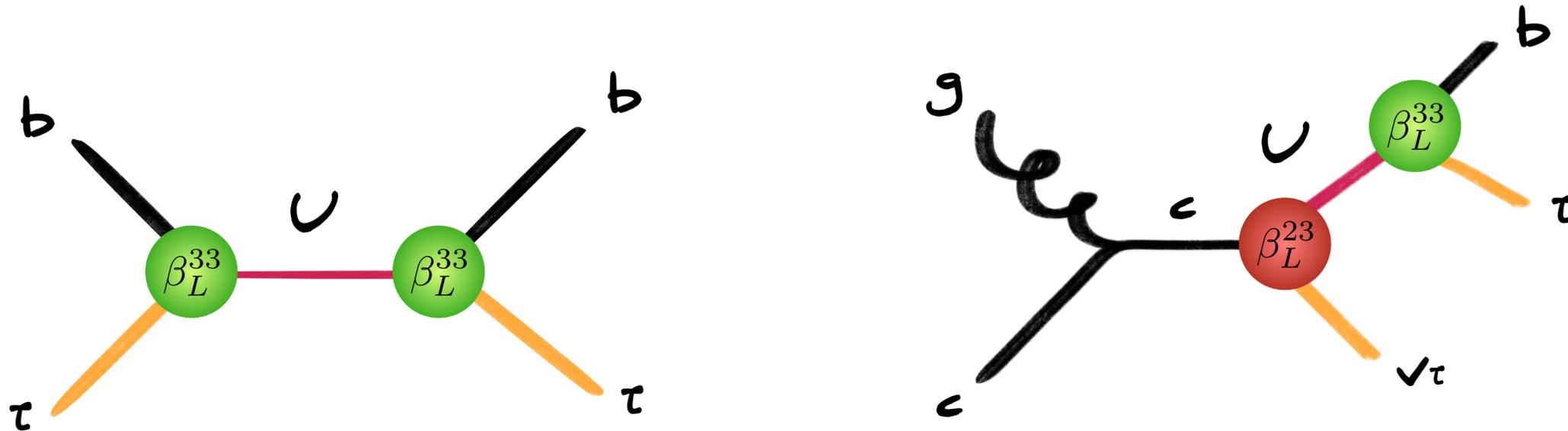
Parameters		Branching ratios			
$\beta_L^{33}$	$\beta_L^{23}$	BR ( $U \rightarrow b\tau^+$ )	BR ( $U \rightarrow t\bar{\nu}_\tau$ )	BR ( $U \rightarrow s\tau^+$ )	BR ( $U \rightarrow c\bar{\nu}_\tau$ )
1	0	51%	49%	0%	0%
1	1	25%	22%	25%	27%

b +  $\tau$   
signature

mono-top  
signature

mono-jet  
signature

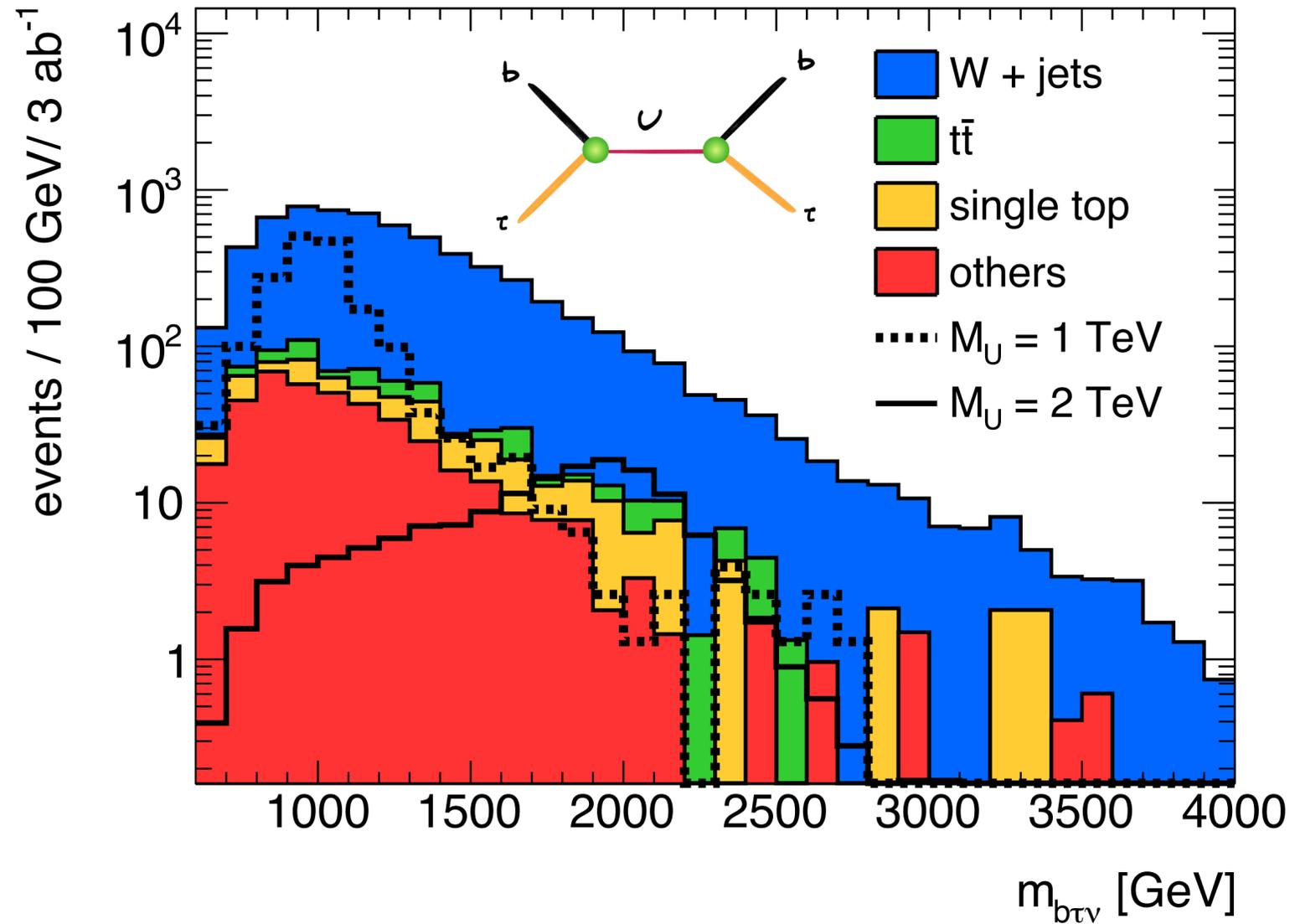
# LQ contributions to $b + \tau$ signature



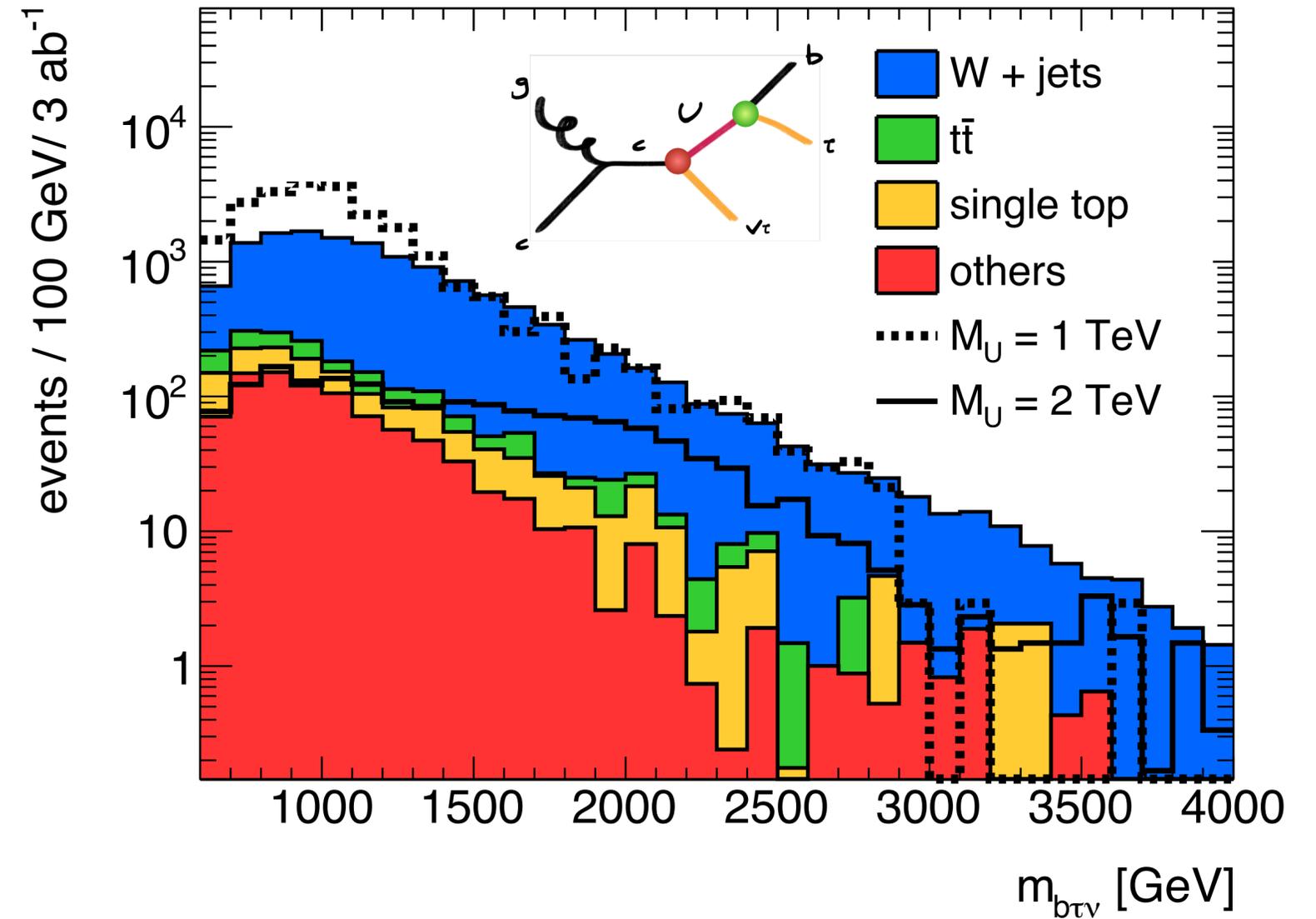
For  $\beta_L^{23} = 0$ ,  $b + \tau$  signal arises only from  $2 \rightarrow 2$  process, while for  $\beta_L^{23} \neq 0$  also  $2 \rightarrow 3$  scattering is relevant. Since two topologies lead to final states with very different kinematic features, it is essential to develop two separate search strategies for them

