

# The Hubble Tension as a Signal of low-scale Leptogenesis and Neutrino Mass Generation

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[ArXiv:1909.04044](https://arxiv.org/abs/1909.04044), EPJC 80 (2020) 4, 294

[ArXiv:2004.01470](https://arxiv.org/abs/2004.01470), NuPhys19 Proceedings with Sam Witte

ArXiv:2102.XXXX, to appear soon

**IPPP, Durham University**

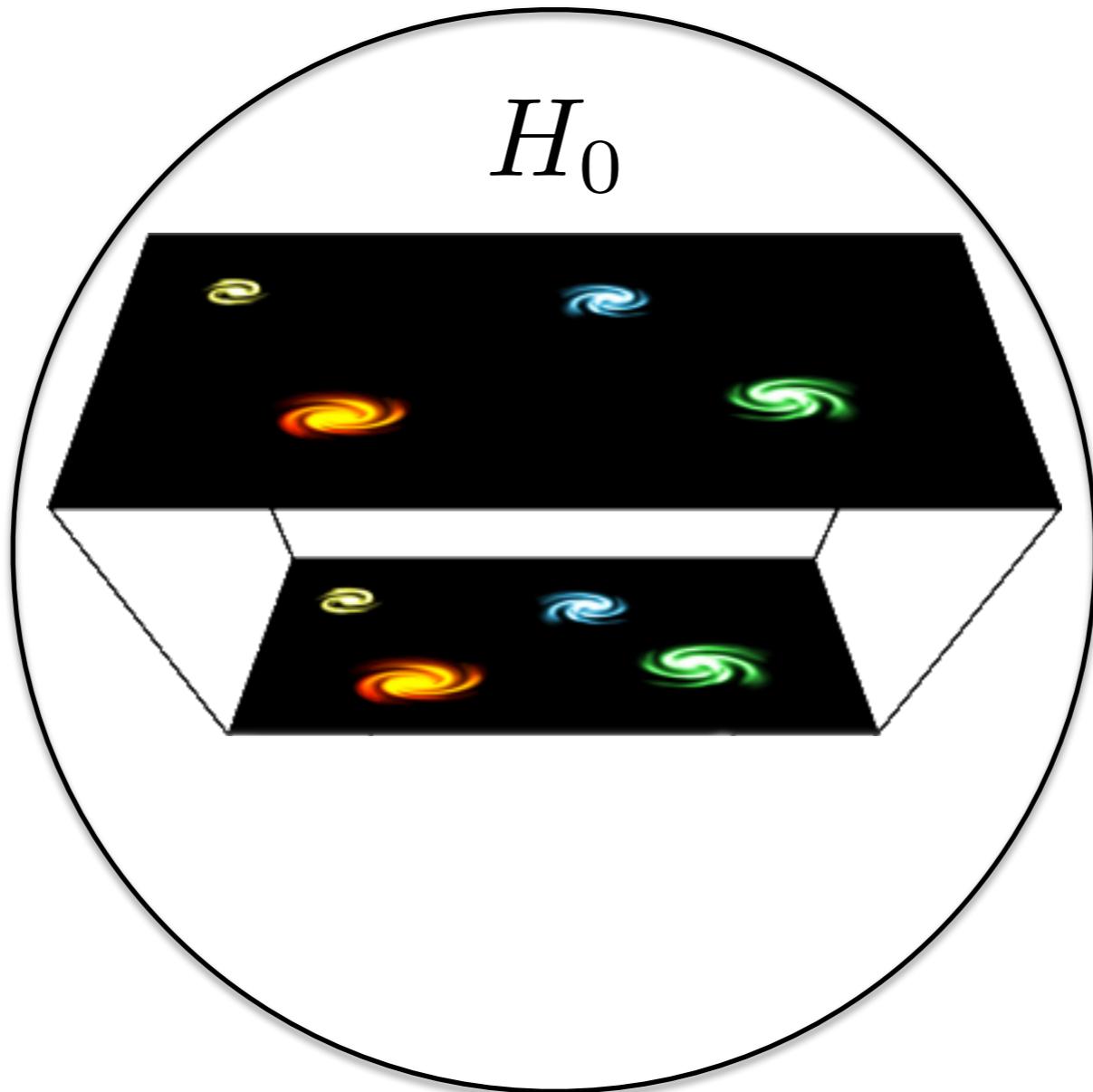
**14-01-2021**



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**Alexander von Humboldt**  
Stiftung / Foundation



**Riess *et al.* 1903.07603**

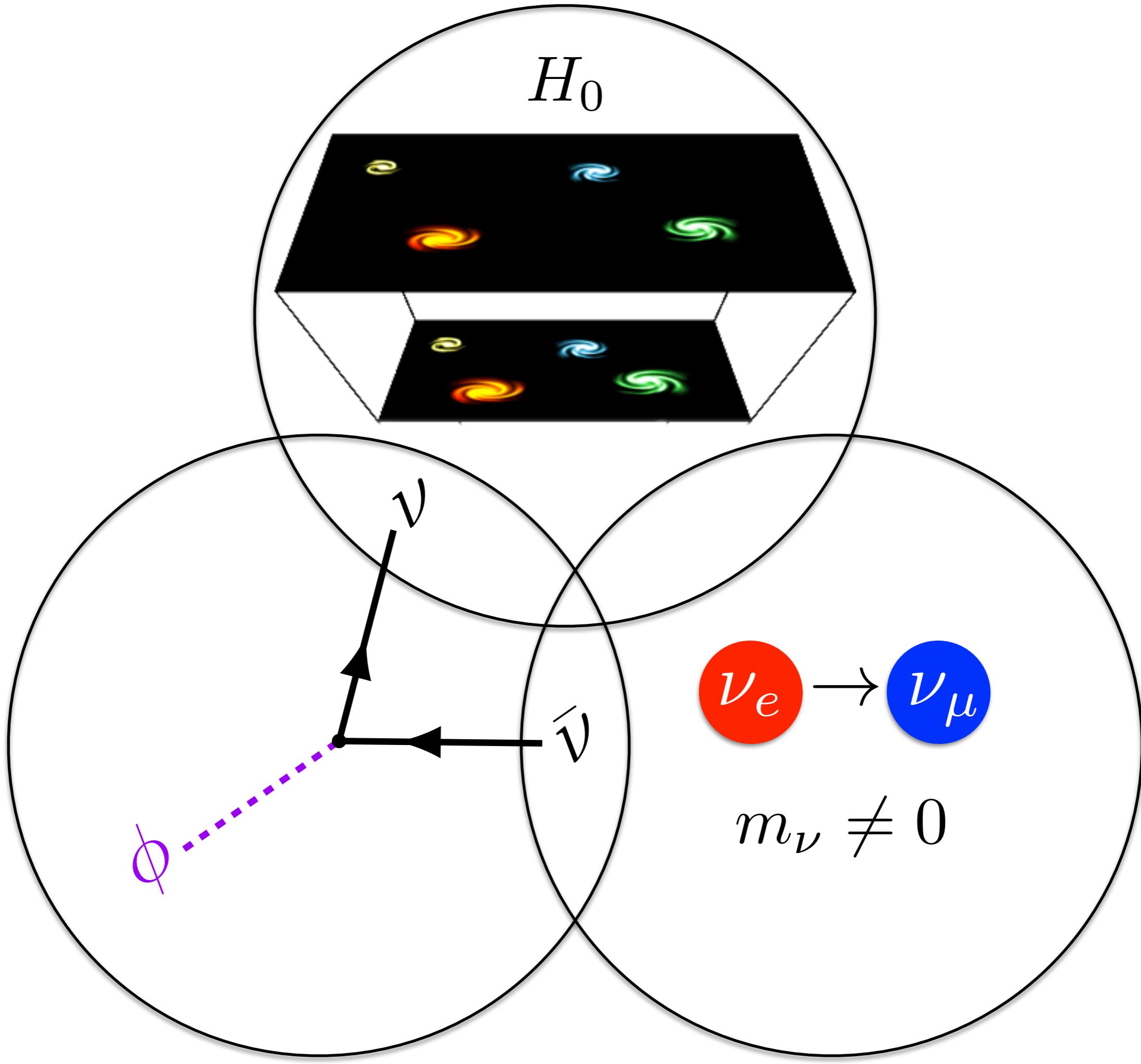
### Local Measurements

$$H_0 = 74.03 \pm 1.42 \text{ km s}^{-1} \text{ Mpc}^{-1}$$

**4.4  $\sigma$  tension  
within  $\Lambda$ CDM!**

### CMB Measurements

$$H_0 = 67.36 \pm 0.54 \text{ km s}^{-1} \text{ Mpc}^{-1}$$



# The Hubble tension and the Majoron

## Escudero & Witte 19'

- Can substantially ameliorate the tension for:
- Couplings from seesaw
- Mass from Planck-scale physics
- $\Delta N_{\text{eff}} \sim 0.5$ ?

$$\begin{aligned}m_\phi &\sim (0.1 - 1) \text{ eV} \\ \lambda &\simeq 10^{-13} - 10^{-12} \\ \Delta N_{\text{eff}} &\sim 0.5\end{aligned}$$

$$\lambda \simeq 10^{-13} \frac{m_\nu}{0.05 \text{ eV}} \frac{\text{TeV}}{v_L}$$
$$v_L \sim (0.1 - 1) \text{ TeV}$$

$$m_\phi \sim (0.1 - 1) \text{ eV}$$

## Escudero & Witte 21'

- $\Delta N_{\text{eff}} \sim 0.5$  as a product of low-scale Baryogenesis
- $H_0$  tension as a signal of  $m_\nu \neq 0$  and Leptogenesis

# Outline

## 1) Neutrinos in $\Lambda$ CDM

Neff

## 2) The Hubble Tension

Neutrinophilic approaches

## 3) The Majoron

The singlet Majoron model

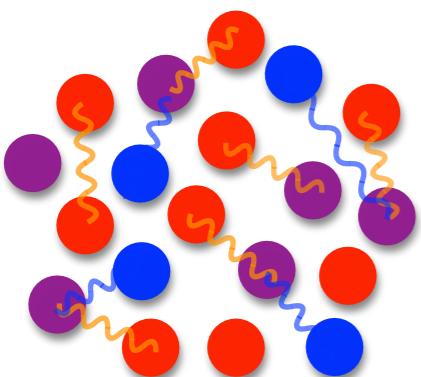
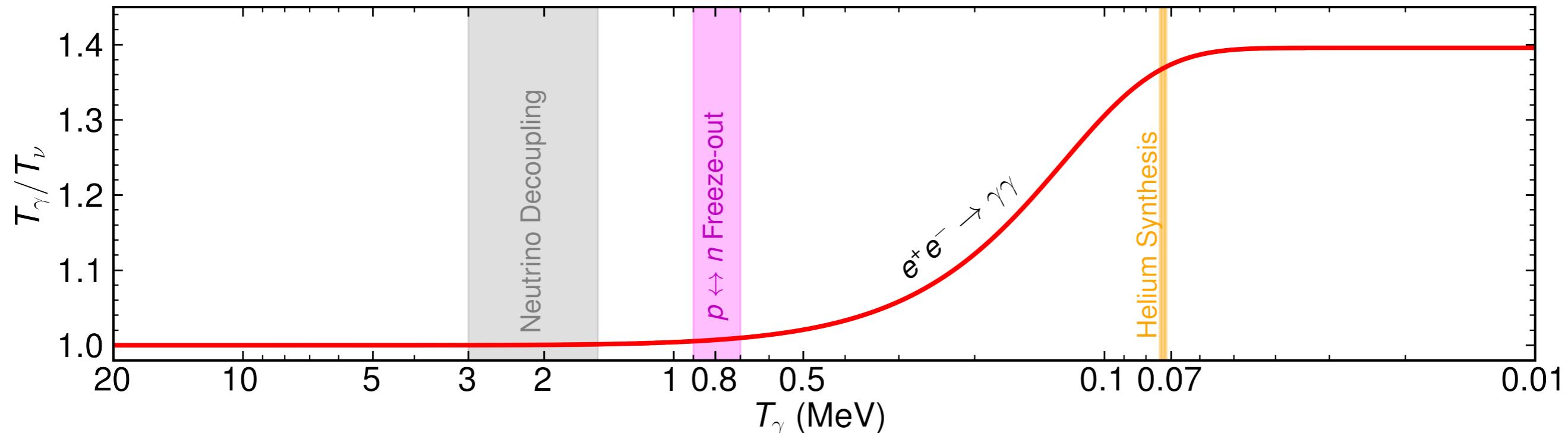
The Majoron as a potential solution to the H0 tension

## 4) ARS Leptogenesis within the Majoron model

## 5) Conclusions and Outlook

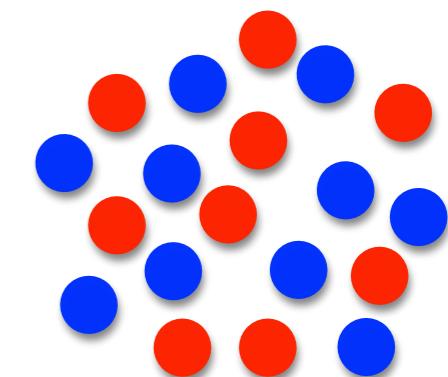
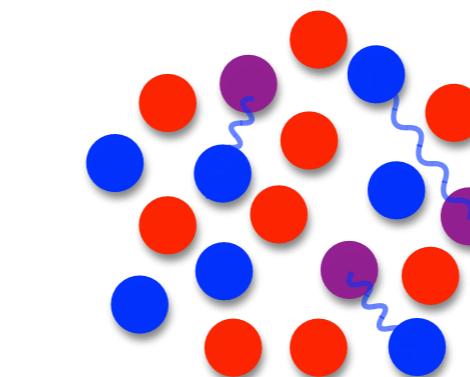
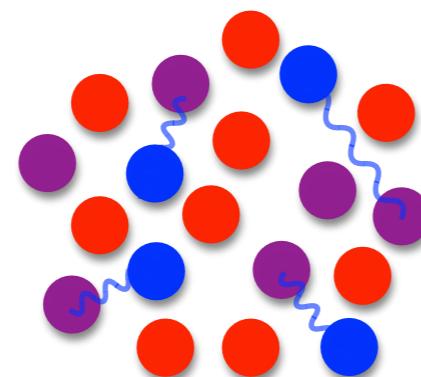
# Neutrino Decoupling

