# Sterile Neutrinos at the LHC in Effective Field Theory

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arXiv: 2010.07305,

with J. de Vries, H. Dreiner, J. Günther, Z. Wang



- Sterile neutrinos and effective filed theory
- Displaced vertices search at the LHC
- Numerical results
- Conclusion

#### Sterile Neutrinos in Effective Field Theory

$$\mathcal{L} = \mathcal{L}_{SM} - \left[ \frac{1}{2} \bar{\nu}_R^c \, \bar{M}_R \nu_R + \bar{L} \tilde{H} Y_\nu \nu_R + \text{h.c.} \right] \\ + \mathcal{L}_{\nu_L}^{(\bar{5})} + \mathcal{L}_{\nu_R}^{(\bar{5})} + \mathcal{L}_{\nu_L}^{(\bar{6})} + \mathcal{L}_{\nu_R}^{(\bar{6})} + \mathcal{L}_{\nu_L}^{(\bar{7})} + \mathcal{L}_{\nu_R}^{(\bar{7})} ,$$

- Explain the neutrino oscillation, dark matter and the asymmetry between matter and antimatter.
- Model independent

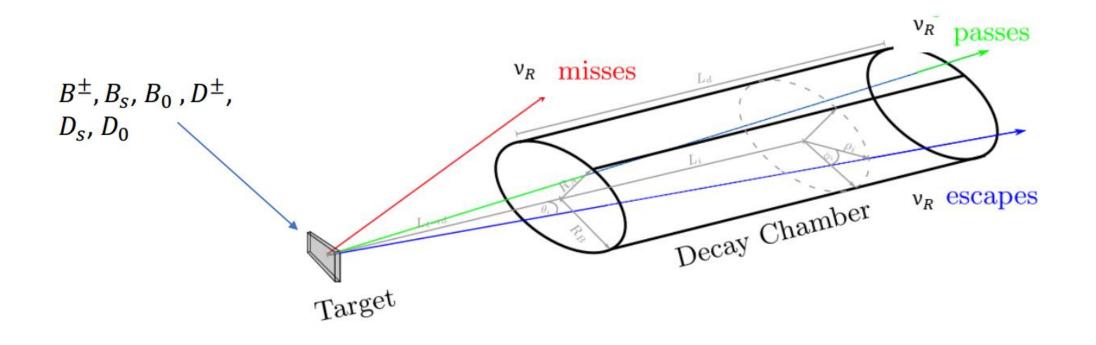
Class 1	$\psi^2 H^3$	Class 4	$\psi^4$
$\mathcal{O}_{L uH}^{(6)}$	$(\bar{L}\nu_R)\tilde{H}(H^{\dagger}H)$	$\mathcal{O}_{du\nu e}^{(6)}$	$(\bar{d}\gamma^{\mu}u)(\bar{\nu}_{R}\gamma_{\mu}e)$
Class 2	$\psi^2 H^2 D$	$\int \mathcal{O}_{Qu\nu L}^{(6)}$	$(ar{Q}u)(ar{ u}_R L)$
$\mathcal{O}_{H u e}^{(6)}$	$(\bar{\nu}_R \gamma^\mu e) (\tilde{H}^\dagger i D_\mu H)$	$\mathcal{O}_{L uQd}^{(6)}$	$(\bar{L}\nu_R)\epsilon(\bar{Q}d))$
Class 3	$\psi^2 H^3 D$	$\mathcal{O}_{LdQ u}^{(6)}$	$(\bar{L}d)\epsilon(\bar{Q} u_R)$
$\mathcal{O}^{(6)}_{ uW}$	$(\bar{L}\sigma_{\mu\nu}\nu_R)\tau^I\tilde{H}W^{I\mu\nu}$		