# Kinematic distributions of $b + \tau$ signal

LHC 14 TeV,  $b + \tau$

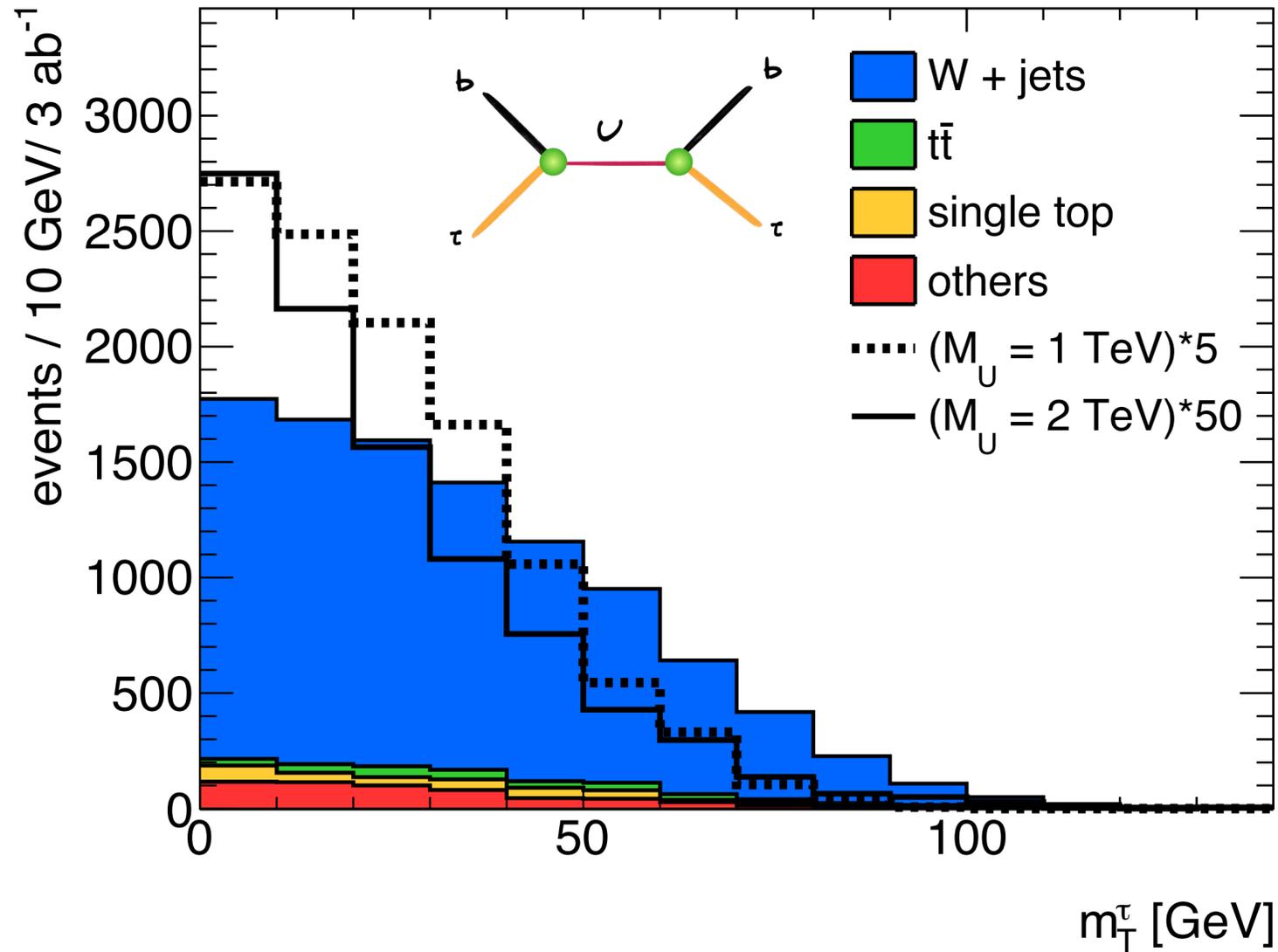


LHC 14 TeV,  $b + \tau$

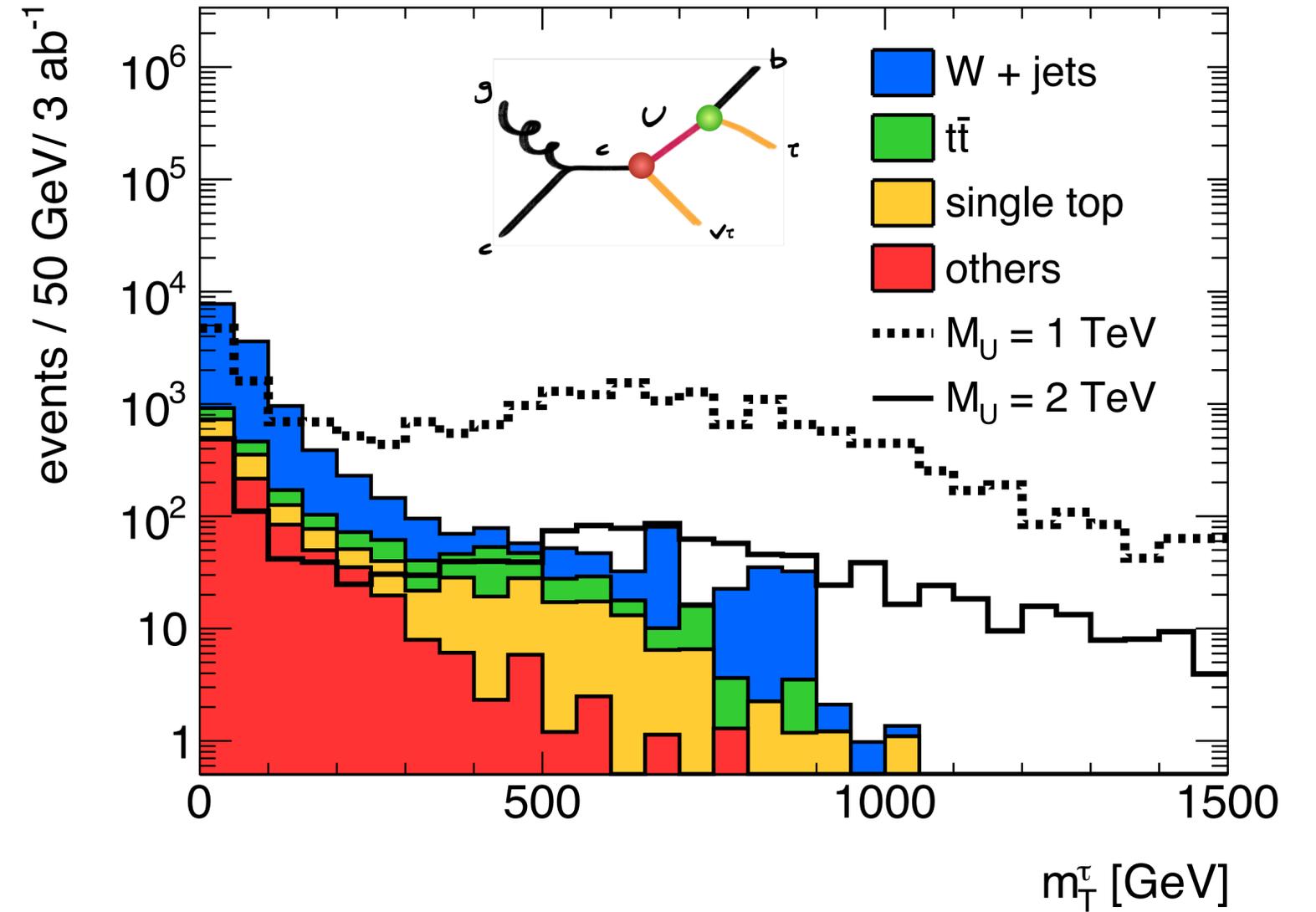


# Kinematic distributions of $b + \tau$ signal

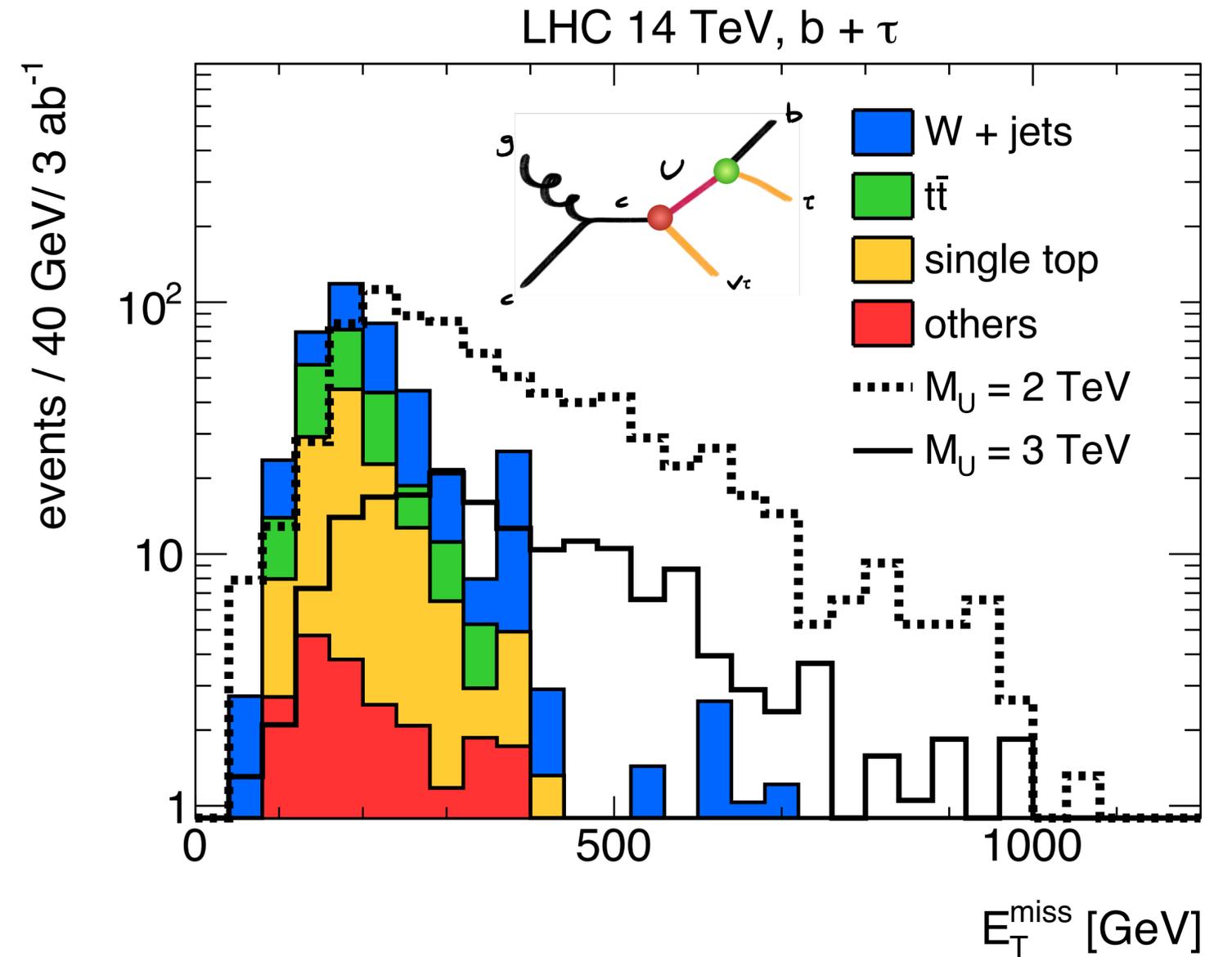
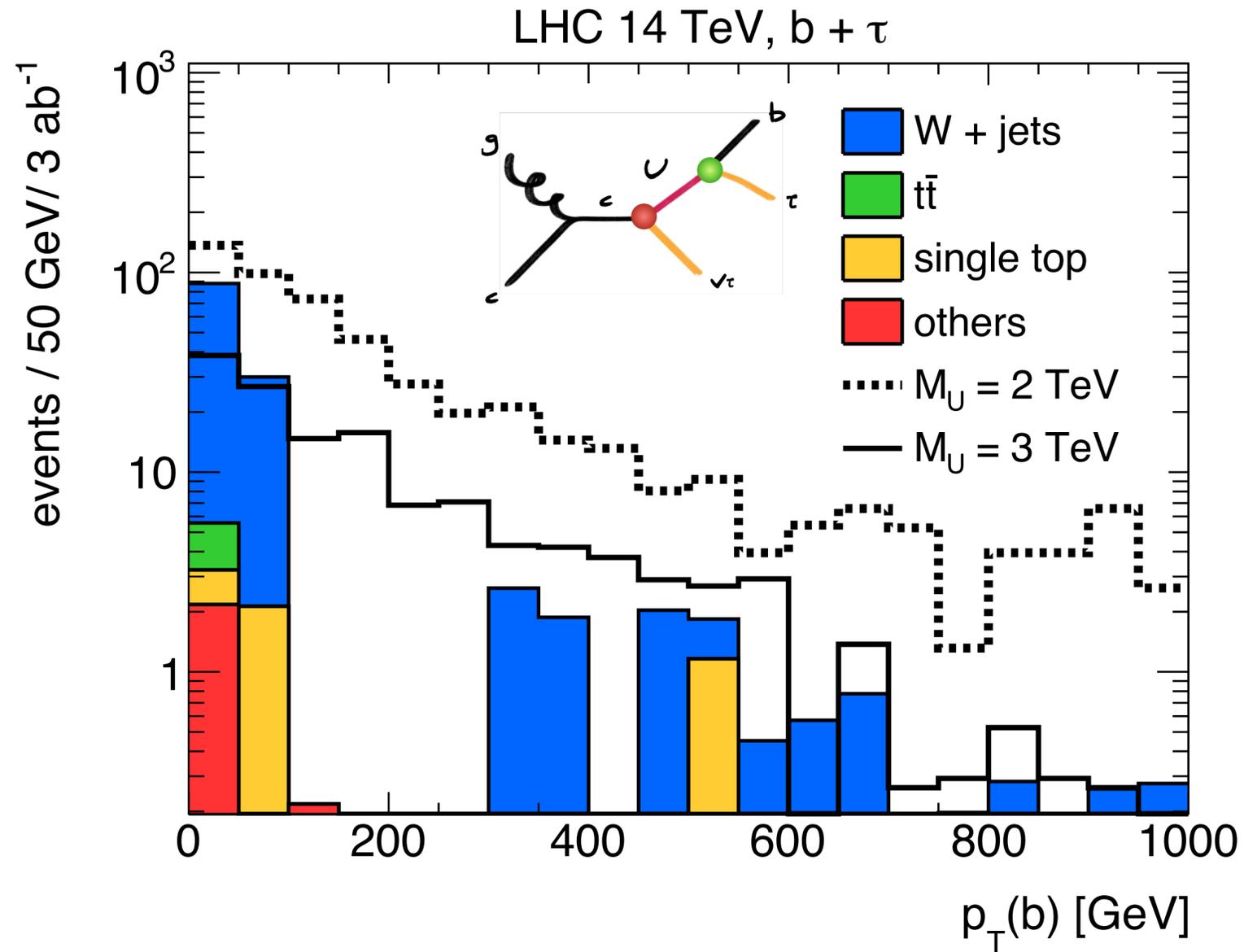
LHC 14 TeV,  $b + \tau$