Evolution in the Standard Model



$$e^+ e^- \leftrightarrow \bar{\nu}_i \nu_i$$

$$e^\pm \nu_i \leftrightarrow e^\pm \nu_i$$



**Neutrinos**



**Electrons**



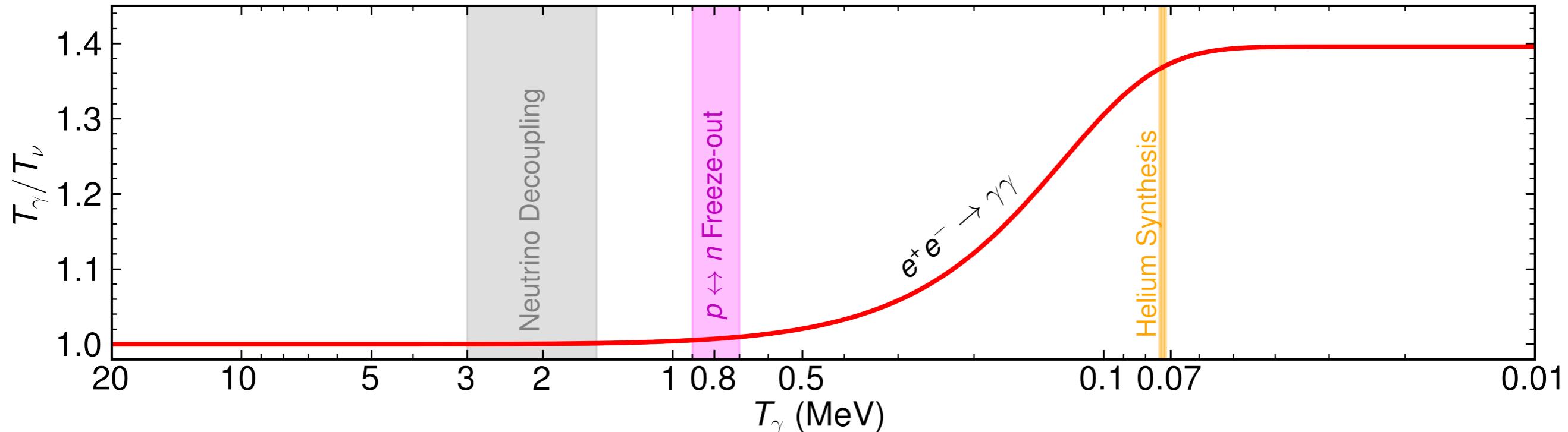
**Photons**



**Z-W (off-shell)**

# Neutrino Decoupling

Evolution in the Standard Model



- $N_{\text{eff}} \equiv \frac{8}{7} \left( \frac{11}{4} \right)^{4/3} \left( \frac{\rho_{\text{rad}} - \rho_{\gamma}}{\rho_{\gamma}} \right)$

$$N_{\text{eff}} = 3 \left( \frac{1.4 T_{\nu}}{T_{\gamma}} \right)^4$$

- $N_{\text{eff}}^{\text{SM}} = 3.044(1)$

de Salas & Pastor 1606.06986  
Bennett, Buldgen, Drewes, Wong 1911.04504  
Escudero 2001.04466

Akita & Yamaguchi 2005.07047  
Froustey, Pitrou & Volpe 2008.01074  
Gariazzo, de Salas, Pastor et al. 2012.02726  
Hansen, Shalgar, Tamborra 2012.03948

## Relic Neutrino Decoupling

$$t \sim 0.1 \text{ s}$$

$$T_{\nu} \sim 2 \text{ MeV}$$

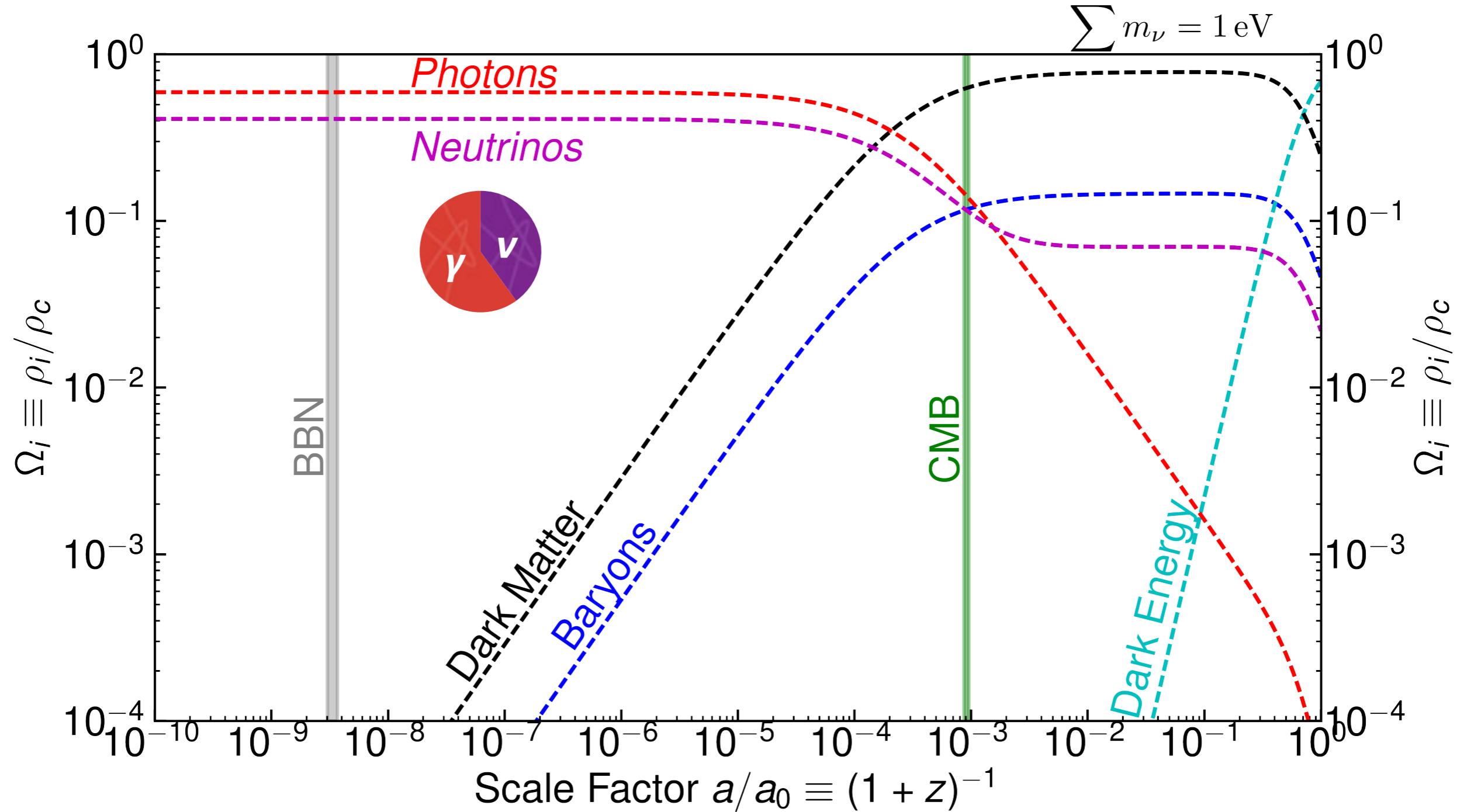
## Why is it not 3?

- Some  $e^+e^-$  heating
- Non-instantaneous decoupling
- QED thermal corrections
- Neutrino Oscillations

Excellent review  
by Dolgov hep-ph/0202122

# Neutrino Evolution

Neutrinos are always a relevant species in the Universe's evolution



**Non-Rel:**  $z_\nu^{\text{non-rel}} \simeq 600 \frac{m_\nu}{0.3 \text{ eV}}$

**DM:**  $\Omega_\nu h^2 = \sum m_\nu / (93.14 \text{ eV})$

# Current Constraints on Neff

## Big Bang Nucleosynthesis

$N_{\text{eff}}$  is, at the moment, constrained by  $Y_p$  as a result of change in the time to temperature relationship

There are various  $Y_p$  measurements (see Pisanti et al. 2011.11537):

$$Y_P = 0.2446 \pm 0.0029 \rightarrow N_{\text{eff}} = 2.84 \pm 0.20$$

$$Y_P = 0.2449 \pm 0.0040 \rightarrow N_{\text{eff}} = 2.86 \pm 0.28 \quad (\text{most used, PDG but with warning})$$

$$Y_P = 0.2551 \pm 0.0022 \rightarrow N_{\text{eff}} = 3.60 \pm 0.17$$

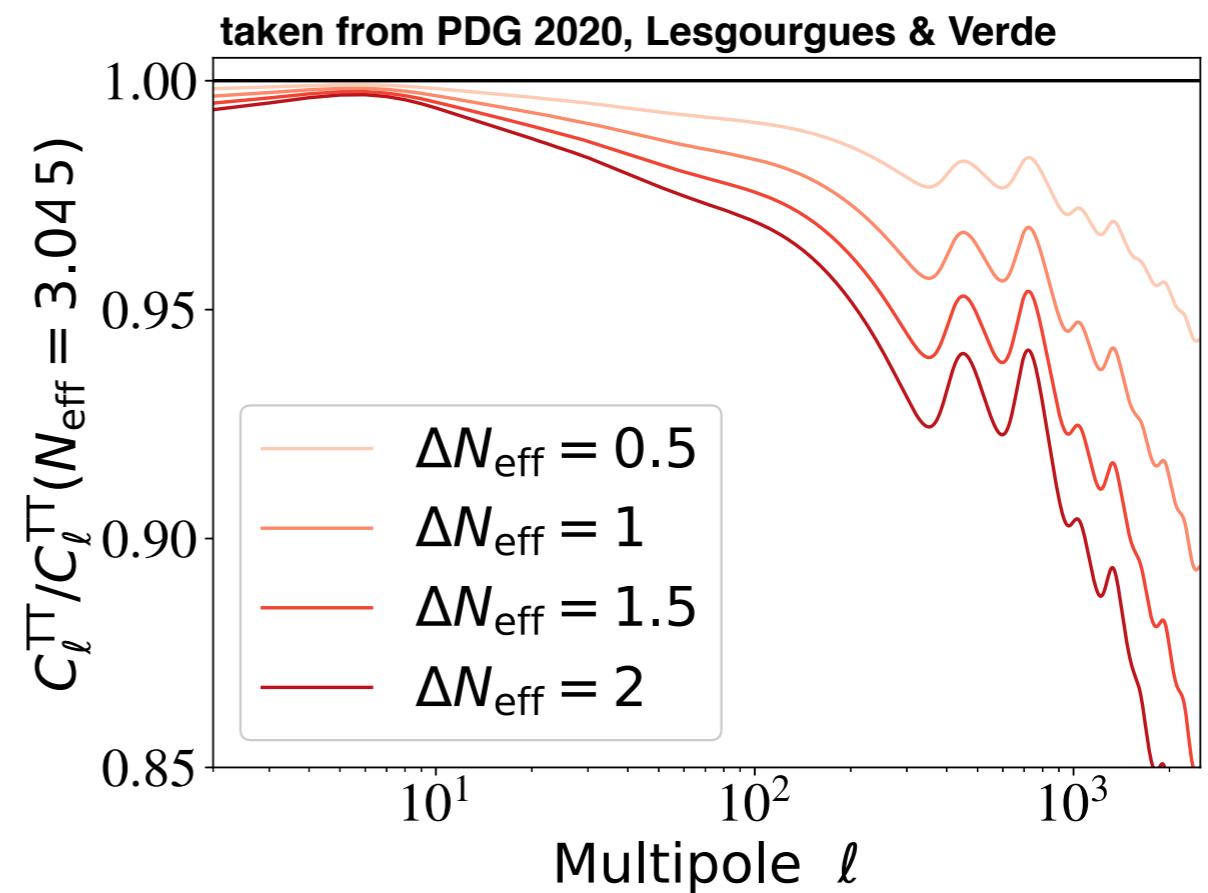
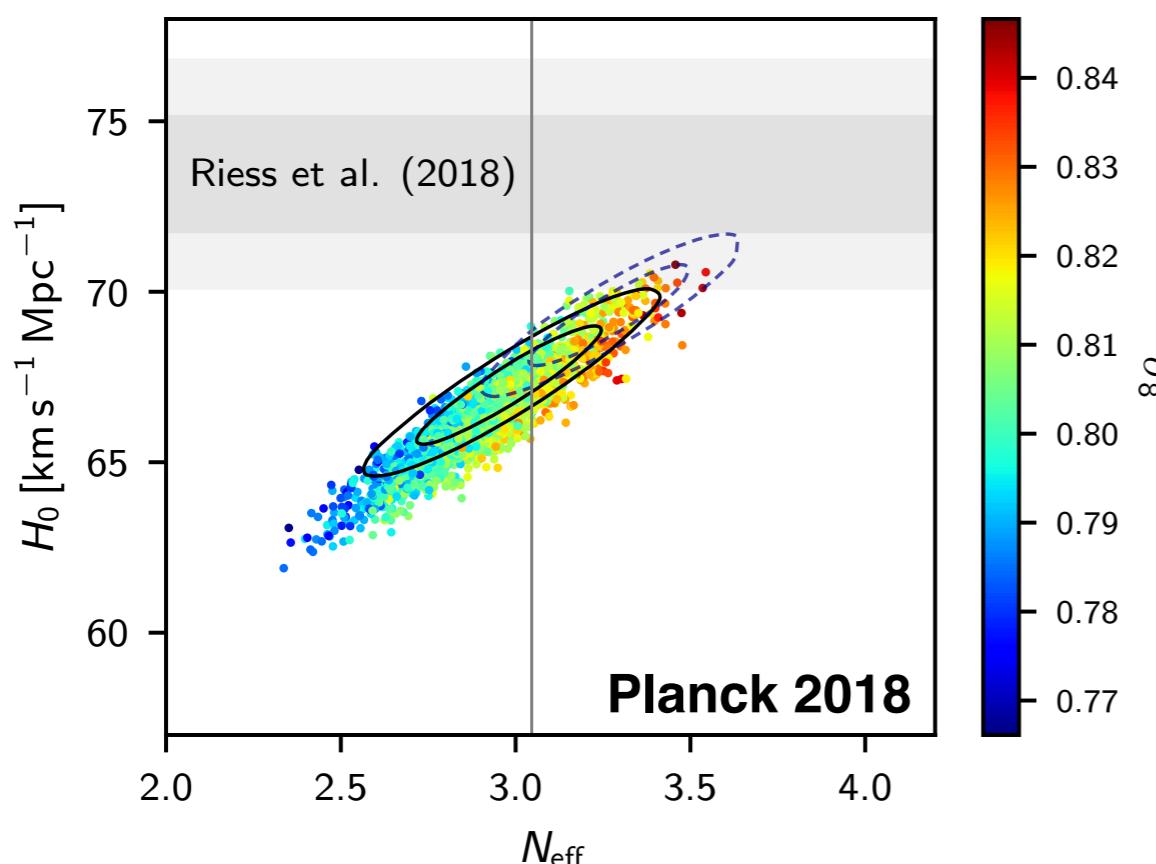
$$Y_P = 0.2436 \pm 0.0040 \rightarrow N_{\text{eff}} = 2.78 \pm 0.28$$

