



The Cosmological Bootstrap

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Based on work with

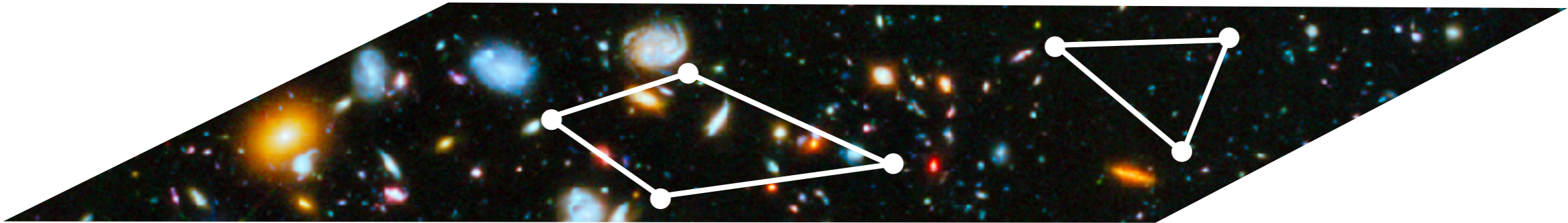
Nima Arkani-Hamed, Wei Ming Chen, Hayden Lee, Manuel Loparco, Guilherme Pimentel, Carlos Duaso Pueyo and Austin Joyce

Interesting related work by

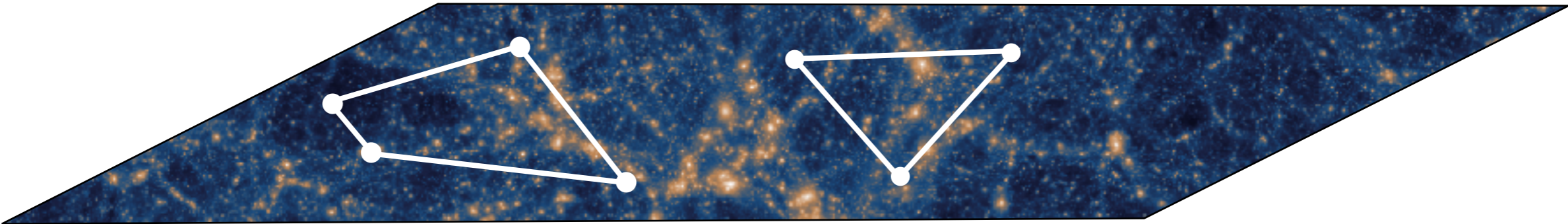
Enrico Pajer, Sadra Jazayeri, David Stefanyszyn, Scott Melville, Harry Goodhew, Paolo Benincasa, Charlotte Sleight, Massimo Taronna

Soner Albayrak, Adam Bzowski, Sebastian Céspedes, Xingang Chen, Claudio Coriano, Anne-Christine Davis, Lorenz Eberhardt, Joseph Farrow, Garrett Goon, Tanguy Grall, Daniel Green, Aaron Hillman, Kurt Hinterbichler, Hiroshi Isono, Alex Kehagias, Savan Kharel, Shota Komatsu, Soubhik Kumar, Arthur Lipstein, Matteo Maglio, Juan Maldacena, Paul McFadden, David Meltzer, Sebastian Mizera, Toshifumi Noumi, Rafael Porto, Suvrat Raju, Antonio Riotto, Gary Shiu, Allic Sivaramakrishnan, Kostas Skenderis, Raman Sundrum, Jakub Supel, Sandip Trivedi, Mark Trodden, Yi Wang, ...

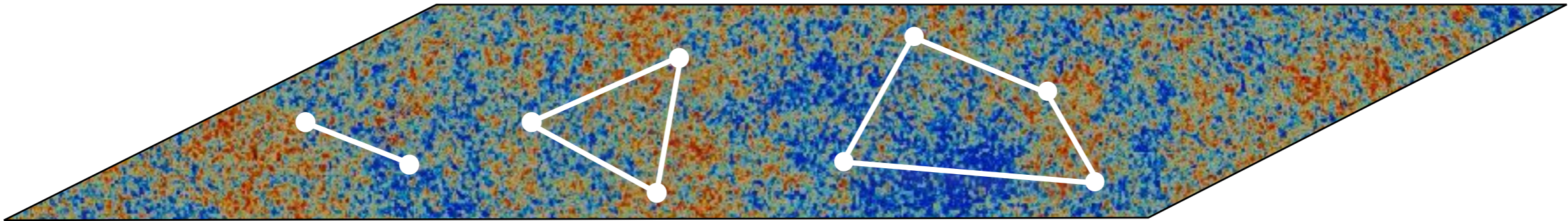
Cosmological structures are not distributed randomly, but display spatial correlations over very large distances:



13.8 billion years

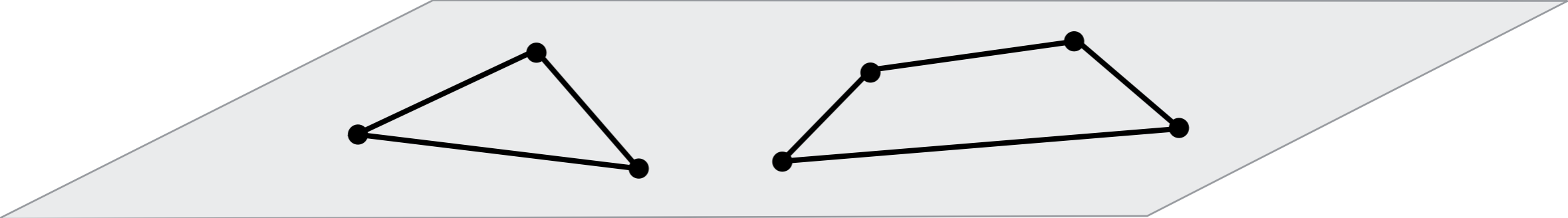
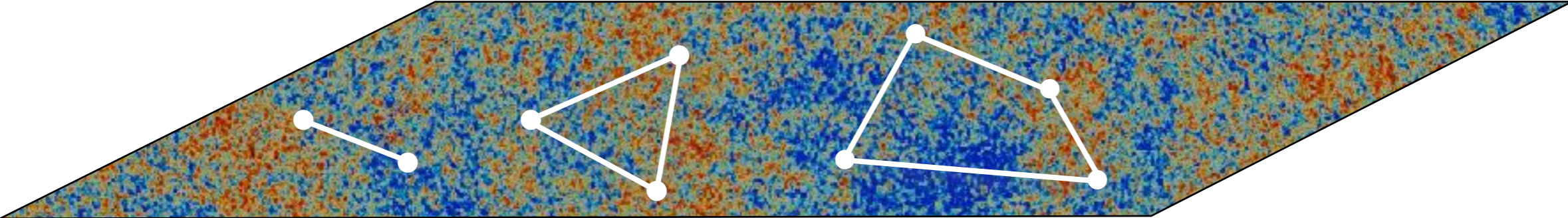