#### Dimension-6 operators with one sterile neutrino

#### After EWSB and rotating to the neutrino mass basis,

$$\mathcal{L}_{\text{mass}}^{(6,7)} = \frac{2G_F}{\sqrt{2}} \Biggl\{ \bar{u}_L \gamma^{\mu} d_L \left[ \bar{e}_L \gamma_{\mu} C_{\text{VLL}}^{(6)} \nu + \bar{e}_R \gamma_{\mu} C_{\text{VLR}}^{(6)} \nu \right] + \bar{u}_R \gamma^{\mu} d_R \bar{e}_R \gamma_{\mu} C_{\text{VRR}}^{(6)} \nu \bar{u}_L d_R \bar{e}_L C_{\text{SRR}}^{(6)} \nu + \bar{u}_R d_L \bar{e}_L C_{\text{SLR}}^{(6)} \nu + \bar{u}_L \sigma^{\mu\nu} d_R \bar{e}_L \sigma_{\mu\nu} C_{\text{TRR}}^{(6)} \nu + \frac{1}{v} \bar{u}_L \gamma^{\mu} d_L \bar{e}_L C_{\text{VLR}}^{(7)} i \overleftrightarrow{D}_{\mu} \nu \Biggr\} + \text{h.c.} ,$$

### **Displaced search for sterile neutrino at the LHC**



from 1511.07436, H. Dreiner, J. de Vries

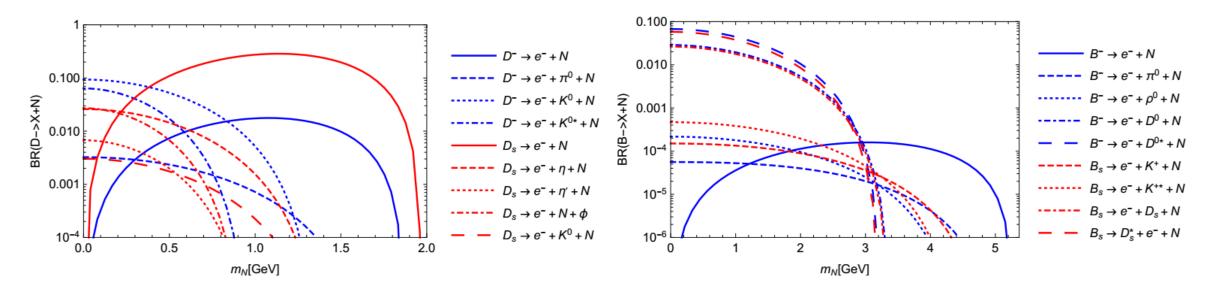
#### **Production of the sterile neutrino:**

(1). *D* and *B* pseudoscalar mesons decay, including leptonic and semileptonic decays.

(2). Direct production with parton collisions, sub-dominant when neutrino is lighter than B meson.

#### Production in the minimal model

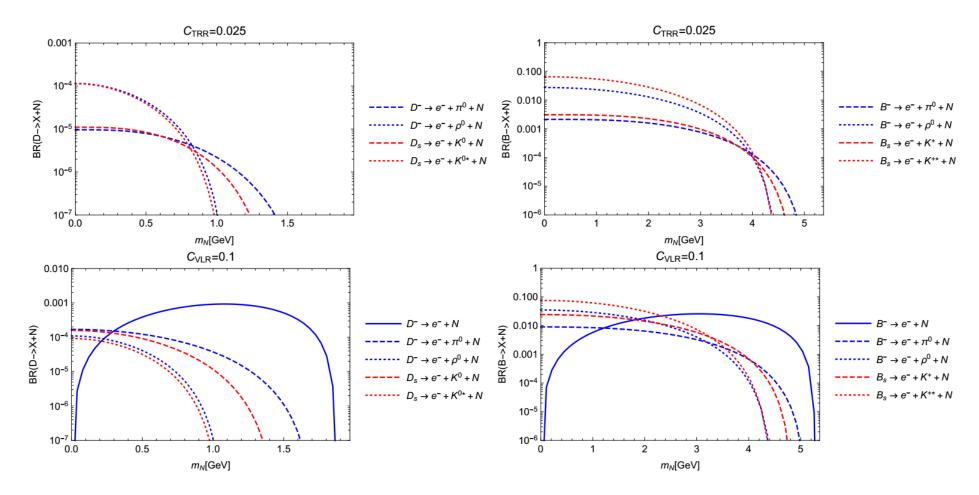
- Only one sterile neutrino
- interact with SM fields via the mixing
- The mixing angle is set to be  $U_{e4} = 1$