**Take away:** likely  $\Delta N_{\text{eff}}^{\text{BBN}} < 0.4$  but it could be that  $\Delta N_{\text{eff}}^{\text{BBN}} = 0.55$

# Current Constraints on $N_{\text{eff}}$

Planck

$N_{\text{eff}}$  is constrained by the high- $\ell$  multipoles,  
i.e. Silk Damping



$$N_{\text{eff}}^{\text{CMB+BAO}} = 2.99 \pm 0.17$$

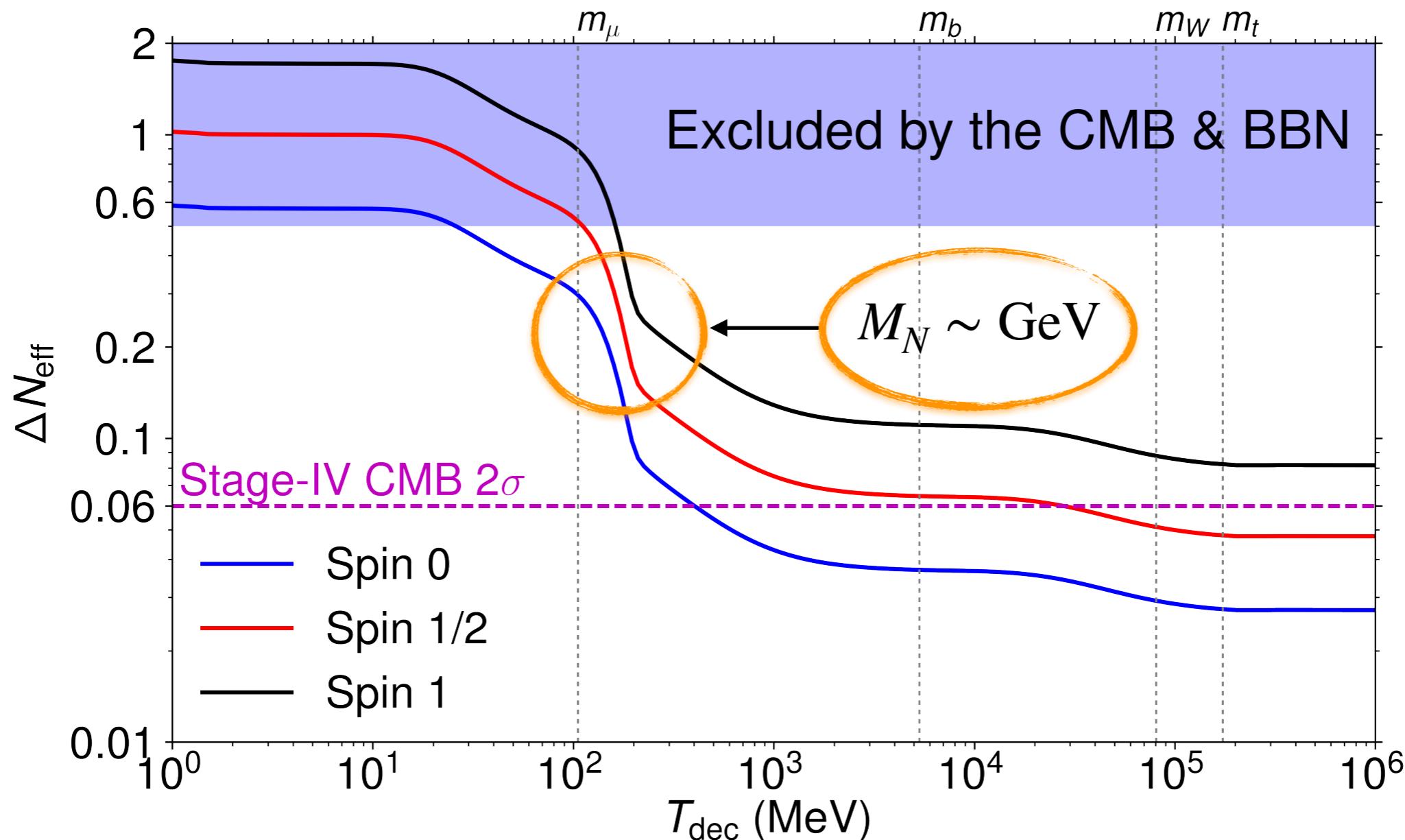
When taking the Riess value as a prior, then the value of  $N_{\text{eff}}$  is

$$N_{\text{eff}}^{\text{CMB+BAO}+H_0} = 3.27 \pm 0.15$$

$\Delta N_{\text{eff}} \sim 0.2\text{-}0.5$  ameliorates the tension but is not favoured by Planck data!

# Neff Scenarios

- **Sterile Neutrino**  $m_N \sim \text{eV}$   $\Delta N_{\text{eff}} = 1$  (e.g. Gariazzo, de Salas & Pastor 1905.11290)
- **Goldstone Bosons** Weinberg 1305.1971
- **Other sterile long-lived particles** Gravitino, axino, hidden sector particles ...



# The Hubble Tension

- The Hubble Tension:

$$H_0 = 74.03 \pm 1.42 \text{ km s}^{-1} \text{ Mpc}^{-1}$$

Riess *et al.* 1903.07603

$$H_0 = 67.36 \pm 0.54 \text{ km s}^{-1} \text{ Mpc}^{-1}$$

Planck 2018 1807.06209

**4.4  $\sigma$  tension within  $\Lambda$ CDM!**

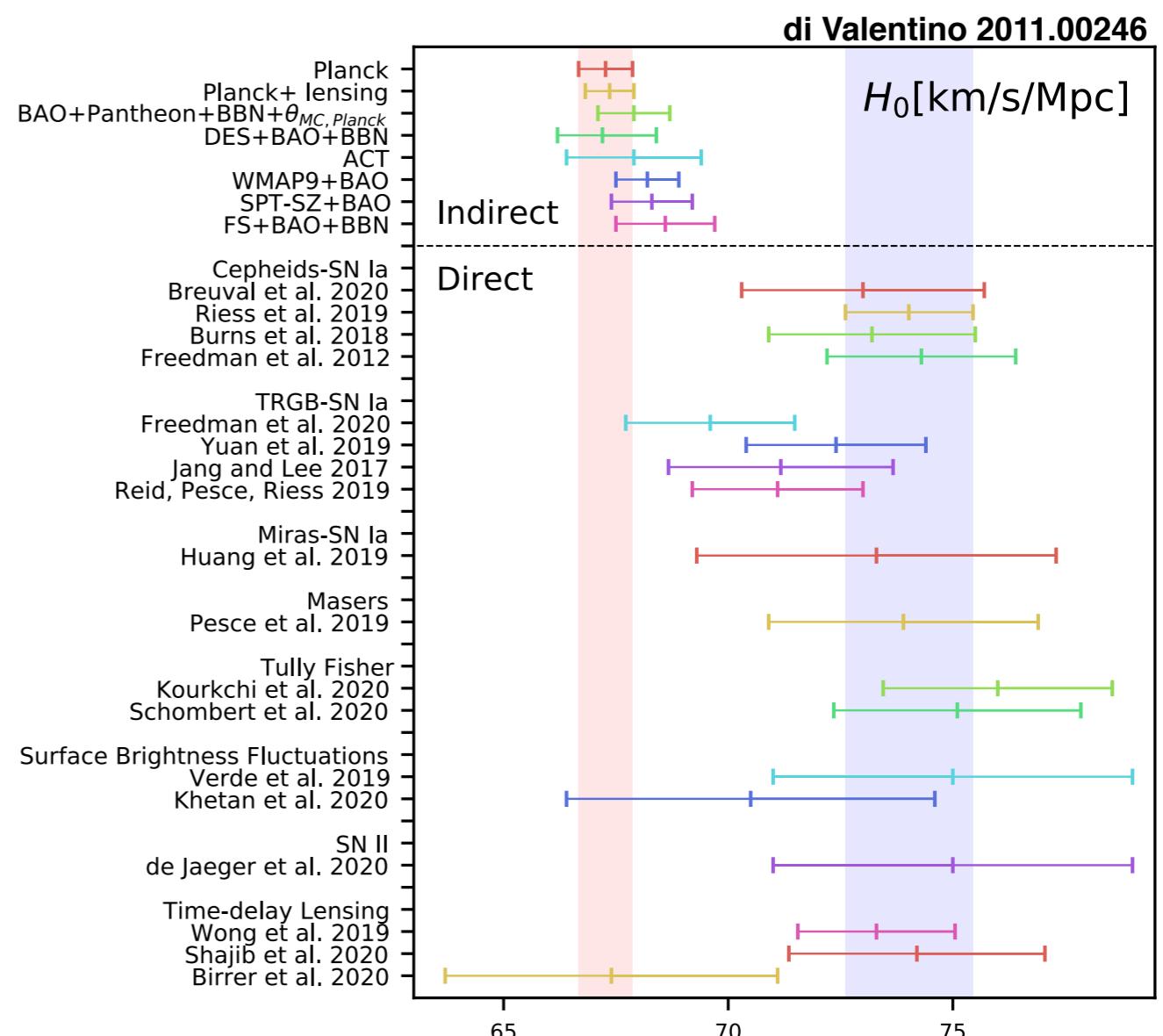
- A pattern has clearly emerged:

- 4-6  $\sigma$  tension depending upon the datasets included

see Verde, Treu, Riess 1907.10625 for a review

- Baryon Acoustic Oscillations point to small  $H_0$

- Cepheids+Type-Ia SN and Strong lensing are among the most precise and they point to  $H_0 \sim (74 \pm 2) \text{ km/s/Mpc}$



# The Hubble Tension

- Possible resolutions:

1) Systematics in the CMB data very unlikely

2) Systematics in local measurements none so far

3) New feature of  $\Lambda$ CDM

- Possibilities Beyond  $\Lambda$ CDM (Knox and Millea 1908.03663):

1) Late Universe Modifications very unlikely

2) Early Universe Modifications hard but doable

# The Hubble Tension: Theory

- Way to Resolve the Hubble Tension (Knox and Millea 1908.03663):

**Enhance the expansion history of the Universe prior and close to recombination!**

**CMB fixes:**  $\theta_s \equiv r_s/D_M(z_*)$  (0.03% precision)

$$r_s = \int_{z_*}^{\infty} \frac{c_s}{H(z')} dz'$$

**Comoving sound horizon  
(Early Universe)**

$$D_M(z) = \int_0^z \frac{1}{H(z')} dz'$$

$H_0$

**Comoving angular diameter distance  
(Late Universe)**

**By how much?**

$$H_0 \simeq [73.6 + 6.2(\Delta N_{\text{eff}} - 1)] \text{ km/s/Mpc}$$

see Vagnozzi 1907.07569

# Neutrinos and the Hubble Tension

## Why Neutrinos?

- 1) Neutrinos are always a relevant species in the evolution of the Universe**
  
- 2) Neutrino masses are the only Laboratory evidence of Physics Beyond the Standard Model**

# Neutrinos and the Hubble Tension

## ● Dark Radiation

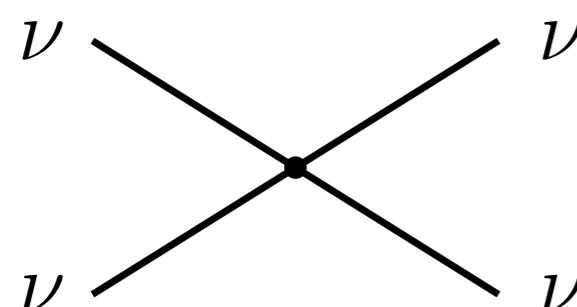
$$\Delta N_{\text{eff}} = 0.23 \pm 0.15$$

(68 % CL, Planck+BAO+H<sub>0</sub>)

**Clear Interpretation**  
H<sub>0</sub> tension from 4.4σ to 3σ  
CMB fit is degraded



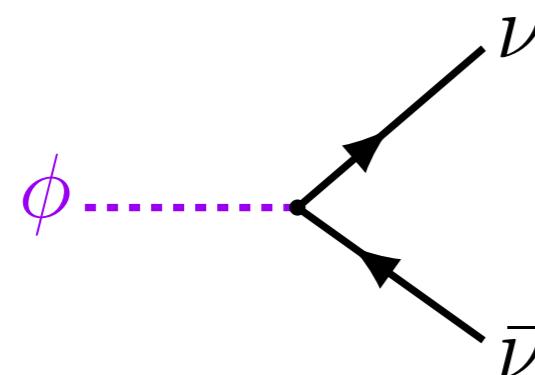
## ● Strong Neutrino Scattering + Dark Radiation Kreisch, Cyr-Racine, Doré 1902.00543



H<sub>0</sub> tension solved if TEEE data is ignored  
If pol data is included no solution for H<sub>0</sub>  
Almost excluded by Lab data (Blinov++1905.02727)



## ● Light Neutrinophilic Scalar + Dark Radiation Escudero & Witte 1909.04044



H<sub>0</sub> tension from 4.4σ to 2.5σ  
CMB fit is not degraded  
Direct connection with Seesaw  
Ad hoc ΔN<sub>eff</sub> ~ 0.5



# Neutrinos and the Hubble Tension

## ● Early Dark Energy sourced by neutrinos

Sakstein & Trodden 1911.11760

Nice way to solve the coincidence problem



Use  $\sum m_\nu = 1.5 \text{ eV}$  (10% of DM) which can be dangerous



But the Cl's have not been calculated yet ...