few billion years

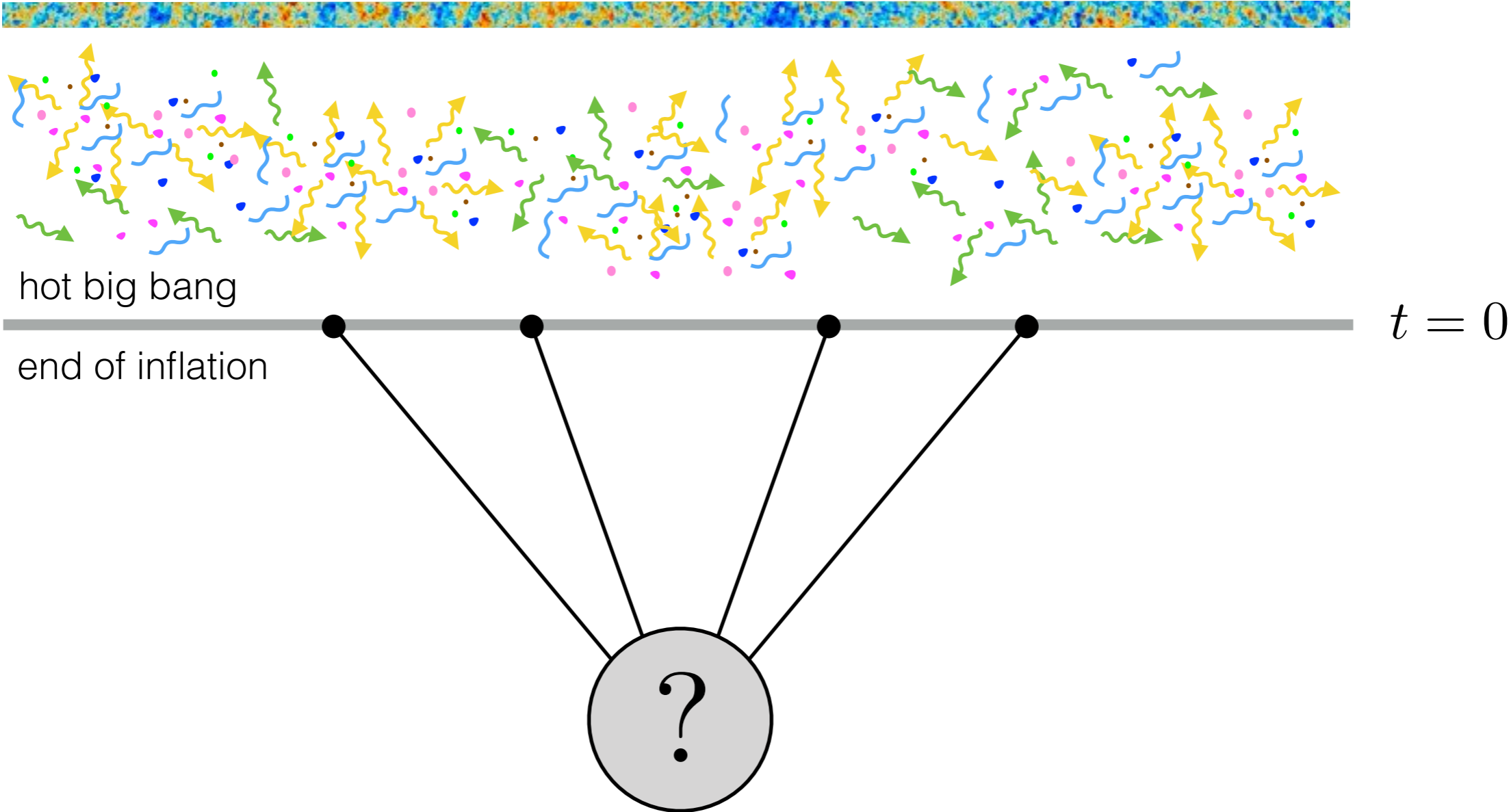


380 000 years

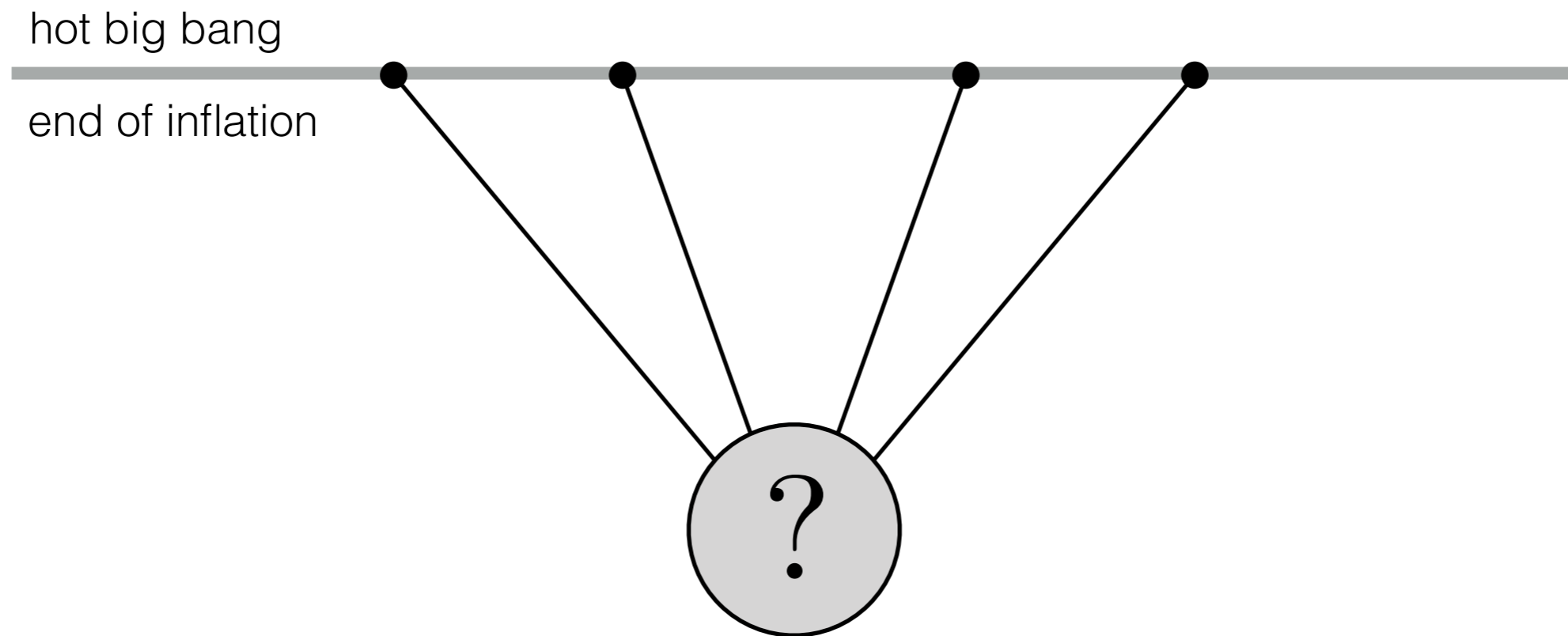
These correlations can be traced back to the beginning of the hot Big Bang:



However, the Big Bang was not the beginning of time, but the end of an earlier high-energy period:

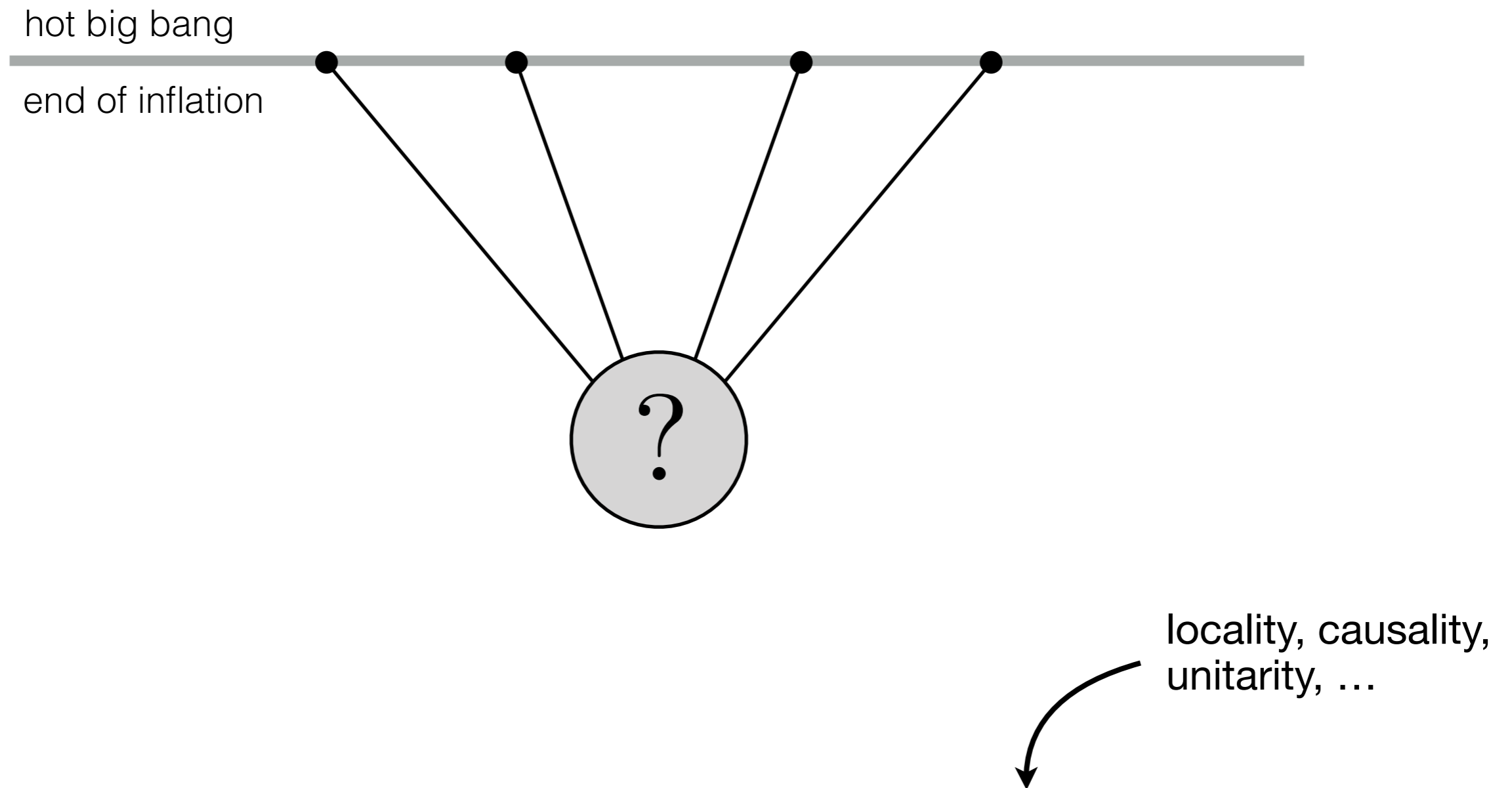


What exactly happened before the hot Big Bang?
How did it create the primordial correlations?