### Production from dimension-6 operators only

{ijkl}=21e4

*{ijkl}=13e4* 



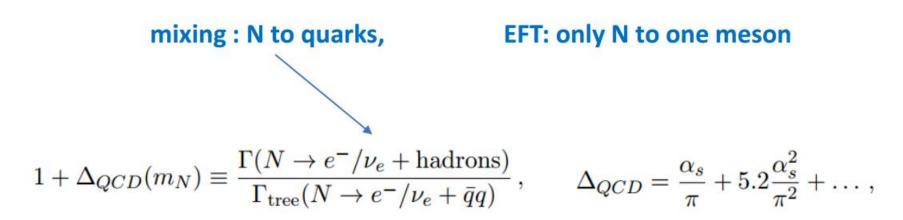
# The decay of sterile neutrino

- **1.** N → leptons (minimal scenario only)
- 2.  $N \rightarrow P/V + e$ ,

P:  $\pi$ , K, D and  $D_s$ . V:  $\rho$ ,  $D^*$ ,  $D_s^*$ ,  $K^*$ 3. N $\rightarrow P^0/V^0 + \nu_e$  (minimal scenario only)

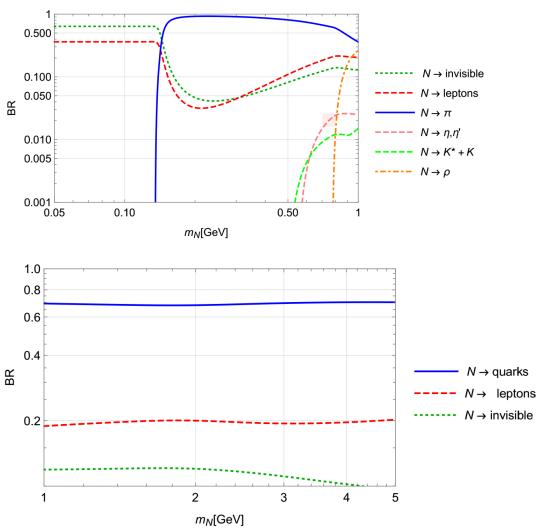
 $P^0$ : π<sup>0</sup>, η, η', η<sup>c</sup>.  $V^0$ : ρ<sup>0</sup>, ω, φ, J/ψ

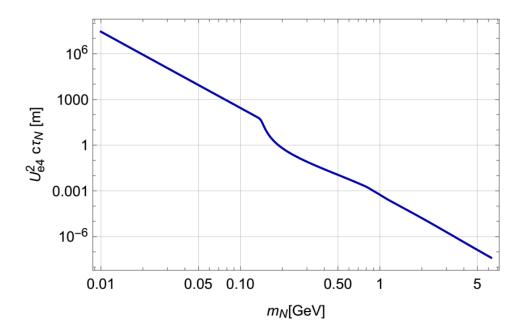
4.  $N \rightarrow 2$  and more mesons



### Decay in the minimal model

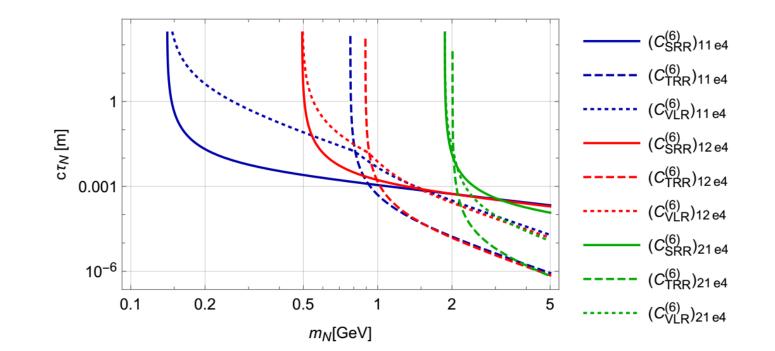
• Sterile neutrinos decay only through the mixing with  $v_e$ .