Some progress has been made Carrillo González et al. 2011.09895



## ● An eV-scale Sterile Neutrino interacting with a pseudoscalar

Archidiacono, Hannestad, Hansen & Tram 1404.5915, 1508.02504

Archidiacono, Gariazzo, Giunti, Hannestad, Hansen, Laveder, Tram 1606.07673

Archidiacono, Gariazzo, Giunti, Hannestad, Tram 2006.12885

Clearly motivated by short-baseline neutrino experiments



Nice idea to try to avoid the cosmo problems with  $m_s \sim \text{eV}$



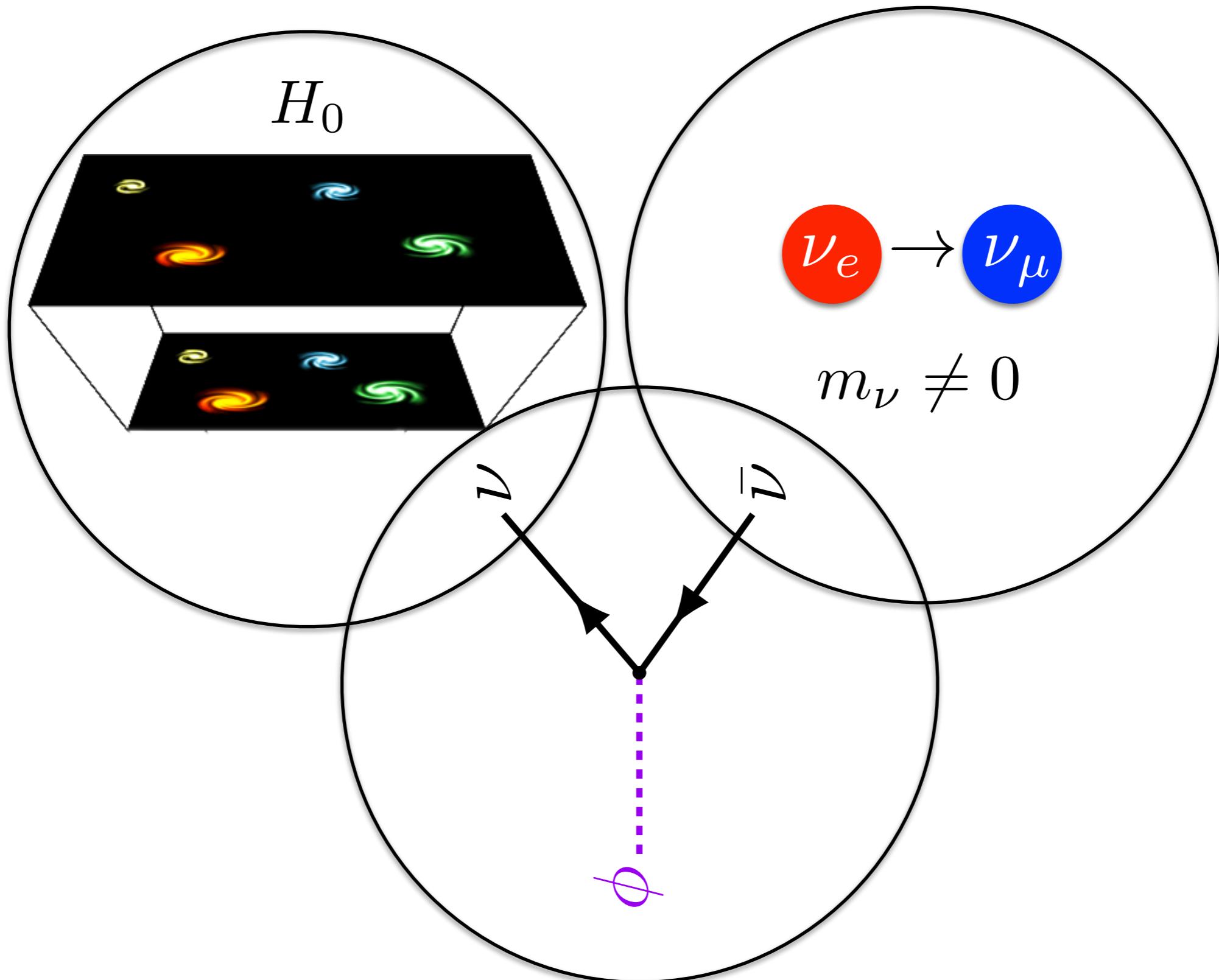
The Hubble Tension could be solved if  $\Delta N_{\text{eff}} = 1$



But that leads to a bad CMB fit  $\Delta\chi^2 = 13 - 32$



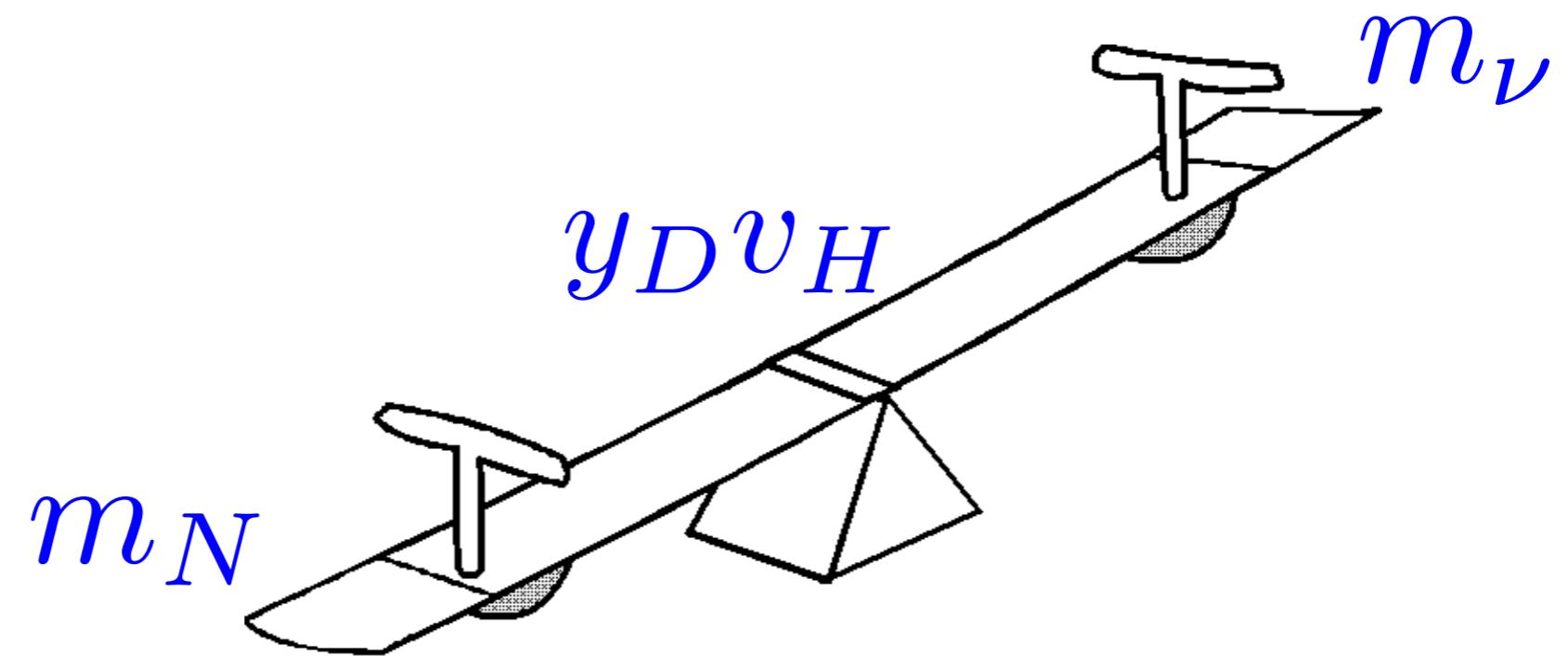
# The Idea:



# The Seesaw Mechanism

Minkowski, Yanagida, Gell-Mann, Ramond, Slansky, Glashow, Mohapatra, Senjanovic, Schechter, Valle

Type-I seesaw



Neutrinos are very light Majorana particles:

$$m_\nu \simeq 0.03 \text{ eV} \left( \frac{y_D}{10^{-6}} \right)^2 \frac{\text{TeV}}{M_N}$$

# The Scenario

## Spontaneously Broken Symmetry Global $U(1)_L$

Chikashige, Mohapatra, Peccei (1981)

**Sterile Neutrinos**  $\mathcal{L} = -\frac{\lambda_{N_{ij}}}{\sqrt{2}} \Phi \bar{N}_{R,i} N_{R,j}^c - h_{\alpha i} \bar{L}_L^\alpha H N_{Ri} + \text{h.c.},$   $L[\Phi] = 2$   
 $L[N] = 1$

SSB:  $\Phi \rightarrow v_L/\sqrt{2} \longrightarrow M_N = \lambda_N v_L \xrightarrow{\text{seesaw}} m_\nu \simeq h^2 v_H^2 / (2M_N)$

**Scalar Sector**  $V_\Phi = -\mu_\Phi^2 \Phi^\dagger \Phi + \lambda_\Phi (\Phi^\dagger \Phi)^2 - \lambda_{\Phi H} (H^\dagger H) (\Phi^\dagger \Phi)$

|  |                                     |  |
|--|-------------------------------------|--|
| $\Phi = \frac{v_L + \rho}{\sqrt{2}} e^{i\phi/v_L}$ | $\rho \equiv \text{CP-even scalar}$ | $m_\rho^2 = 2\lambda_\Phi v_L^2$           |
|  | $\phi \equiv \text{Majoron}$        | <b>pseudo-Goldstone:</b> $m_\phi \simeq 0$ |

**Interactions**  $\mathcal{L}_{\text{eff}} = -\frac{\lambda_N}{2} [\rho \bar{N} N - i\phi \bar{N} \gamma_5 N]$   $\lambda_\nu \ll \lambda_{N\nu} \ll \lambda_N$   
 $- \frac{\lambda_{N\nu}}{2} [\rho \bar{N} \nu - i\phi \bar{N} \gamma_5 \nu] + \text{h.c.}$   $\lambda_\nu \simeq |\theta| \lambda_{N\nu} \simeq |\theta|^2 \lambda_N$   
 $- \frac{\lambda_\nu}{2} [\rho \bar{\nu} \nu - i\phi \bar{\nu} \gamma_5 \nu]$   $|\theta|^2 \simeq 5 \times 10^{-11} \frac{m_\nu}{0.05 \text{ eV}} \frac{1 \text{ GeV}}{M_N}$

# The Scenario

## Spontaneously Broken Symmetry Global $U(1)_L$

Chikashige, Mohapatra, Peccei (1981)

**The Majoron:**  $\phi$

$$\mathcal{L}_{\text{int}} = i\lambda \phi \bar{\nu} \gamma_5 \nu$$

**Very weakly interacting:**

$$\lambda \simeq 10^{-13} \frac{m_\nu}{0.05 \text{ eV}} \frac{\text{TeV}}{v_L} \quad (\text{seesaw})$$

**Extremely feebly interacting with matter:**  $\lambda_{\phi ee} \sim 10^{-20}$

**Dimension-5 Planck suppressed operators:**

$$m_\phi \sim v_L \sqrt{\frac{v_L}{M_{\text{Pl}}}} \lesssim 0.1 \text{ keV}$$

$$\Delta V = \beta (\Phi^\star \Phi)^2 \frac{\Phi^\star + \Phi}{M_{\text{Pl}}}$$

Rothstein, Babu, Seckel hep-ph/9301213  
Akhmedov, Berezhiani, Mohapatra, Senjanovic hep-ph/9209285

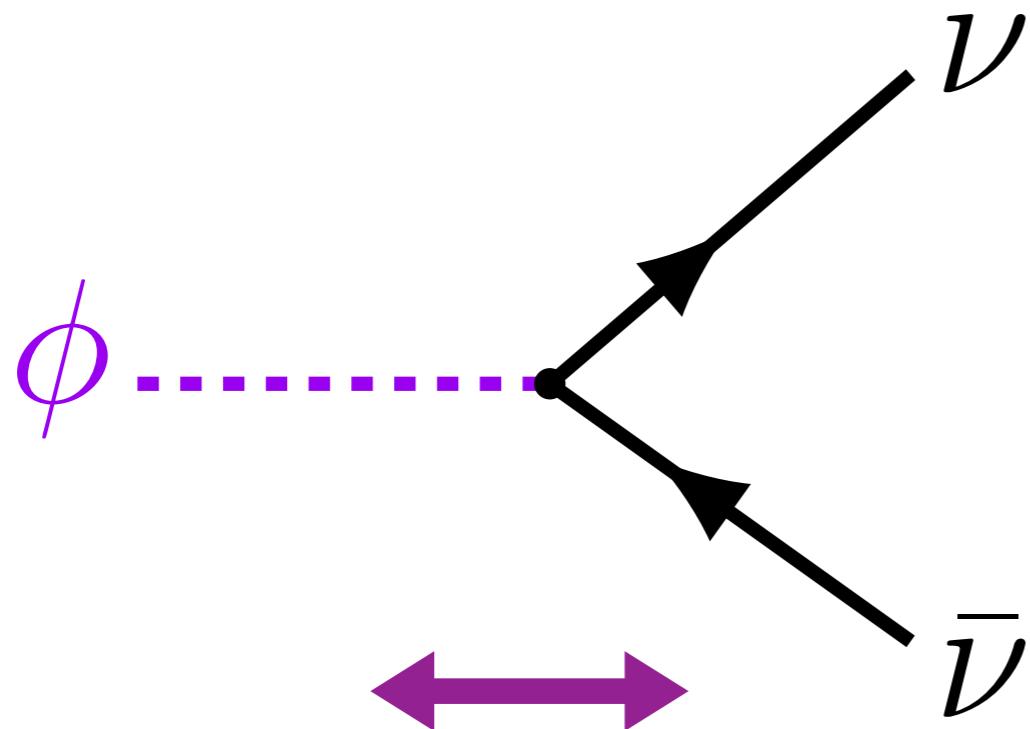
**Parameter Space:**

$$10^{-15} < \lambda < 10^{-3}$$
$$0.1 \text{ eV} < m_\phi < \text{MeV}$$

And assume that  $n_\phi = 0$  at BBN

# Cosmological Implications

Only Relevant Process:



**provided**  $\Gamma_\phi \geq H(T_\nu = m_\phi/3)$

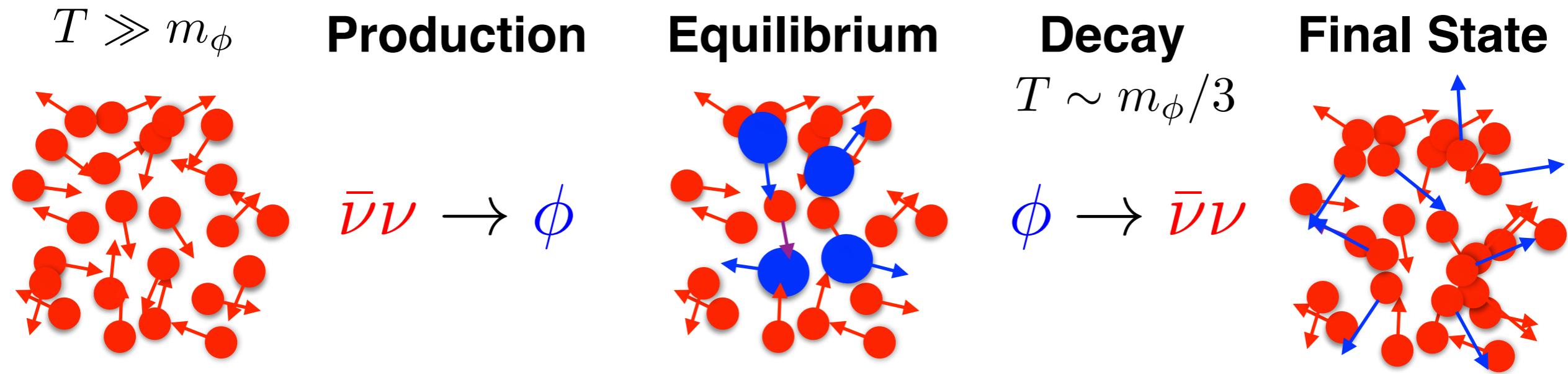
Two main effects:

Chacko, Hall, Okui, Oliver  
hep-ph/0312267

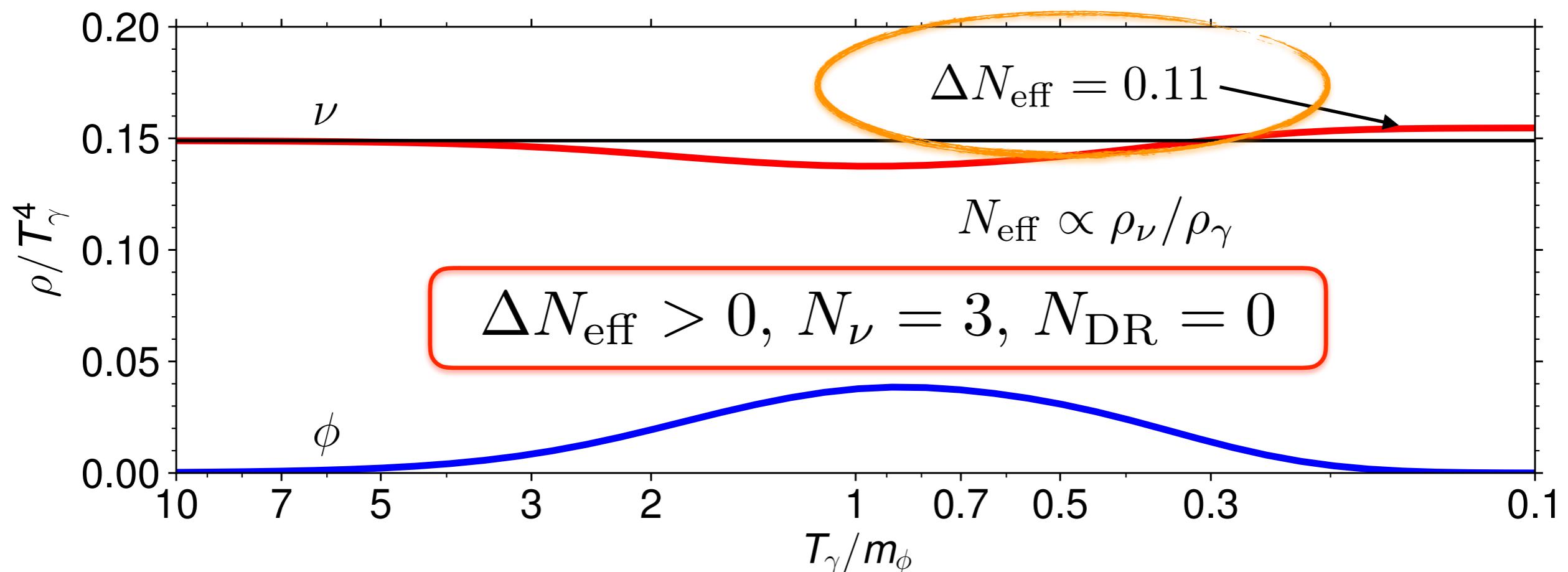
- Non-standard expansion history
- Erase the neutrino anisotropic stress

- We solve the Boltzmann equation for the background  
Escudero 1812.05605, 2001.04466
- We include the neutrino-majoron Boltzmann hierarchy in CLASS

# Cosmological Implications

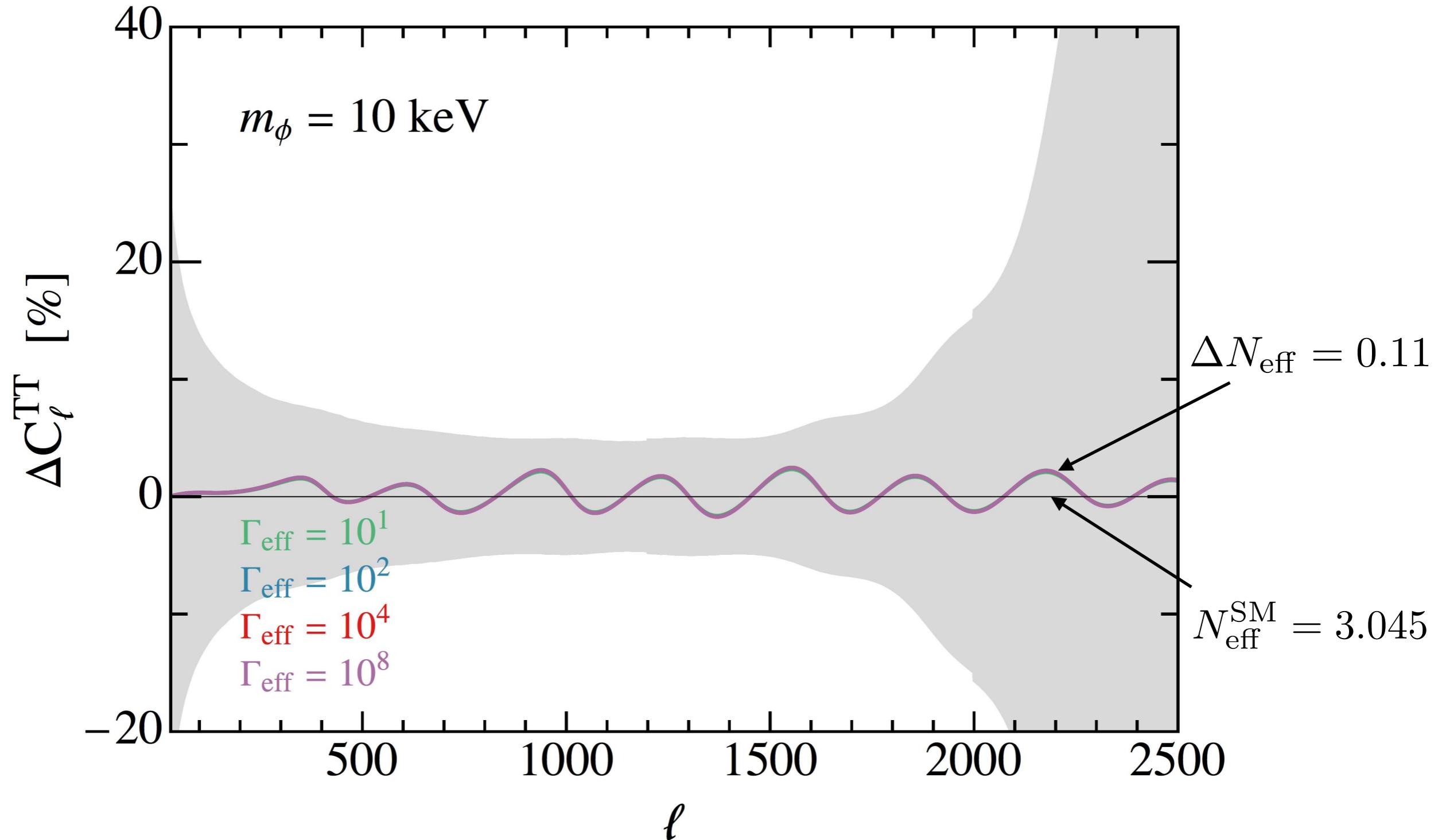


$$\Gamma_\phi \simeq H(T_\nu = m_\phi/3)$$



# Effects on the CMB

$$\Gamma_{\text{eff}} = \left( \frac{\lambda}{4 \times 10^{-12}} \right)^2 \left( \frac{1 \text{ keV}}{m_\phi} \right)$$



# Effects on the CMB

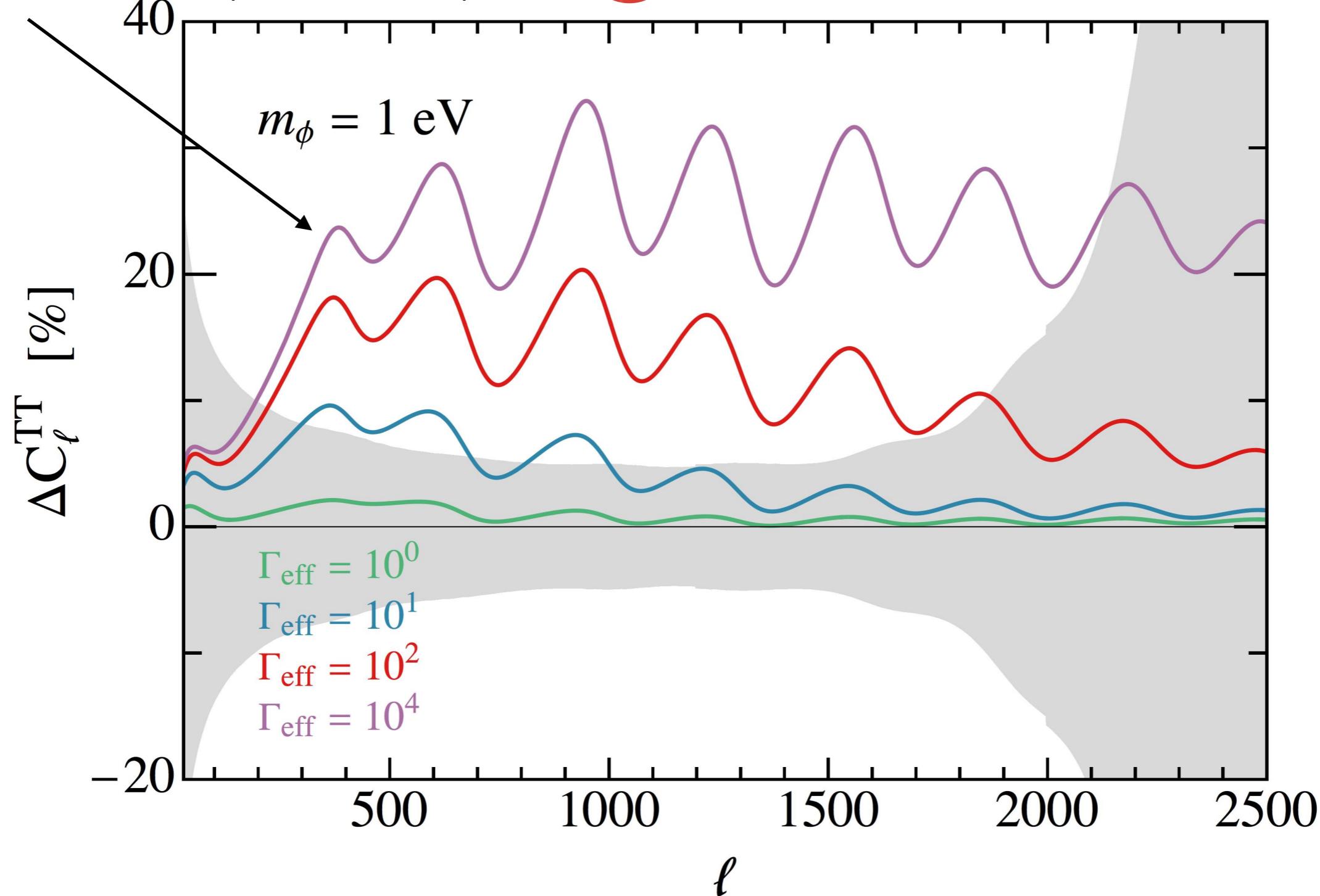
see Bashinsky and Seljak astro-ph/0310198

$$\sigma_\nu \rightarrow 0$$

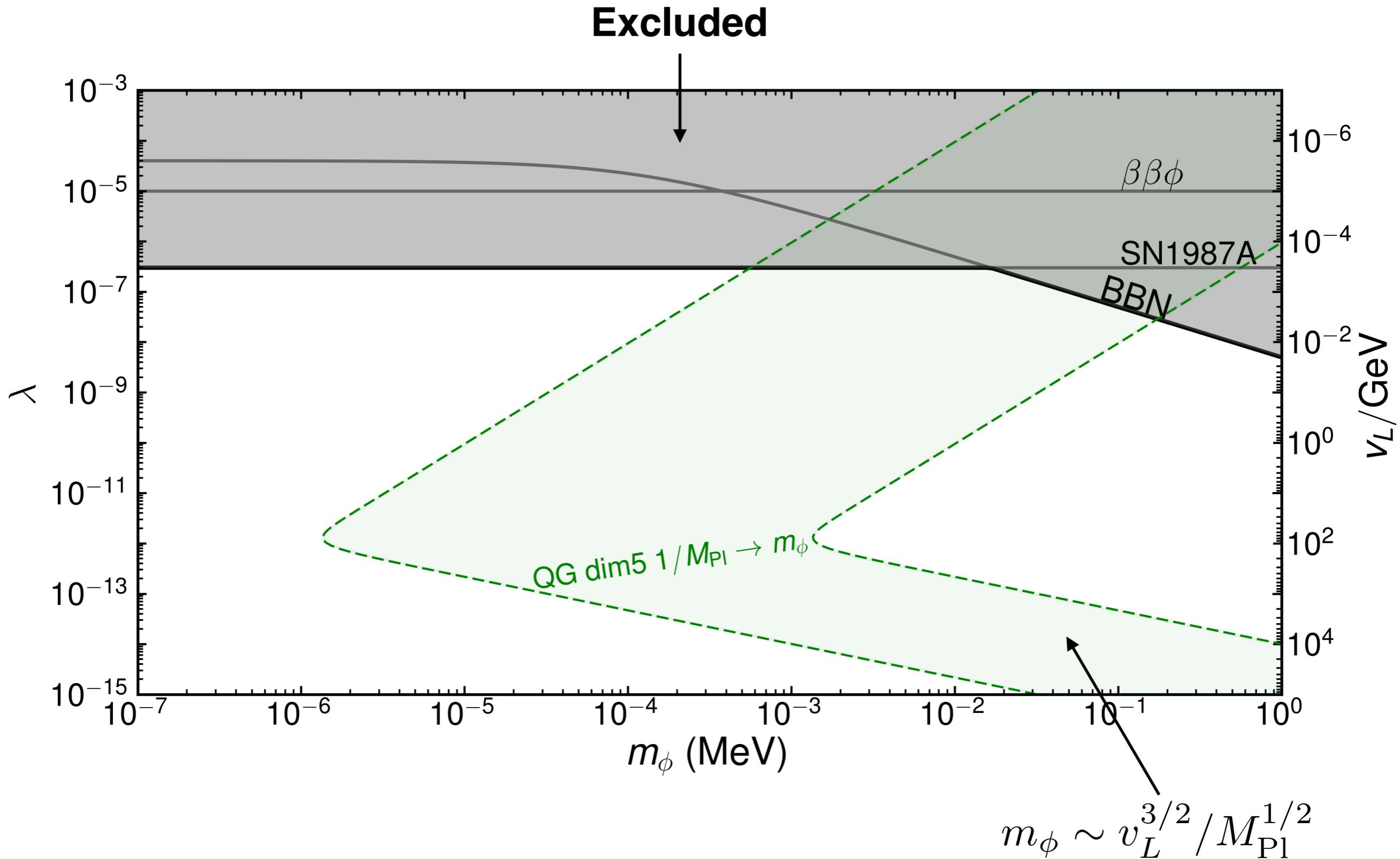
$$\delta G_{\mu\nu} = 8\pi G \delta T_{\mu\nu}$$