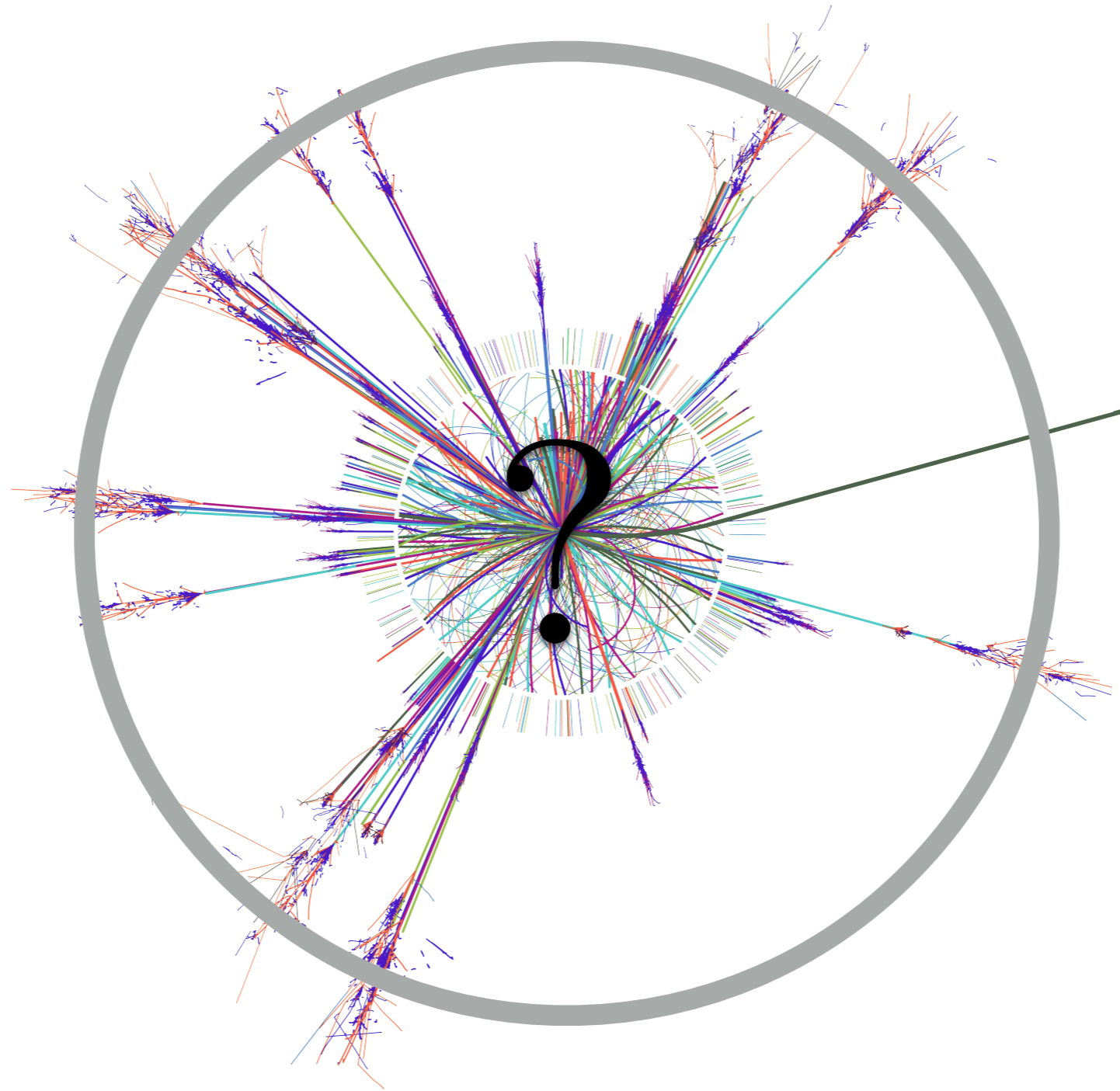
The challenge is to extract this information from the pattern of correlations in the late universe.

What is the space of consistent correlations?