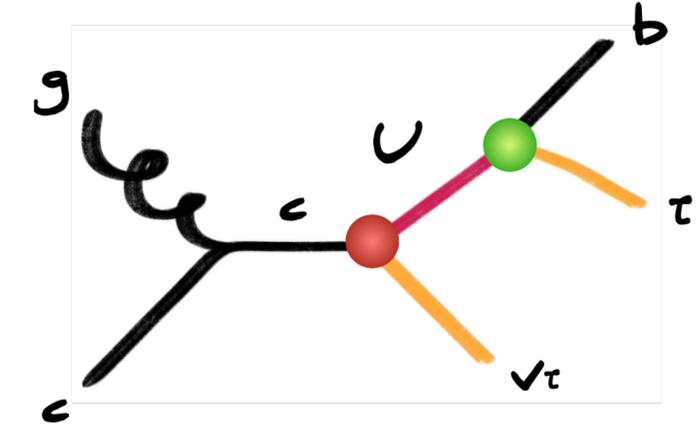
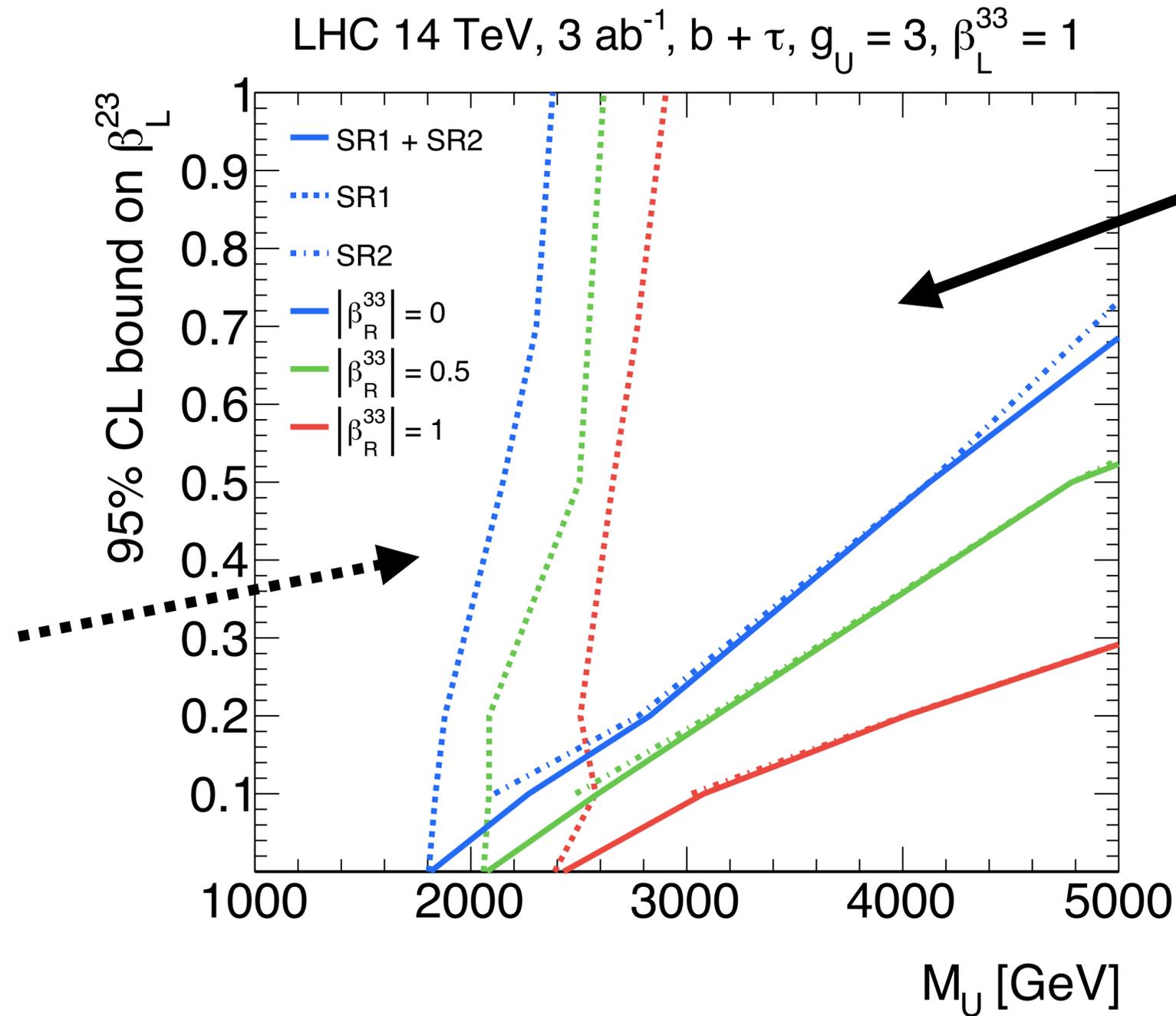
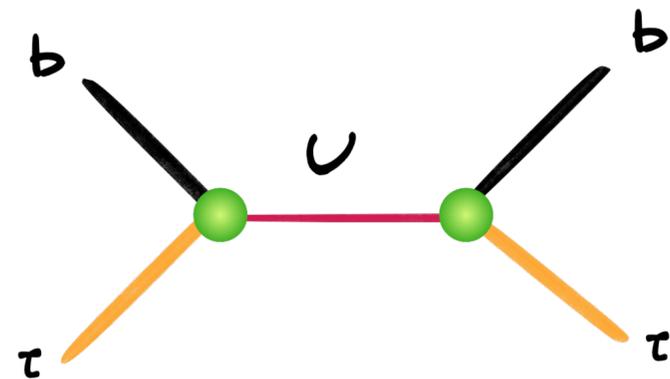
LHC 14 TeV,  $b + \tau$