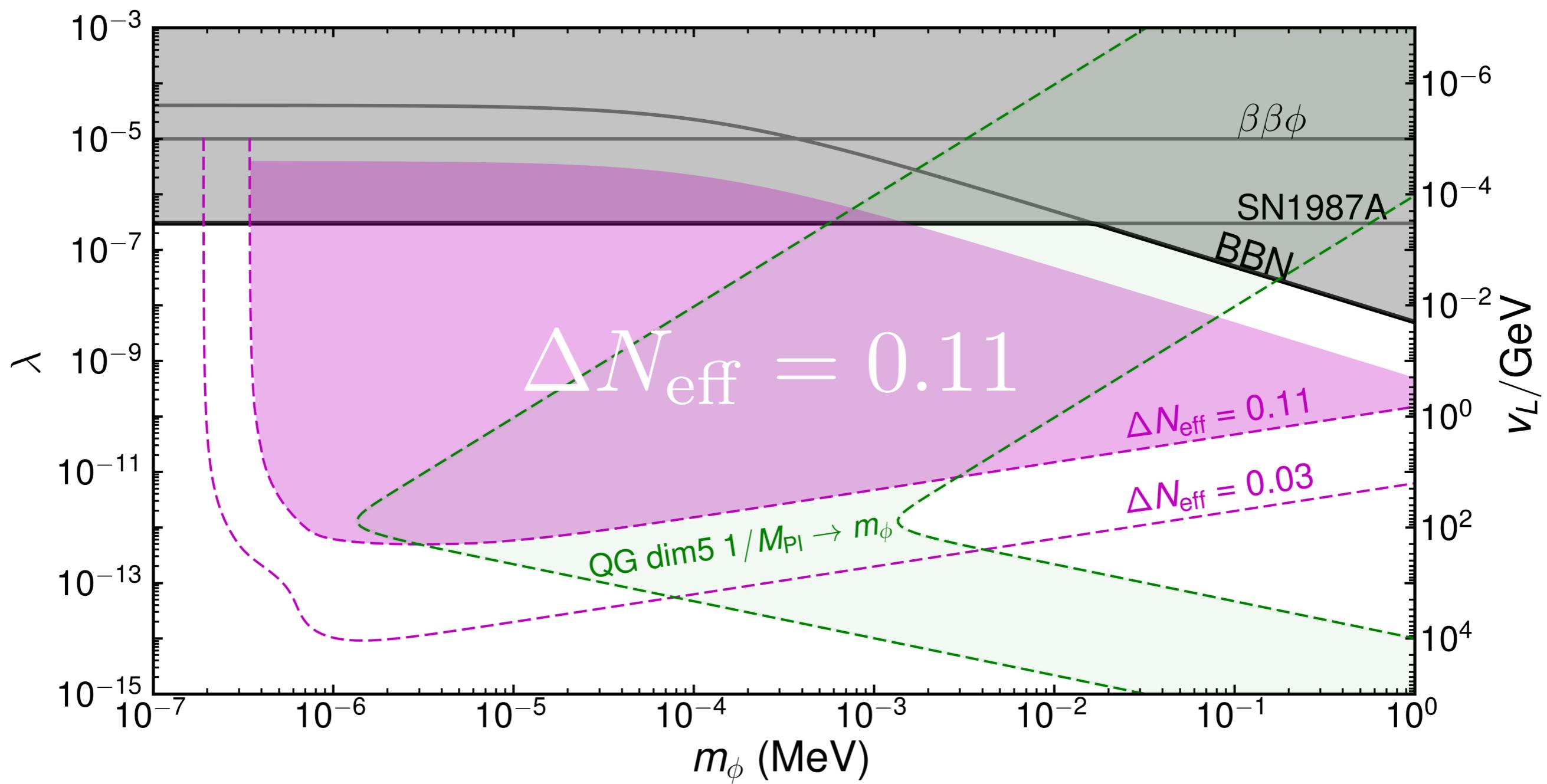
$$\Gamma_{\text{eff}} = \left( \frac{\lambda}{4 \times 10^{-12}} \right)^2 \left( \frac{1 \text{ keV}}{m_\phi} \right)$$



# Parameter Space

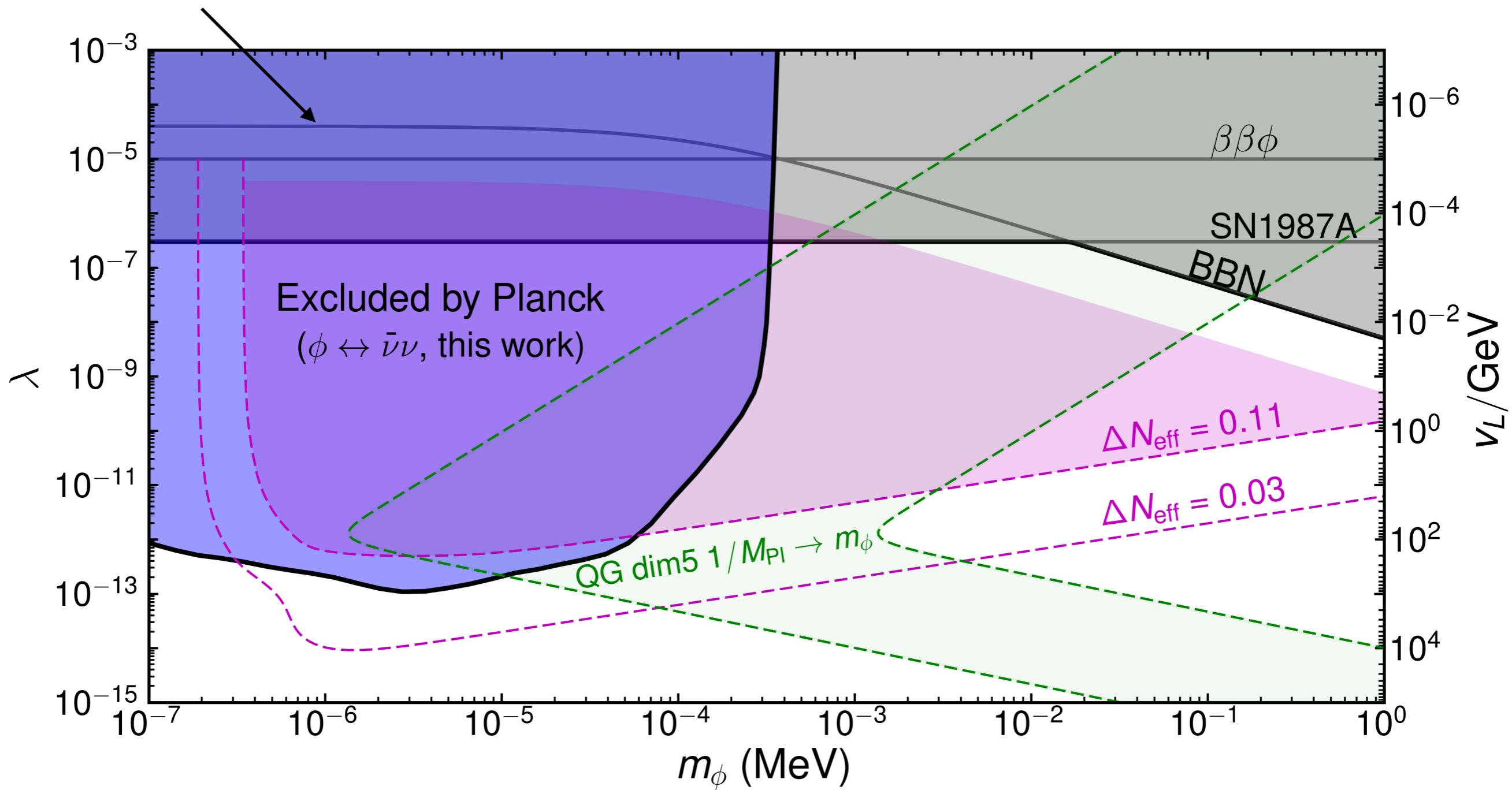


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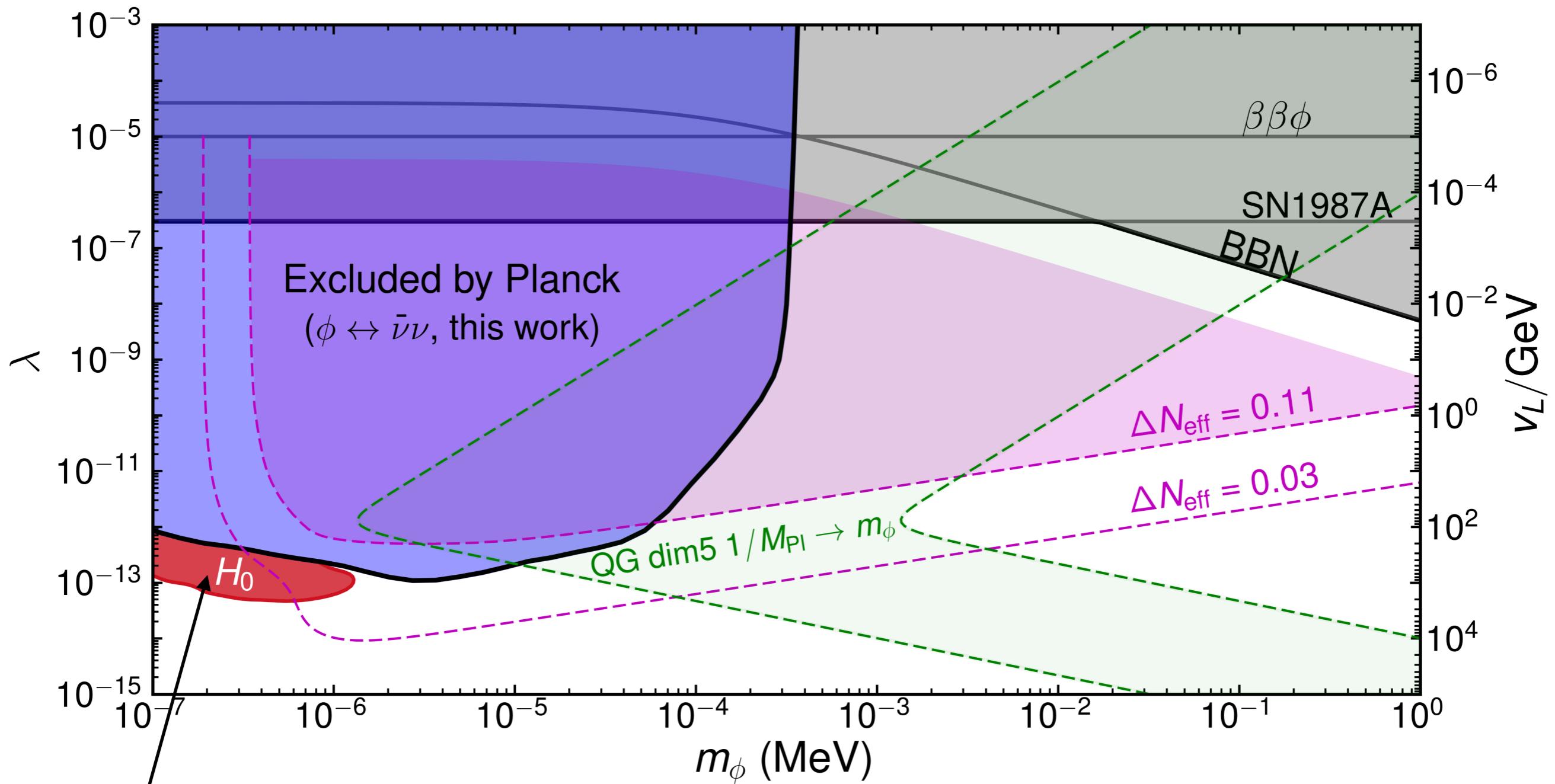


# Parameter Space

Full MCMC to Planck 2018 data

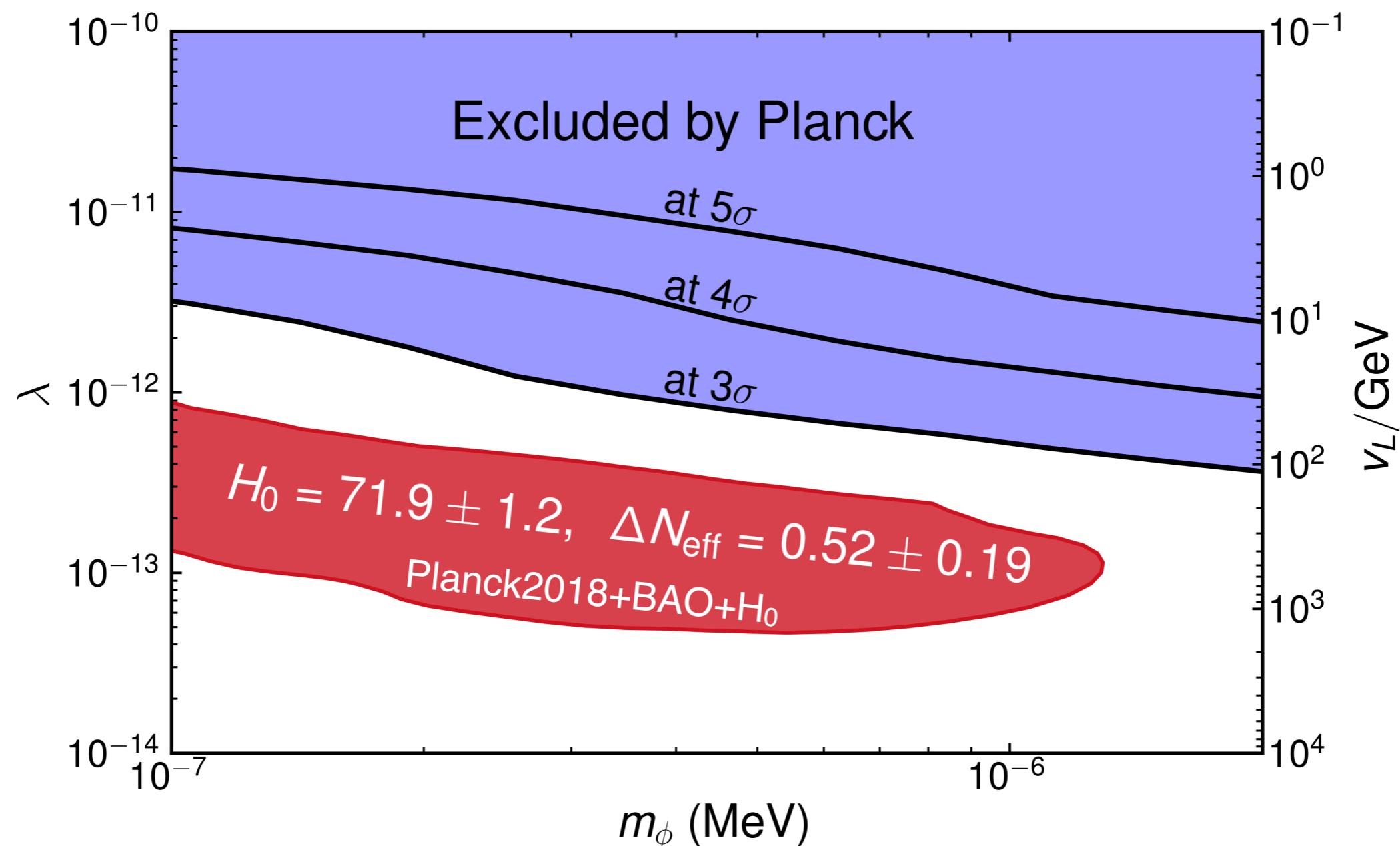


# Parameter Space



**1 $\sigma$  preference when including  $H_0$  in the fit and an additional  $\Delta N_{\text{eff}}$**

# Parameter Space for $H_0$



- Requires a positive  $\Delta N_{\text{eff}} \sim 0.5$
- $H_0$  Thanks to the  $\nu - \phi$  interactions Planck 2018 fit is not degraded wrt  $\Lambda\text{CDM}$
- Very close to the electroweak scale  $v_L \sim (0.1 - 1) \text{ TeV}$