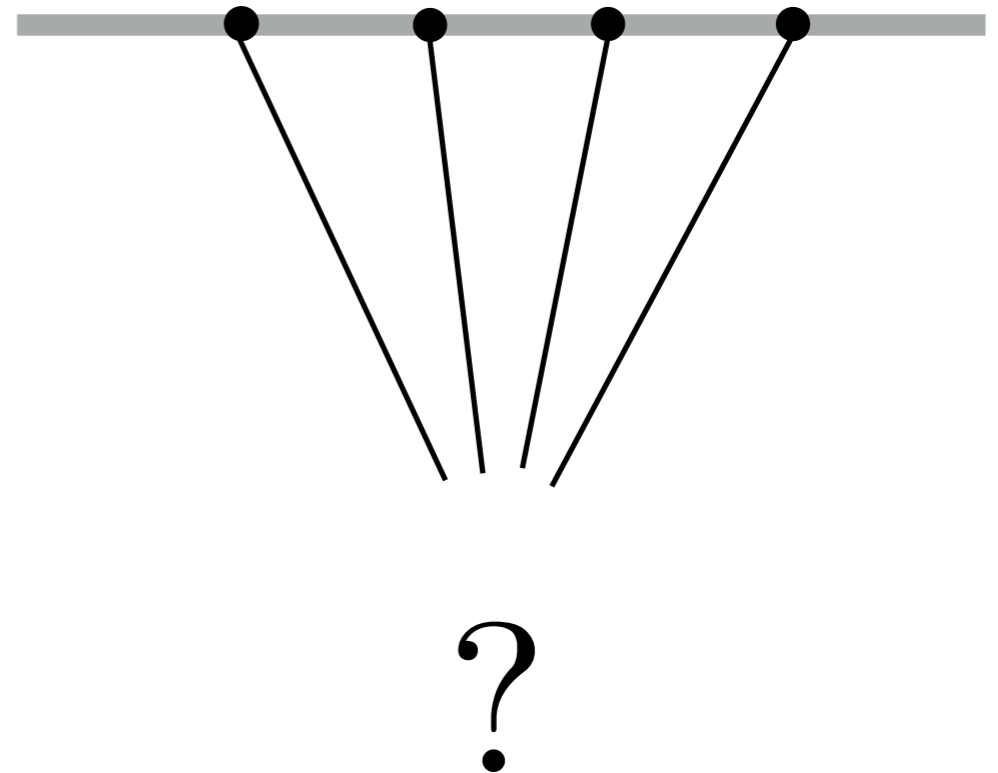
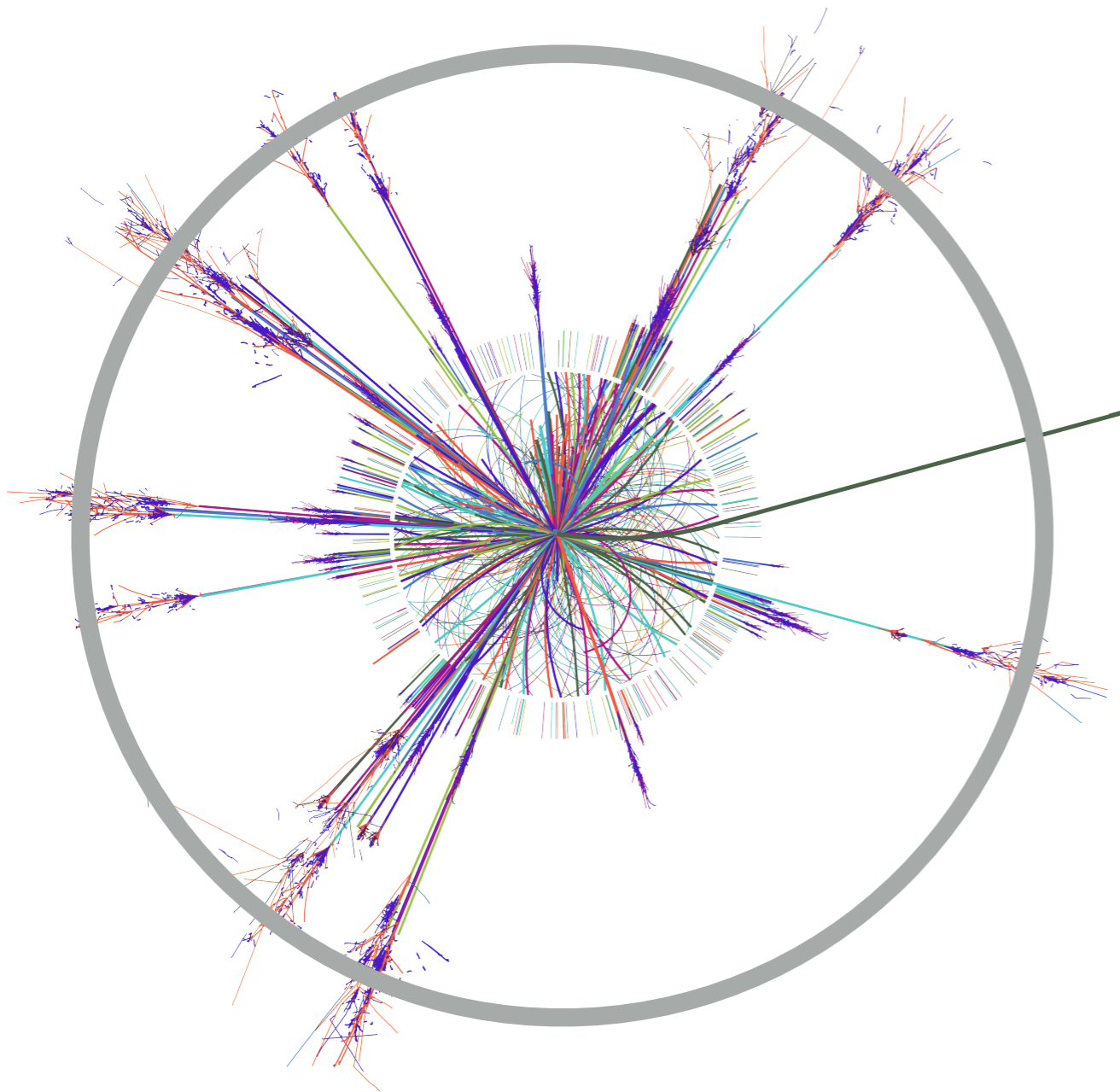
- What correlations are consistent with basic **physical principles**?
- Can these correlations be **bootstrapped** directly?