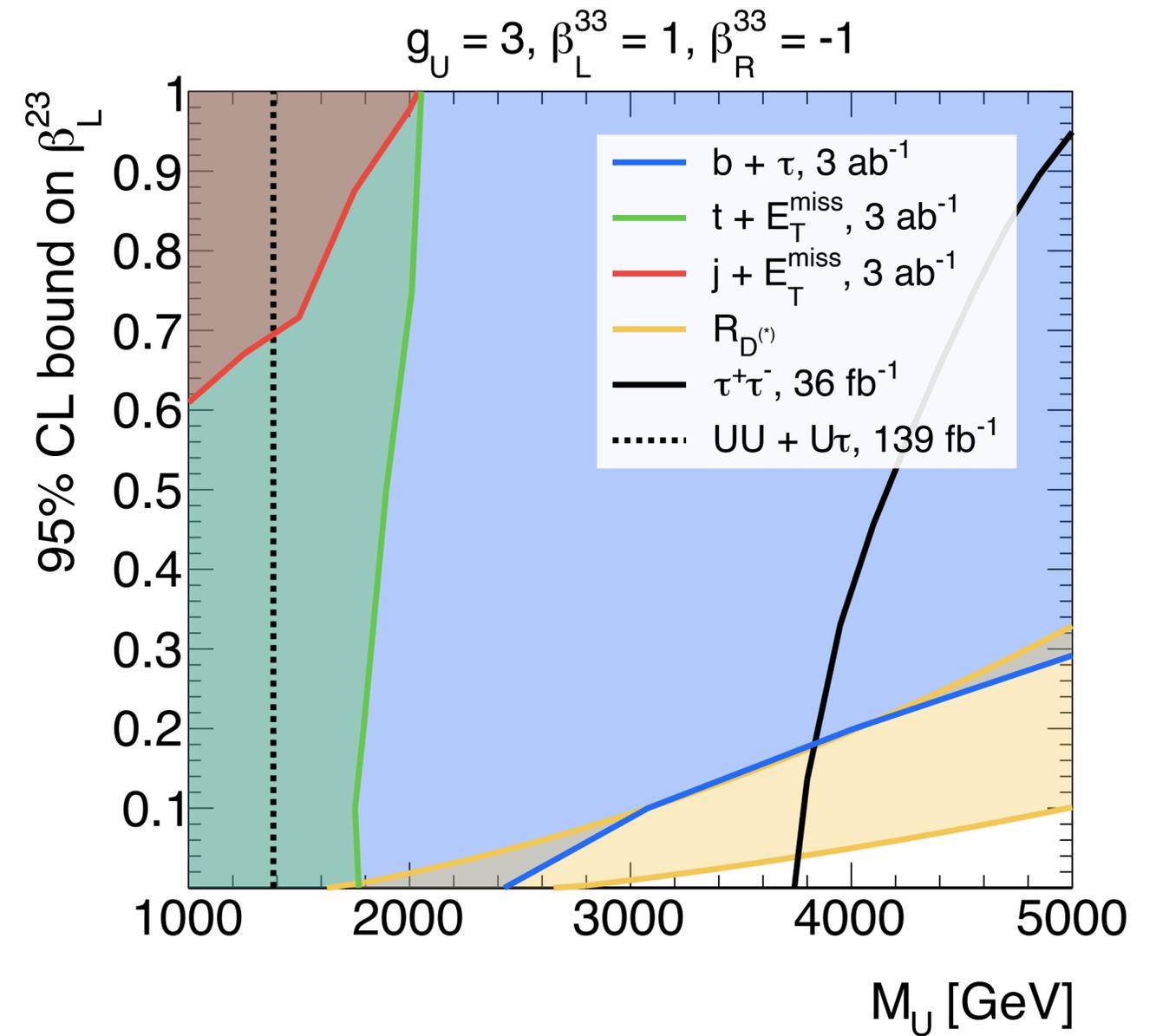
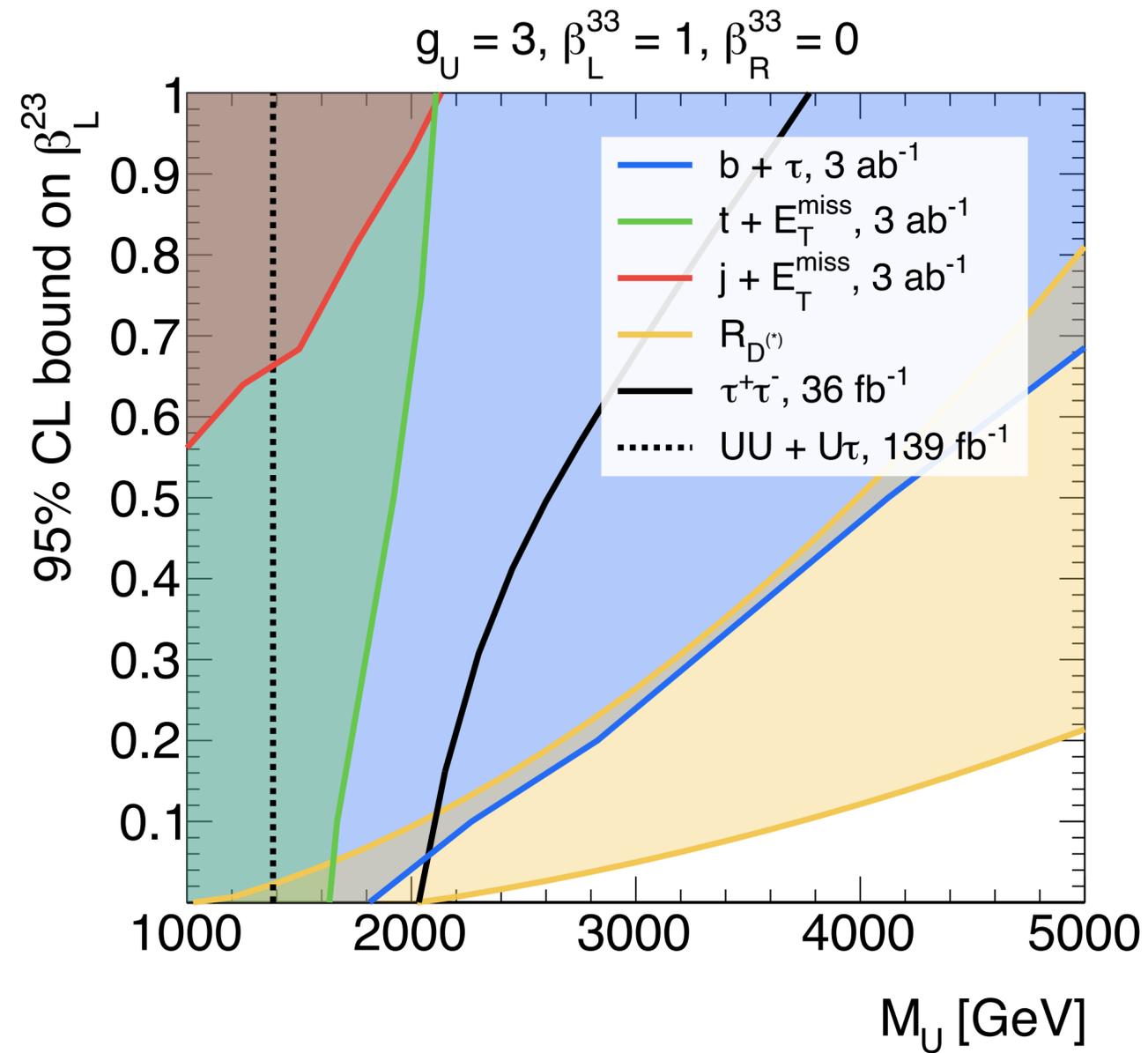
# Kinematic distributions of $b + \tau$ signal