# Summary of 2019

- The Majoron and the Hubble tension
  - Couplings from seesaw and mass from gravity
  - Planck sets very stringent constraints
  - Ameliorates  $H_0$  tension via  $\Delta N_{\text{eff}} = 0.11$
  - CMB S4 experiments will test this region  $\sigma(N_{\text{eff}}^{\text{CMB}}) \sim 0.03$
  - May significantly reduce the tension if:
    - $m_\phi \sim (0.1 - 1) \text{ eV}$
    - $v_L \sim (0.1 - 1) \text{ TeV}$
    - $\Delta N_{\text{eff}} \sim 0.5$
  - 😊  $\Delta N_{\text{eff}} \sim 0.5$  is somewhat ad hoc
  - 😊 Now we have a very good reason for it!

# Leptogenesis with the Majoron

## ● The Majoron and Baryogenesis

- Electroweak Baryogenesis: Cohen, Kaplan, Nelson 90'-91'
- Thermal Leptogenesis: Aristizabal-Sierra, Tortola, Valle, Vicente 1405.4706
- Resonant Leptogenesis: Pilaftsis 0805.1677
- ARS Leptogenesis: Caputo, Hernández & Rius 1807.03309

Let's stick with ARS because since  $\nu_L \sim \nu_H$  we can easily get  $M_N \sim \text{GeV}$

Conclusion from 1807.03309 is that:  $\nu_L > 10^5 \text{ GeV}$  for successful ARS Leptogenesis

I will argue that that is not necessarily true!

I will further argue that N decays can lead to a primordial  $\Delta N_{\text{eff}}$

# ARS Leptogenesis

## ● Baryogenesis via Sterile Neutrino Oscillations

Akhmedov, Rubakov & Smirnov, hep-ph/9803255.

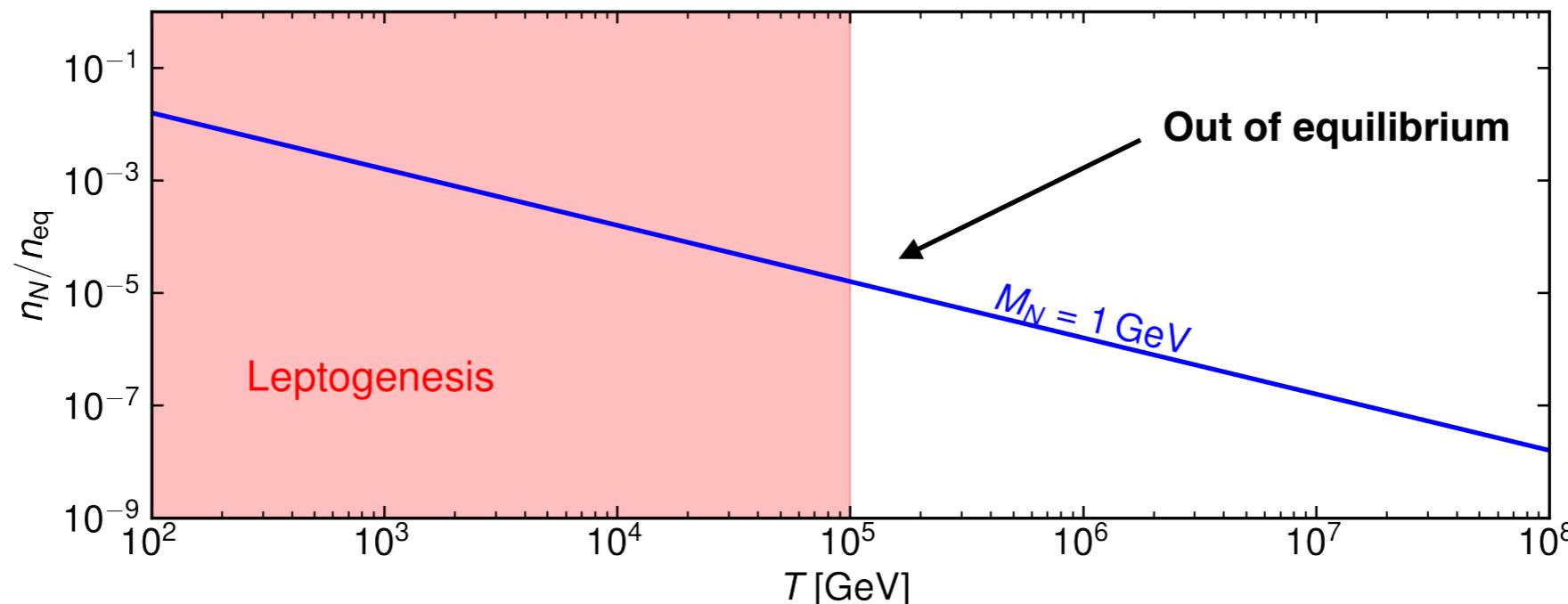
See also Asaka & Shaposhnikov, hep-ph/0505013, Shaposhnikov 0804.4542

Nice review by Hernández et al. 1711.02862

## ● The Idea & key Ingredients:

1) Reheating leads to only SM particles, i.e.  $n_N(T = \infty) = 0$

2) Sterile Neutrinos are slowly produced via their Yukawas, e.g.  $Q_3 \bar{t} \rightarrow \ell N$



# ARS Leptogenesis

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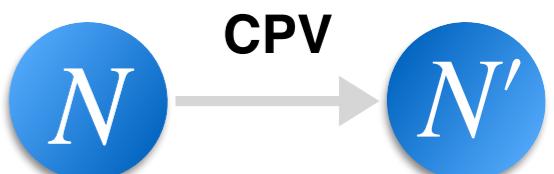
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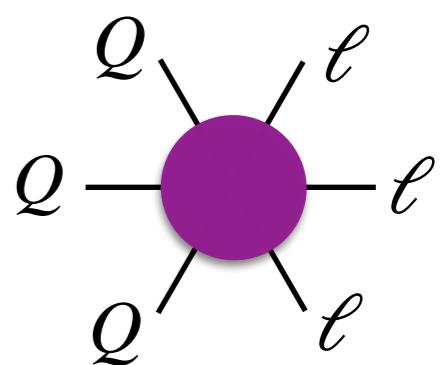
- 1) Reheating leads to only SM particles, i.e.  $n_N(T = \infty) = 0$
- 2) Sterile Neutrinos are slowly produced via their Yukawas, e.g.  $Q_3 \bar{t} \rightarrow \ell N$
- 3) These sterile neutrinos undergo CP violating oscillations

Majorana masses are small and lepton number is conserved

However, via the interactions with leptons an effective asymmetry is produced in the standard model because  $L = L_e + L_\mu + L_\tau + L_{N_1} + L_{N_2} = 0$



- 4) Baryogenesis occurs due to sphaleron freeze-out at  $T_{\text{SPH}} \sim 130 \text{ GeV}$



## ● Parameter Space $n = 2$ :

see e.g. Eijima, Shaposhnikov & Timiryasov, 1808.10833

$0.1 \text{ GeV} \lesssim M_N \lesssim 10 \text{ GeV}$

$\Delta M/M_N \sim 10^{-7} - 10^{-5}$

# ARS Leptogenesis + Majoron

## ● Necessary requirements:

- 1) Sterile Neutrinos cannot be in thermal equilibrium prior to  $T_{\text{SPH}} \sim 130 \text{ GeV}$
- 2) Sterile Neutrinos need to oscillate

$$T_{\text{Lepto}} \sim 0.17 (M_N \Delta M M_{\text{Pl}})^{1/3} \sim 2 \times 10^4 \text{ GeV} \left( \frac{M_N}{10 \text{ GeV}} \right)^{2/3} \left( \frac{\Delta M}{10^{-6} M_N} \right)^{1/3}$$

## ● Threats in the Majoron model:

### Threat to 1)

$$\phi\phi \leftrightarrow N\bar{N}$$

$$\rho\rho \leftrightarrow N\bar{N}$$

$$\rho \leftrightarrow \bar{N}N$$

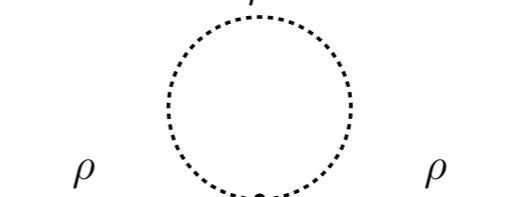
- If  $\rho$ 's are present after reheating, then:  $v_L > 10^5 \text{ GeV}$  Caputo, Hernández & Rius 1807.03309
- Not a problem provided if  $\rho$ 's and  $\phi$ 's are not populated during reheating
- There are no  $\rho$ 's if  $\lambda_{\Phi H} < 10^{-7}$  ( $\rho\rho \leftrightarrow HH$ )
- The smallness is respected by quantum corrections  $\beta(\lambda_{\Phi H})^{(1)} \propto \lambda_{\Phi H}$   $\beta(\lambda_{\Phi H})^{(2)} \propto h_N^2 \lambda_N^2$

### Threat to 2) Symmetry could be unbroken!

$$M_N = \lambda_N \langle \Phi \rangle_\rho$$

- It is fine provided:

$$|\lambda_{\Phi H}| < 10^{-7} \frac{v_L}{1 \text{ TeV}} \sqrt{\frac{10^5 \text{ GeV}}{T_c}}$$



$$m_\rho^2(T) \sim \frac{n_\rho}{n_\rho^{\text{eq}}} \lambda_\Phi \Phi^2$$

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## ● Conclusion

**ARS Leptogenesis is compatible with the singlet Majoron model!**

This works for low  $\nu_L$  provided:

No  $\rho$ 's and  $\phi$ 's are produced during reheating

$$|\lambda_{\Phi H}| < 10^{-7} \frac{\nu_L}{1 \text{ TeV}} \sqrt{\frac{10^5 \text{ GeV}}{T_c}}$$

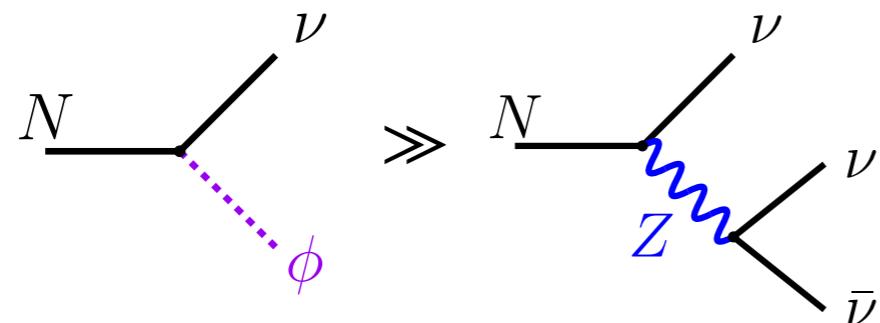
(which is stable under RGE flow)

# Production of Majoron population

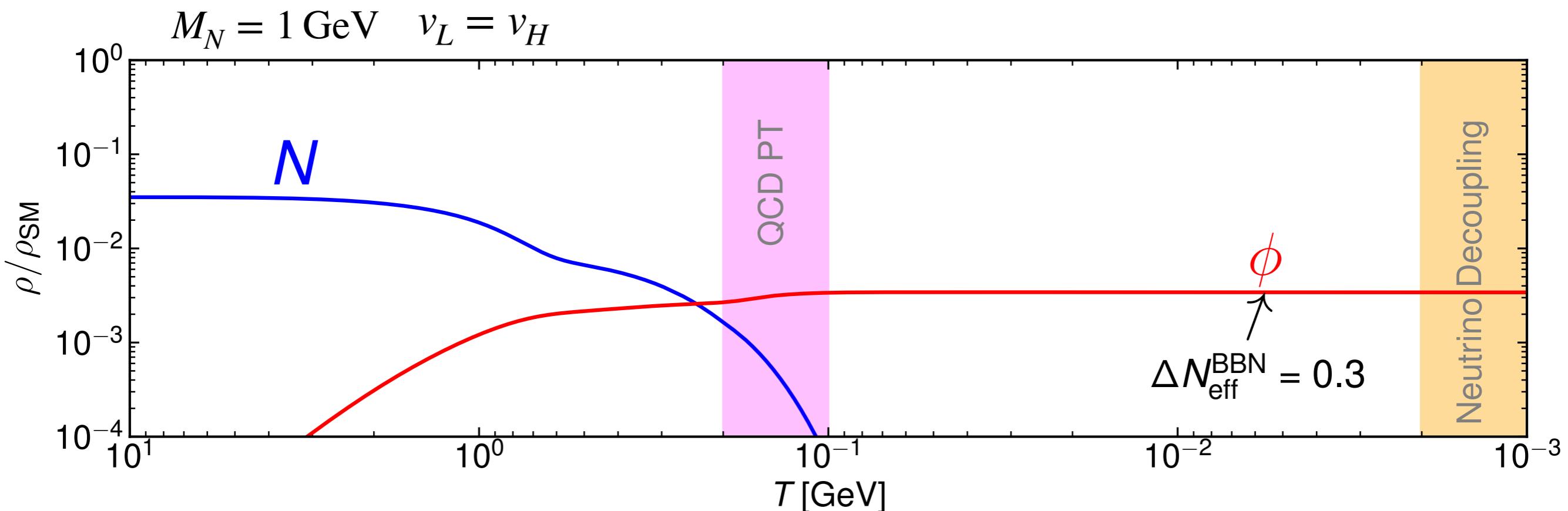
## ● Evolution after Baryogenesis

Sterile Neutrinos that give mass to the active neutrinos thermalize    Ghiglieri & Laine, 1605.07720