The bootstrap perspective has been very influential for **scattering amplitudes**:



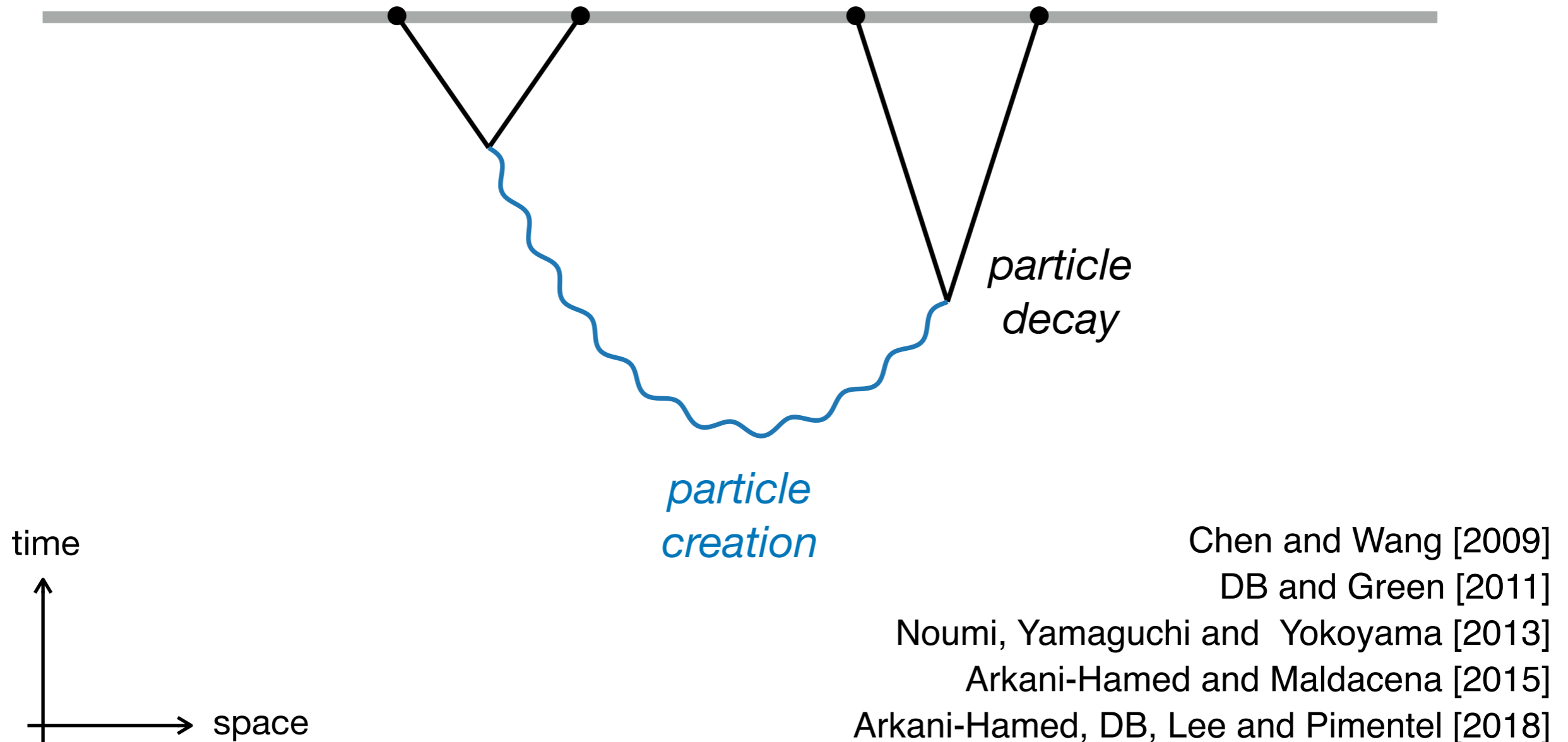
In that case, the rules of **quantum mechanics** and **relativity** are very constraining.

Does a similar **rigidity** exist for cosmological correlators?



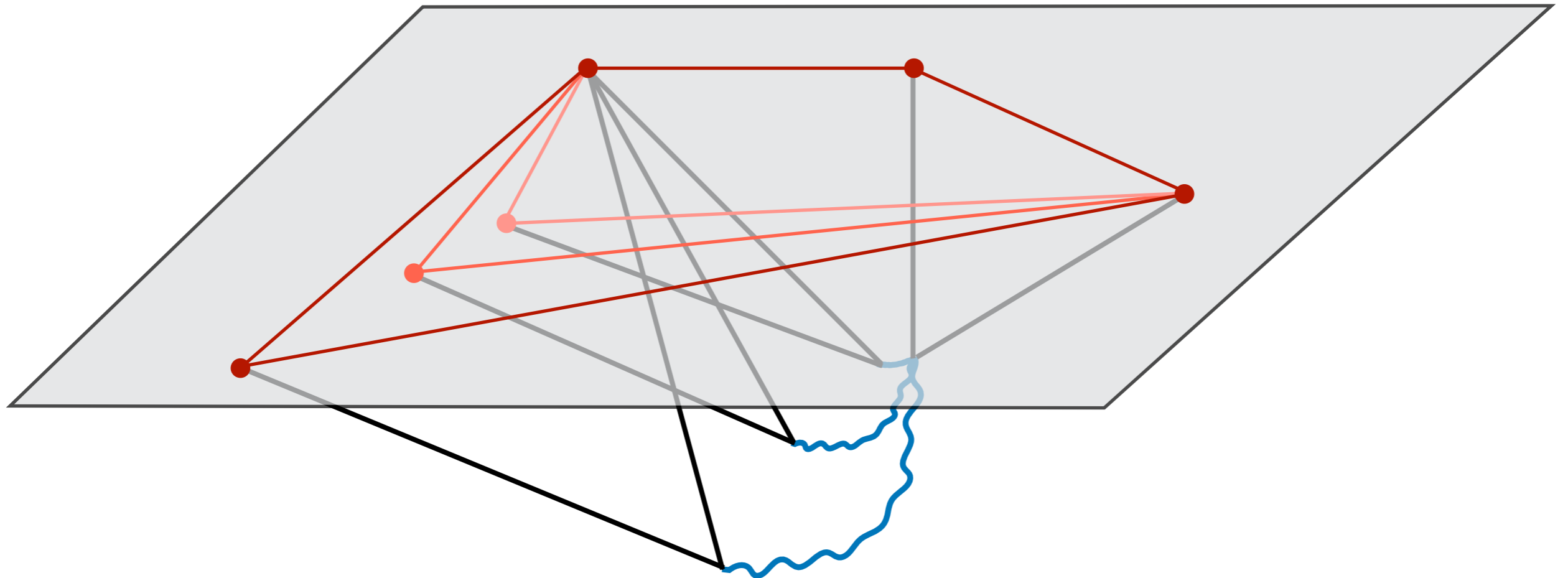
Goal: Develop an understanding of cosmological correlators that parallels our understanding of flat-space scattering amplitudes.