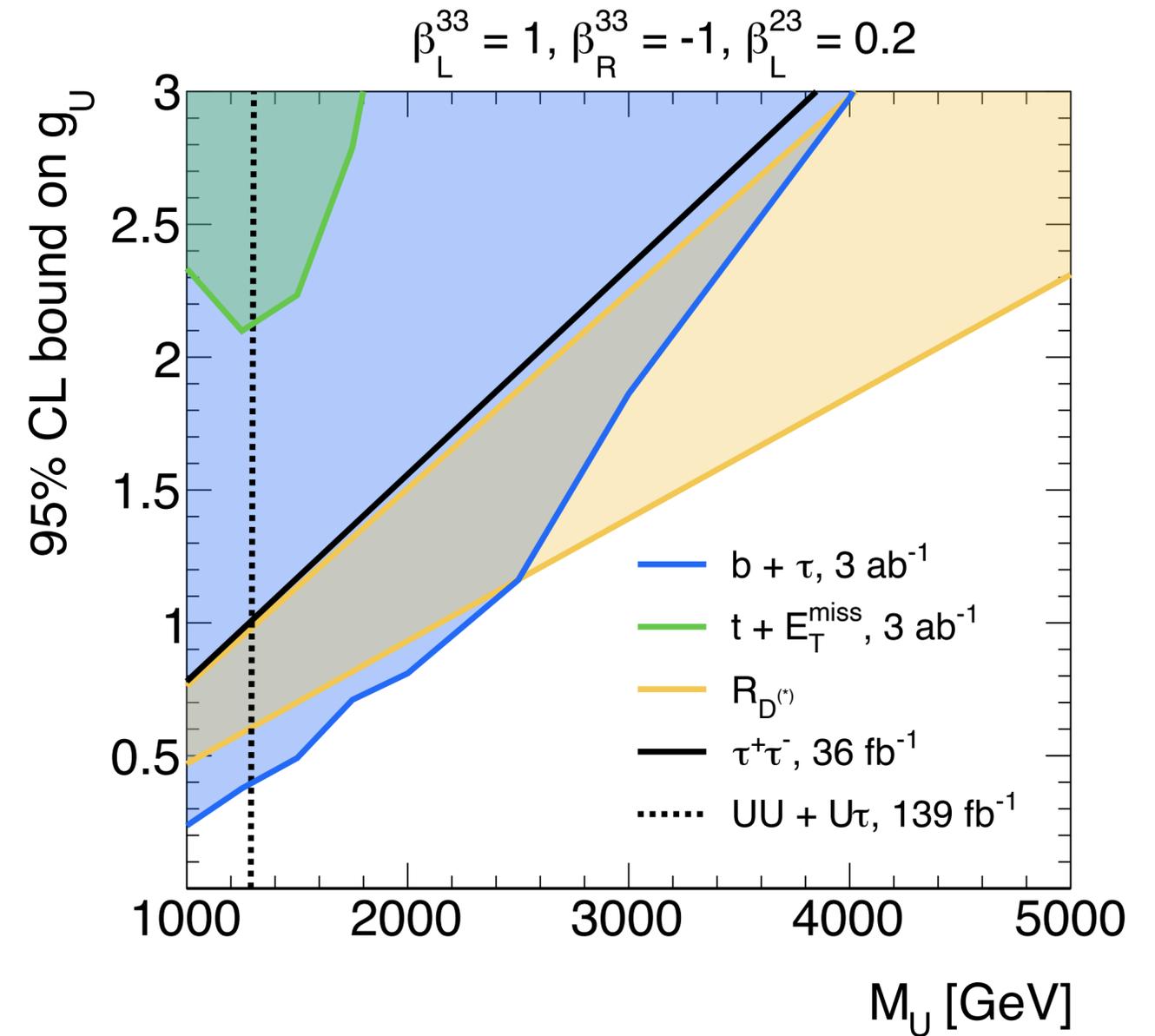
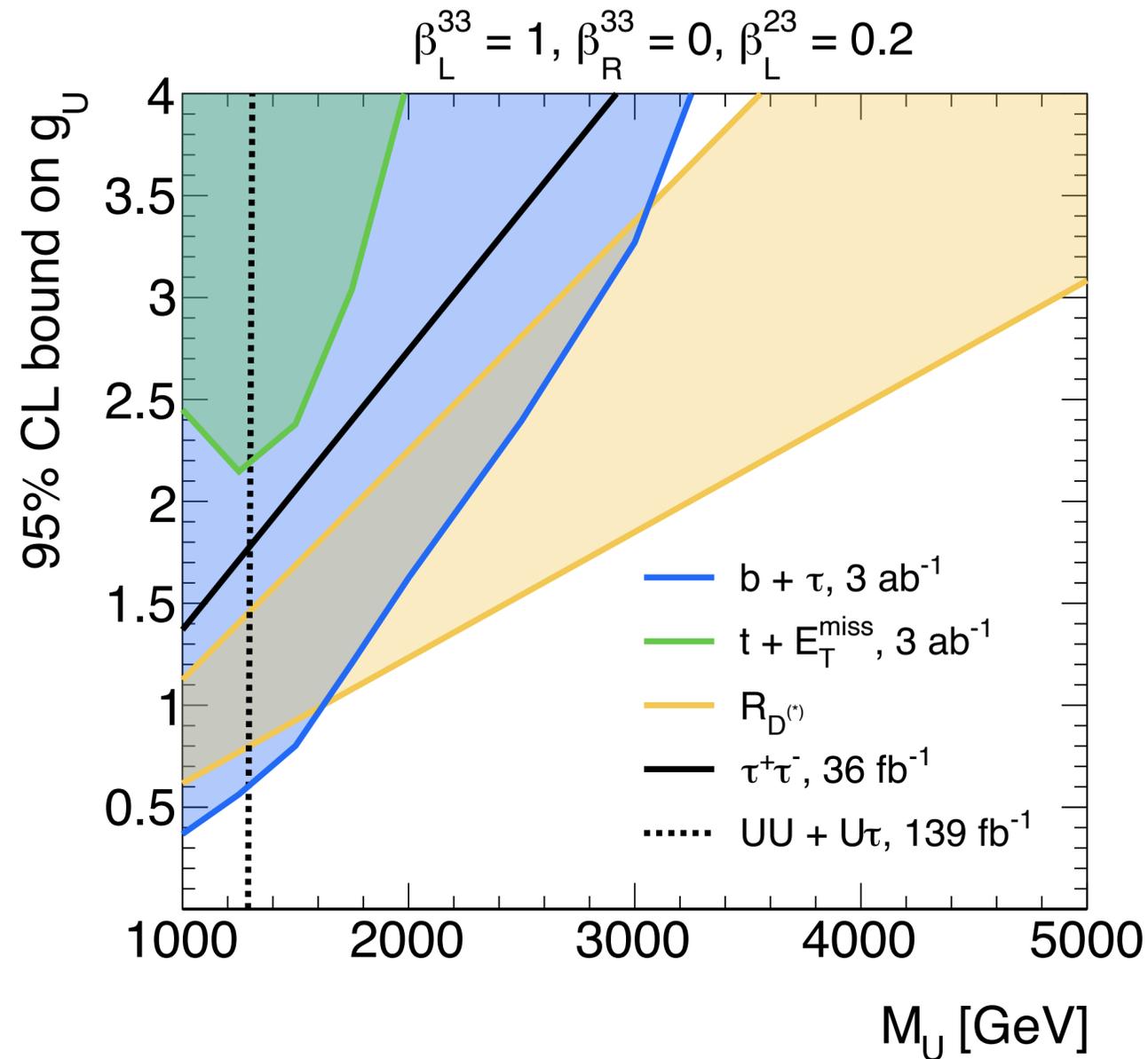
# b + $\tau$ constraints from $2 \rightarrow 2$ & $2 \rightarrow 3$ signal