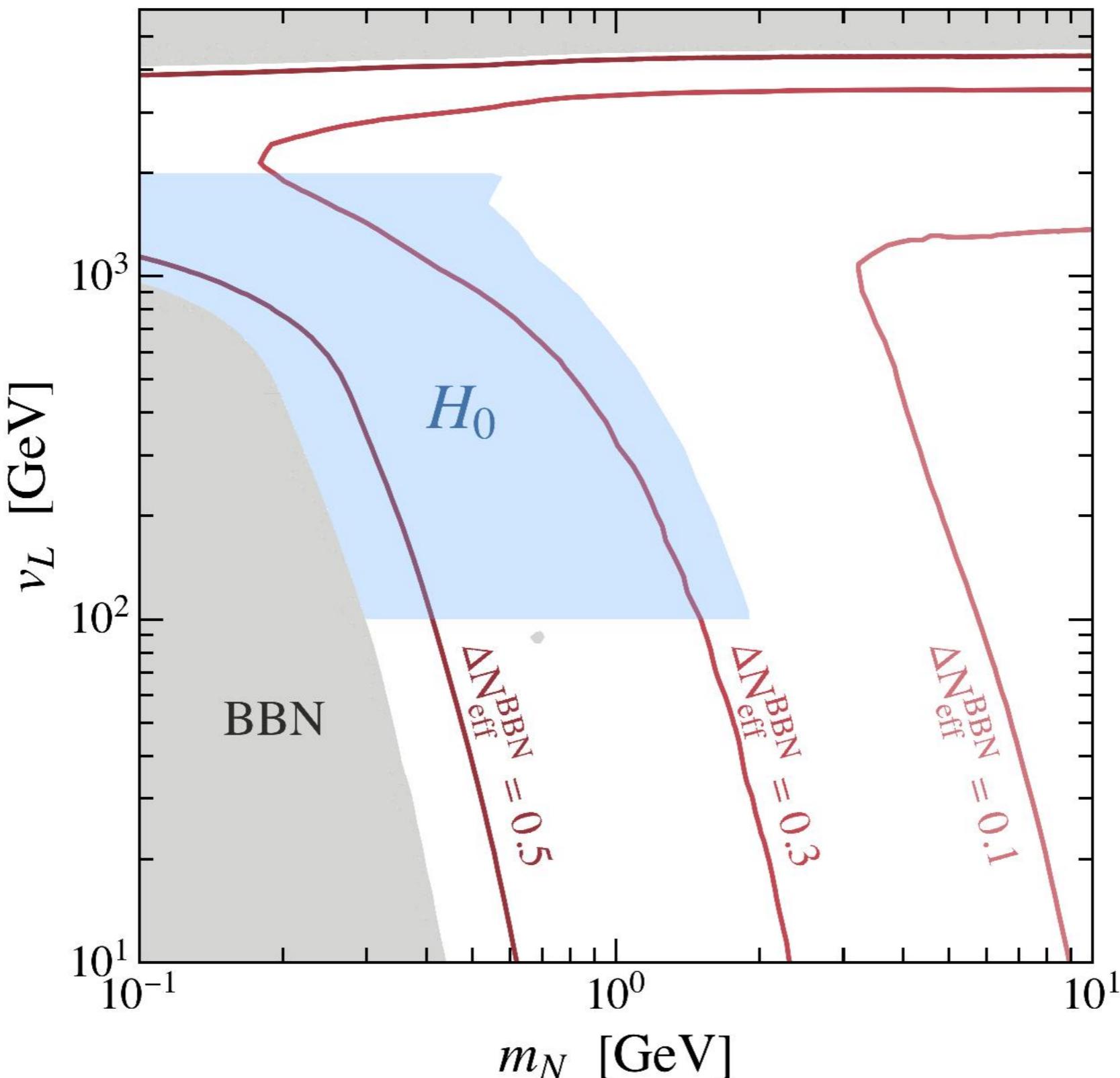
In the majoron model sterile neutrinos have a new decay mode



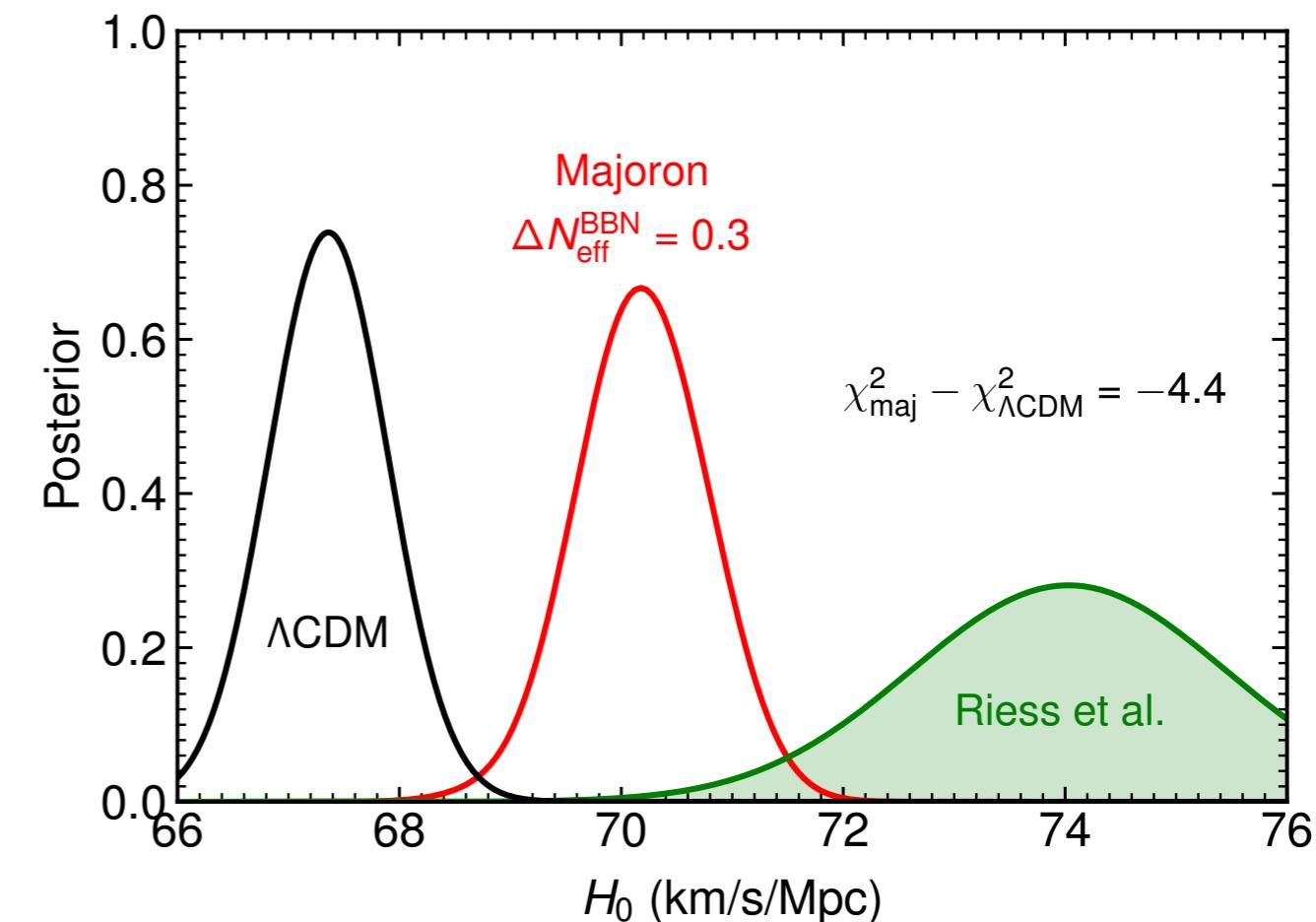
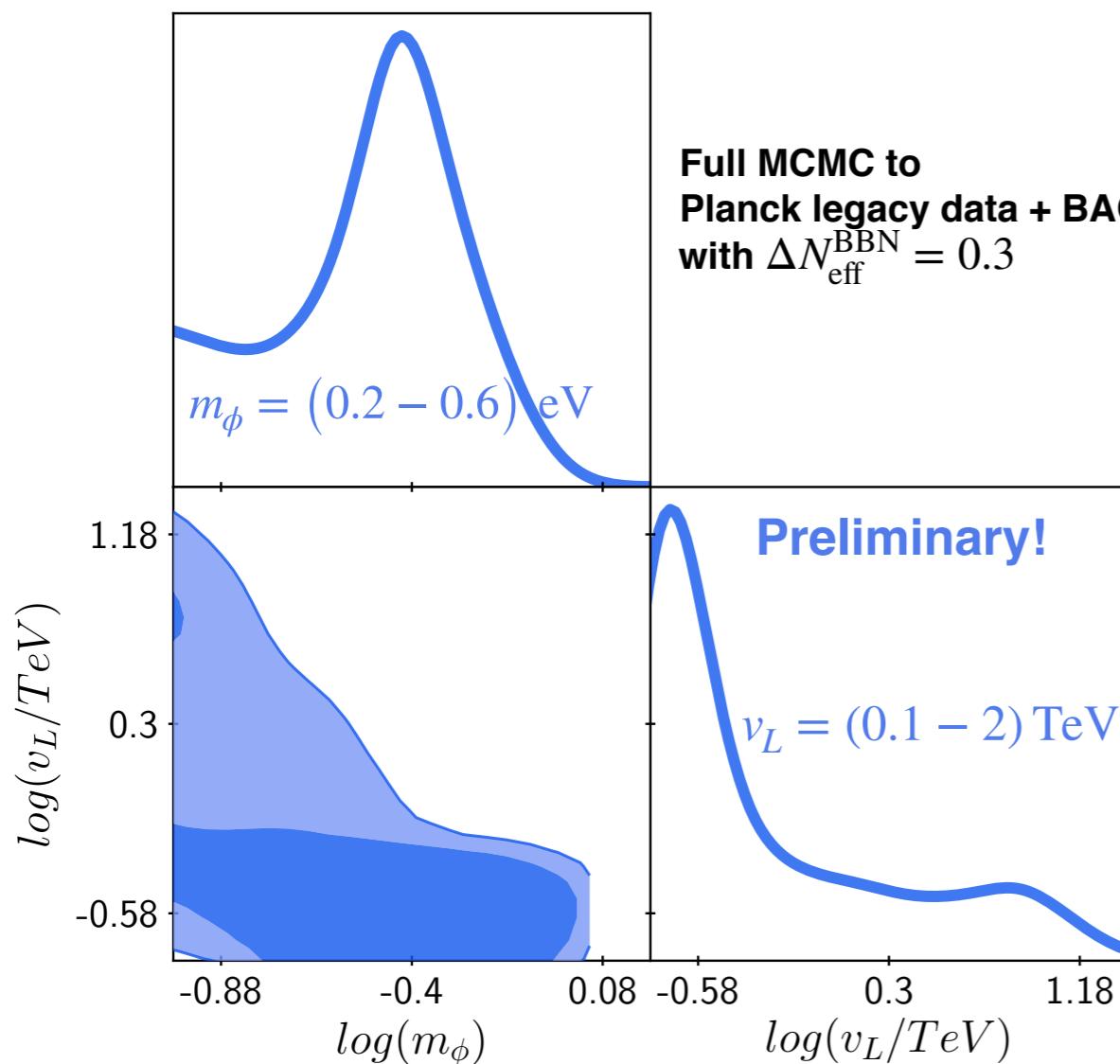
$$\frac{\Gamma(N \rightarrow \nu\phi)}{\Gamma(N \rightarrow \text{SM})} \simeq 4 \times 10^3 \left( \frac{1 \text{ GeV}}{M_N} \right)^2 \left( \frac{1 \text{ TeV}}{v_L} \right)^2$$



# Parameter Space



# Key Results



- **$H_0$  tension down to  $2.5\sigma$**

$$m_\phi = (0.2 - 0.6) \text{ eV}$$

$$v_L = (0.1 - 2) \text{ TeV}$$

- **Excellent fit to Planck data**

$$\Delta\chi^2 = \chi_{\text{maj}}^2 - \chi_{\Lambda\text{CDM}}^2$$

$$\Delta\chi^2 = -4.4$$

# Summary of Escudero & Witte 21'

- ARS can work in the Majoron model for  $\nu_L > 1 \text{ GeV}$

provided  $|\lambda_{\phi H}| < 10^{-7} \frac{\nu_L}{1 \text{ TeV}} \sqrt{\frac{10^5 \text{ GeV}}{T_c}}$  which is protected under RGE flow

- Sterile Neutrinos can provide just the right primordial majoron population
- A full Planck Legacy data analysis shows that:

for  $\Delta N_{\text{eff}}^{\text{BBN}} = 0.3$

$$m_\phi = (0.2 - 0.6) \text{ eV}$$

$$\nu_L = (0.1 - 2) \text{ TeV}$$

$$H_0 = (70.2 \pm 0.6) \text{ km/s/Mpc}$$

$$M_N = (0.2 - 2) \text{ GeV}$$

This makes the tension  $4.4\sigma \rightarrow 2.5\sigma$  but with a better CMB fit than  $\Lambda\text{CDM}$ !

# Conclusions

- **The Hubble Tension:**

- Could be pointing to physics beyond  $\Lambda$ CDM

- There are several neutrino related approaches

- Further data on the way: Gaia, Strong Lensing, GW, DESI ...

- **The Majoron**

- Can substantially relax the tension down to the level of  $2.5\sigma$

- Directly linked to the neutrino mass mechanism

- Mass can be understood from planck-scale effects

- Parameter space to solve the tension very well motivated:  $\nu_L \sim \nu_H$

- The solution can be sourced by low-scale Leptogenesis!

# Outlook

- Could the  $\Delta N_{\text{eff}}^{\text{BBN}} = 0.5$  case fully solve the tension?  
We are investigating this at the moment!
- Including Neutrino masses in the analysis  
This can be relevant to the  $S_8$  tension  
In preparation, Witte & Poulin
- Performing an ARS Leptogenesis calculation  
Never performed in the literature  
Regardless of what happens to  $H_0$ ,  
 $\Delta N_{\text{eff}}$  can be a probe of Leptogenesis!  
In preparation, Ciscar, Escudero & Ibarra

# Acknowledgements

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**Sam Witte!**



# Time for Questions and Comments

**Thank you for your attention!**