The connection to scattering amplitudes is also relevant because the early universe was like a giant **cosmological collider**:



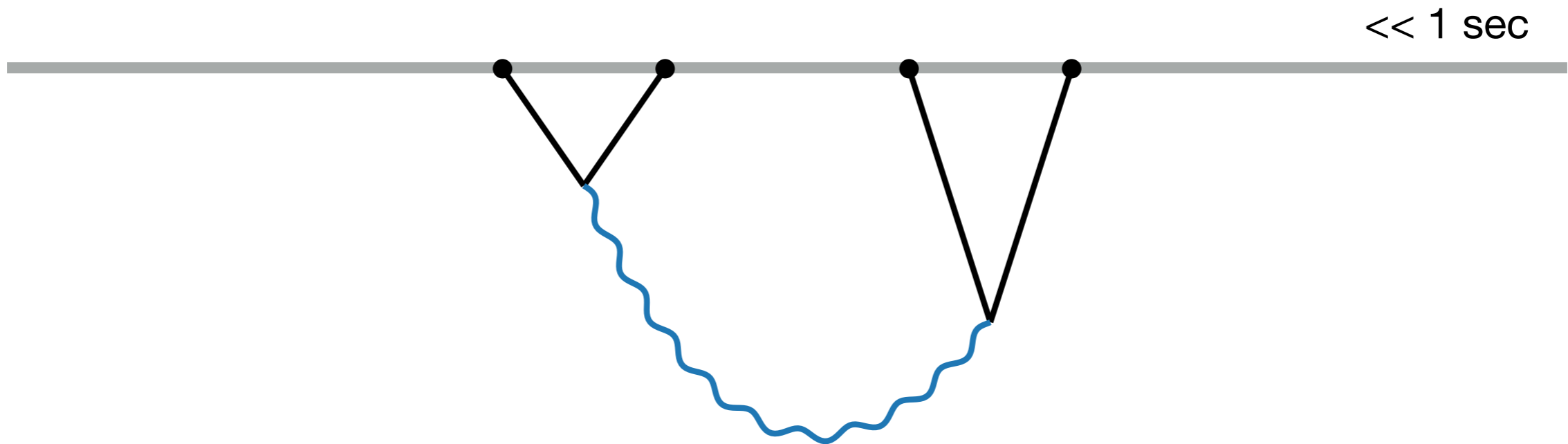
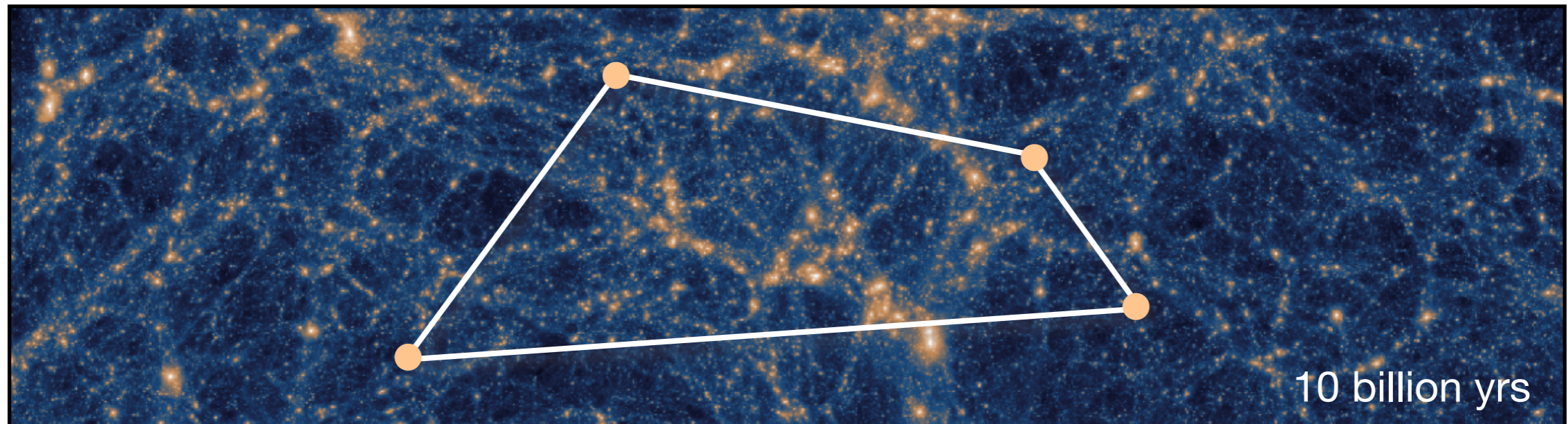
During inflation, the rapid expansion can produce very **massive particles** ($\sim 10^{14}$ GeV) whose decays lead to nontrivial correlations.

These particles are tracers of the inflationary dynamics:



The pattern of correlations *after* inflation contains a memory of the physics *during* inflation (evolution, symmetries, particle content, etc).

At late times, these correlations leave an imprint in the distribution of galaxies:



Goal: Develop a systematic way to predict these signals.

Outline



**Basics of the
Bootstrap**



**Recent
Progress**



**Future
Directions**



Basics of the Bootstrap

The Conventional Approach

How we usually make predictions:

Physical Principles

locality, causality,
unitarity, symmetries



Models

Lagrangians
equations of motion
spacetime evolution
Feynman diagrams



Observables

Works well if we have a well-established theory (Standard Model, GR, ...)
and many observables.