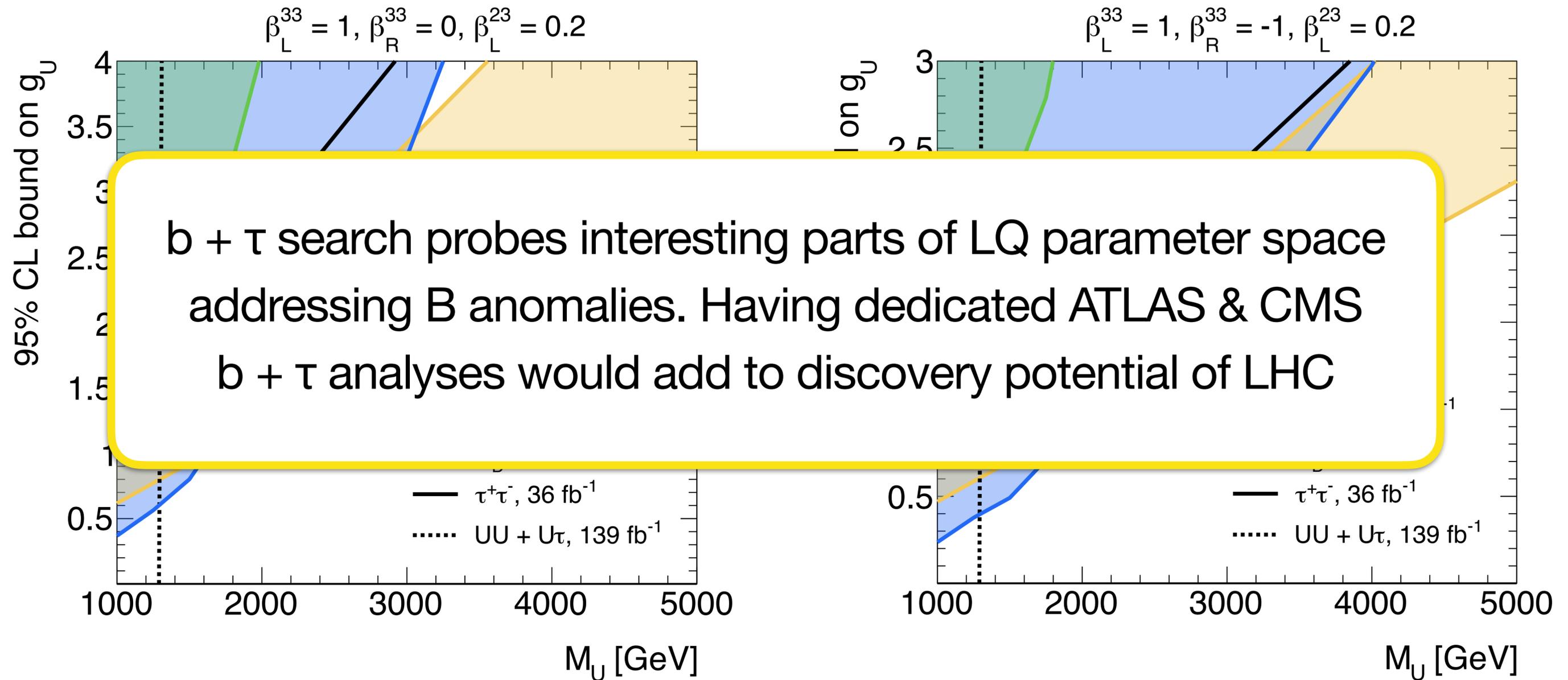
# Constraints from new LQ search strategies