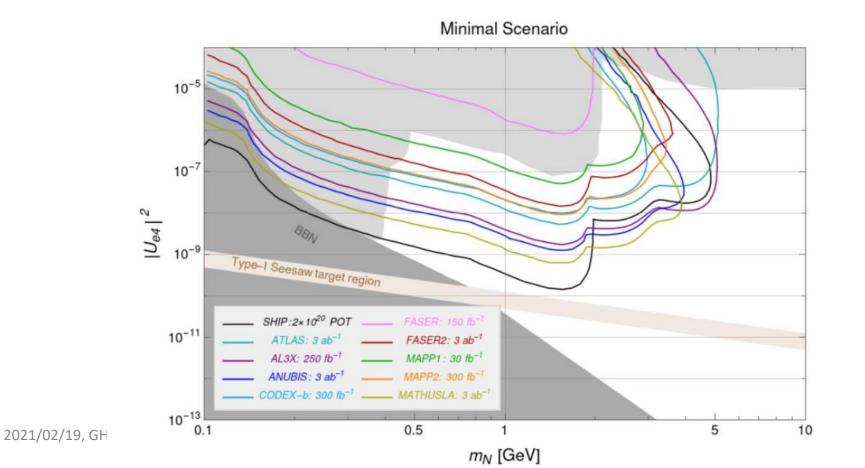
### Decay from dimension-6 operators only

- Turn off the weak interaction
- Turn on only one unit WC at a time
- Each operator can only induce one or two decay channels, N to one meson plus an electron



#### Results in the minimal scenario

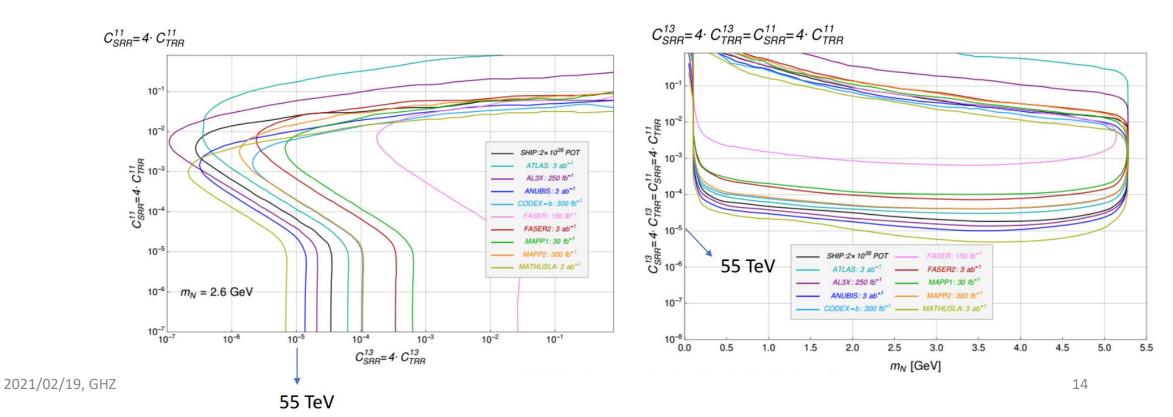
- Switch off higher-dimensional operators
- Consider only one sterile neutrino mixed with  $v_e$  only



See also: 2001.04750 Hirsch, Wang

#### Results in the leptoquark scenario

- It induces two d-6 operators with  $C_{SRR} = 4C_{TRR}$
- Minimal scenario is also included
- The mass of electron neutrino is 0.05 eV



# Conclusion

- We evaluate the sensitivity reach of ATLAS, AL3X, ANUBIS, CODEX-b, FASER, MATHUSLA, MOEDAL\_MAPP, and SHiP
- Searches for displaced vertices of long-lived sterile neutrinos are a good probe of vSMEFT
- SHiP and MATHUSLA can probe scales around 80 TeV, ANUBIS: 55 TeV, FASER2: 25 TeV.
- Our results are competitive with and complementary to the existing limits

### Appendix

 $\langle 0 | \bar{q}_i \gamma^\mu \gamma^5 q_j | M(q) \rangle \equiv i q^\mu f_M \,,$ 

$$\langle M'_P(p') | \bar{q}_1 \gamma^{\mu} q_2 | M_P(p) \rangle = f_+(q^2) \left[ (p+p')^{\mu} - \frac{M^2 - m^2}{q^2} q^{\mu} \right] + f_0(q^2) \frac{M^2 - m^2}{q^2} q^{\mu} , \langle M'_P(p') | \bar{q}_1 q_2 | M_P(p) \rangle = f_S(q^2) , \langle M'_P(p') | \bar{q}_1 \sigma^{\mu\nu} q_2 | M_P(p) \rangle = \frac{2i}{M+m} [ p^{\mu} p'^{\nu} - p^{\nu} p'^{\mu} ] f_T(q^2) ,$$