# Constraints from new LQ search strategies



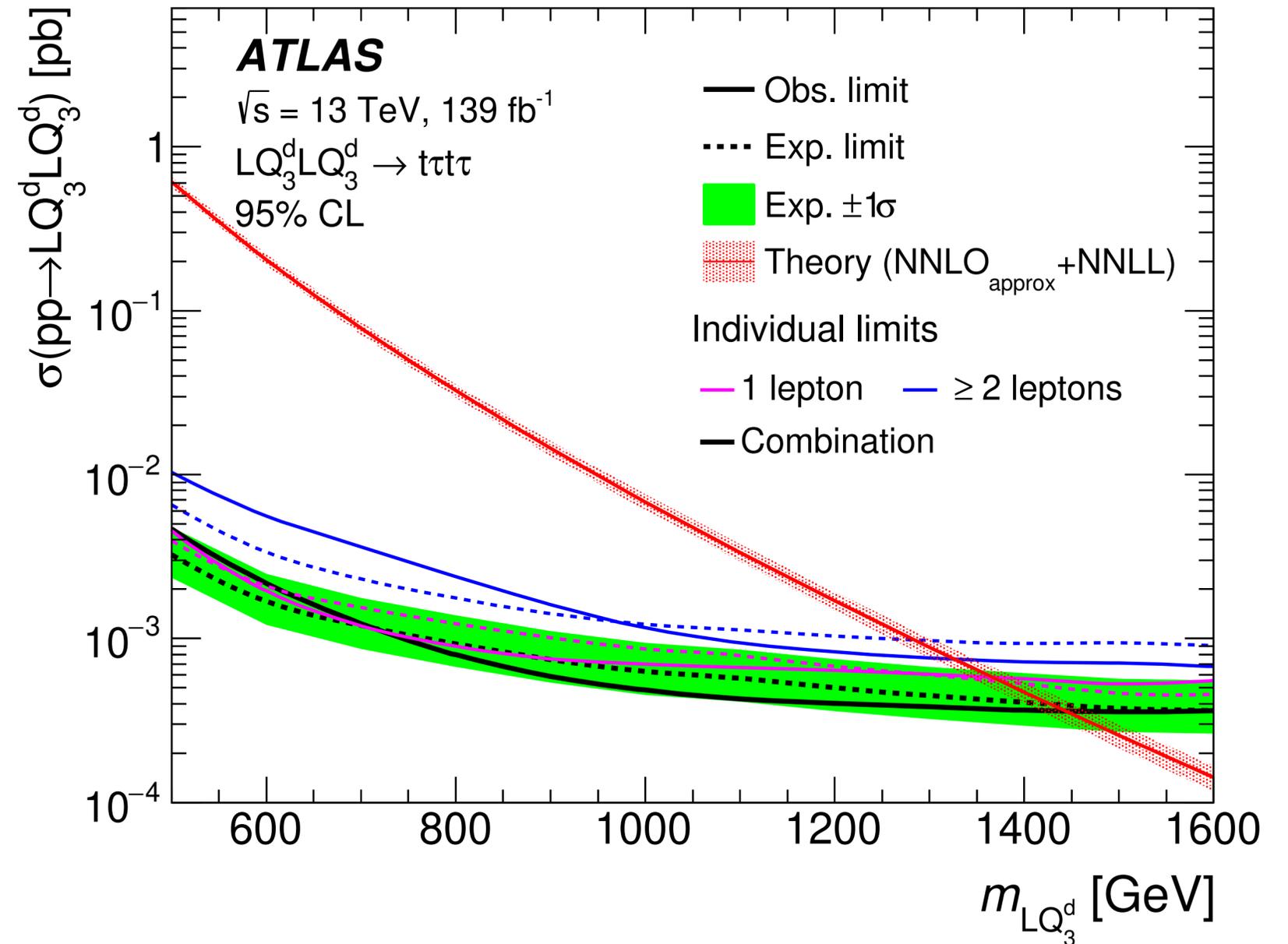
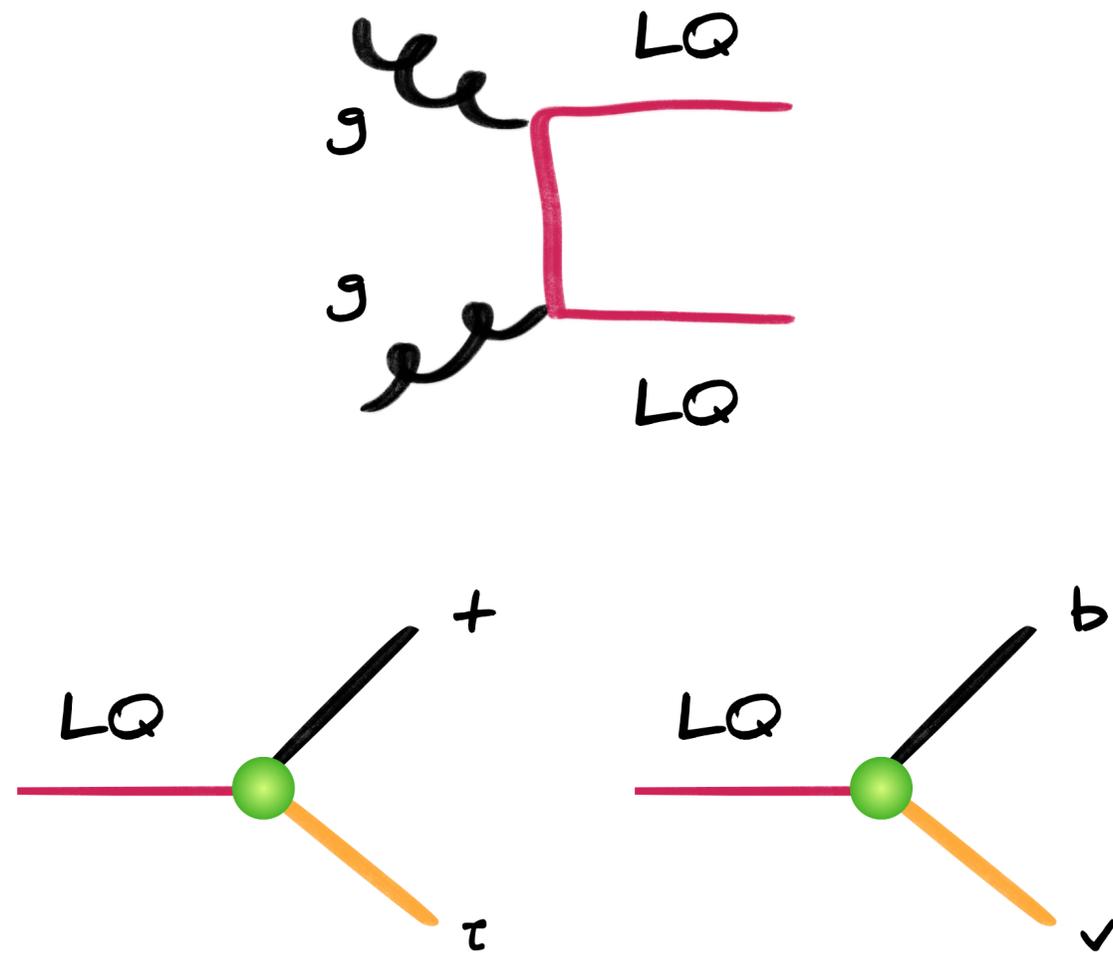
# Constraints from new LQ search strategies



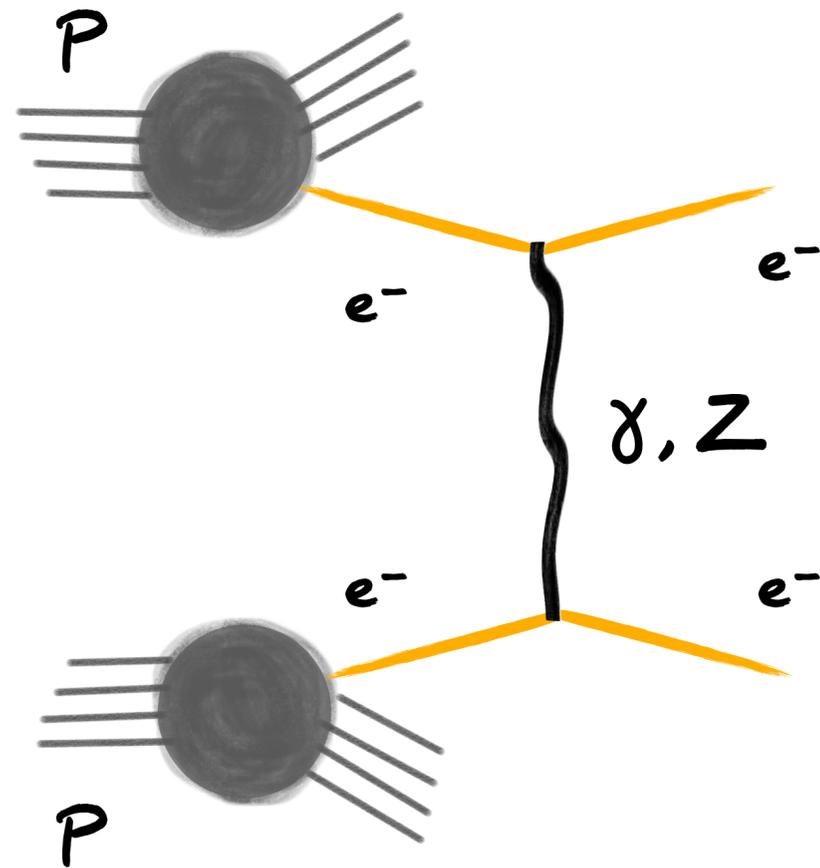
# Summary

- Precision determination of lepton PDFs opens up new ways to test SM (e.g.  $l^\pm l^\pm$  production) & to search for new physics @ the LHC
- Resonant LQ production allows to probe so far unexplored parameter space & has discovery potential
- Further theoretical developments needed to achieve next-to-leading order (NLO) plus parton shower (PS) accuracy for fiducial LQ cross sections

# LQ searches triggered by B anomalies



# Same sign lepton-pair production @ LHC



Signal events after cuts:

$$N_{\text{HL-LHC}}(e^-e^-) \simeq 700,$$

$$N_{\text{HL-LHC}}(\mu^-\mu^-) \simeq 550,$$

$$N_{\text{HL-LHC}}(\tau^-\tau^-) \simeq 250$$

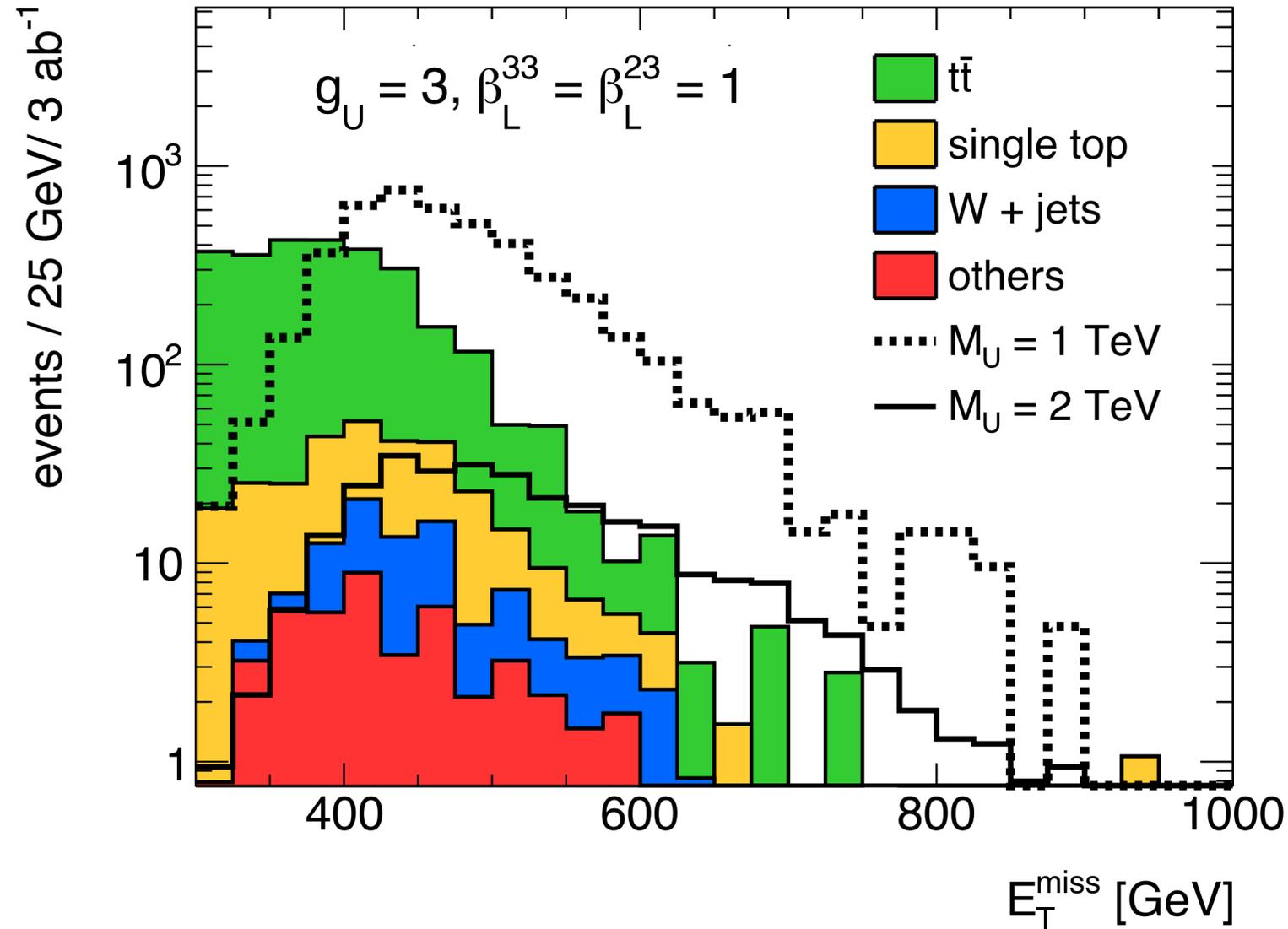
Dominant SM background from W-W-  
production after same cuts close to 0

# Simulation of 1<sup>st</sup> & 2<sup>nd</sup> resonant LQ signals

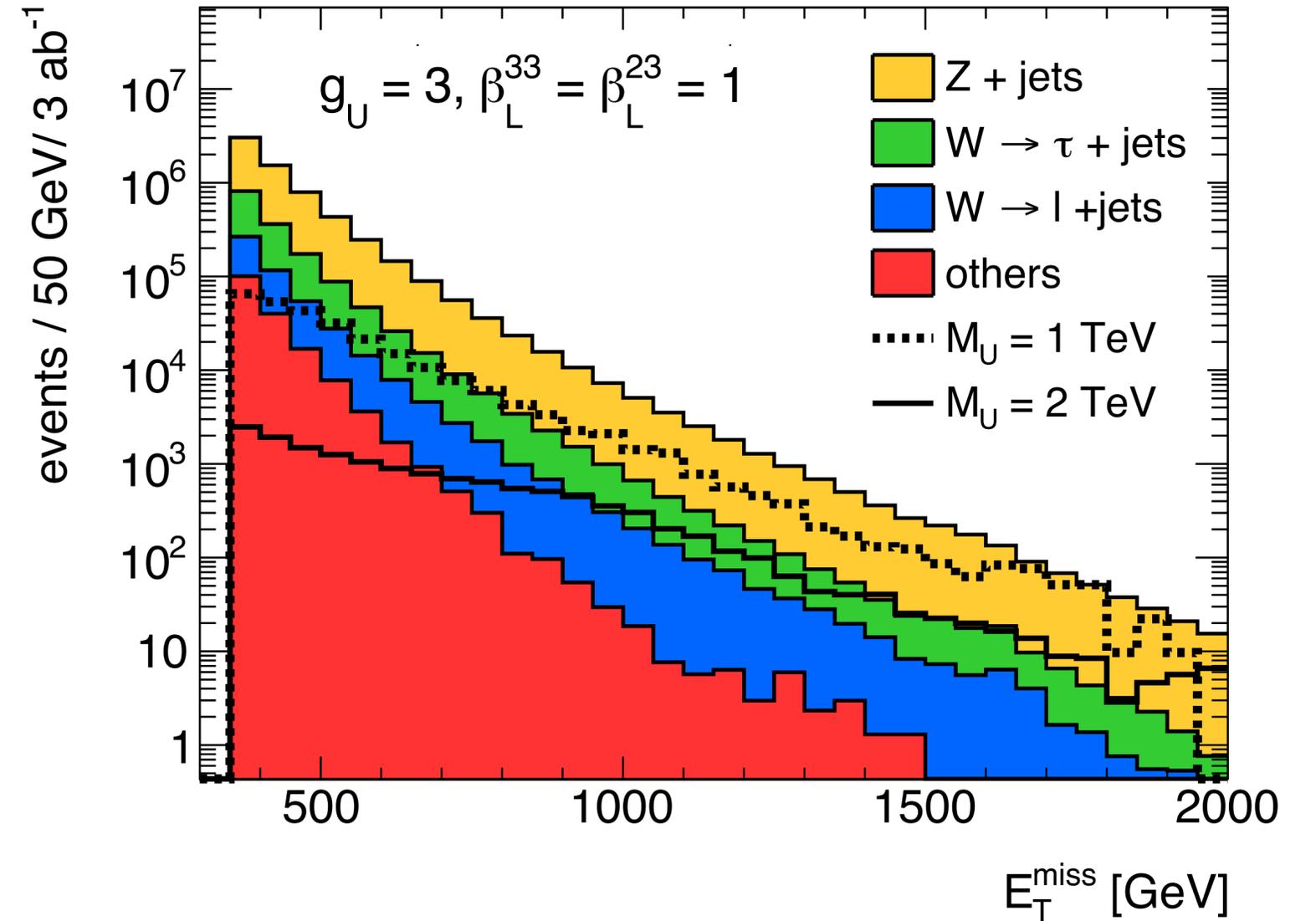
- Since PYTHIA currently cannot handle incoming leptonic partons, initial-state leptons have been replaced by photons to shower events. Our simulations do thus not include leptons but quarks from photon splitting in PS backward evolution
- As a result, jet- & lepton-veto induce a mismodelling of signal strength. By studying process  $q\gamma \rightarrow LQ \rightarrow q\ell^+\ell^-$ , we estimate this effect to be of  $O(10\%)$  & therefore to only mildly affect derived LQ limits
- Above PS issue needs to be resolved before NLO QCD & QED corrections for LQ signal can be correctly included in differential fashion

# Mono-top & mono-jet distributions

LHC 14 TeV, mono-top



LHC 14 TeV, mono-jet



# Prospects of LQ search strategies