The Conventional Approach

Although conceptually straightforward, this has some drawbacks:

- It involves unphysical gauge degrees of freedom.
- Relation between Lagrangians and observables is many-to-one.
- Even at tree level, the computation of cosmological correlators involves complicated time integrals:

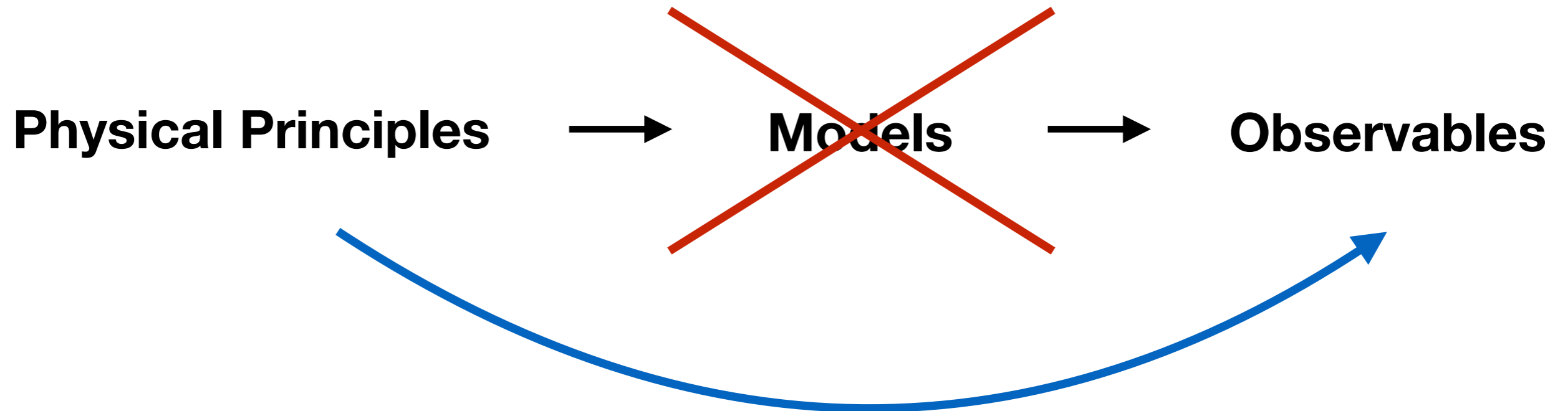
$$F_N = \sum_{\text{Diagrams}} \left(\prod_{\text{Vertices}} dt \right) \left(\text{External propagators} \right) \left(\text{Internal propagators} \right) \left(\text{Vertex factors} \right)$$

Hard to compute!

- Fundamental principles (e.g. locality and unitarity) are obscured.
- To derive nonperturbative correlators and constraints from the UV completion, we need a deeper understanding.

The Bootstrap Method

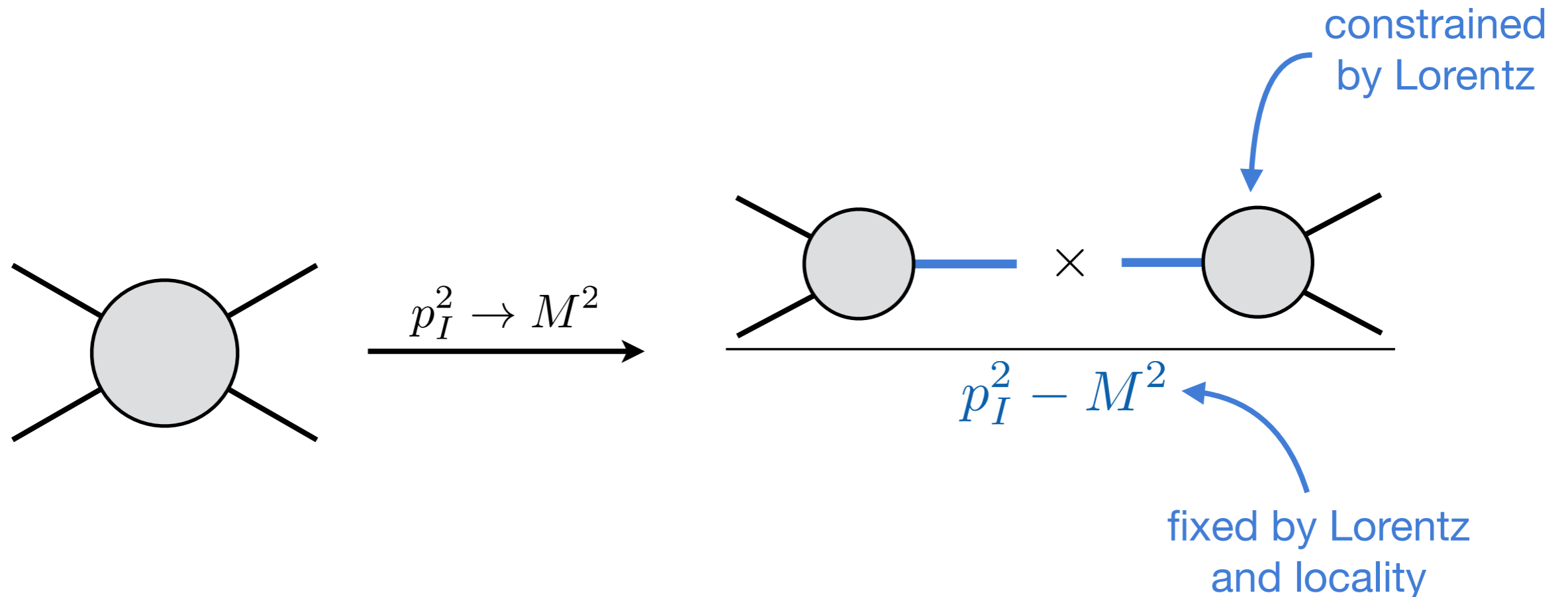
In the bootstrap approach, we cut out the middle man and go directly from fundamental principles to physical observables:



This is particularly relevant when we have many theories (inflation, BSM, ...) and few observables.

S-matrix Bootstrap

Much of the physics of scattering amplitudes is controlled by **SINGULARITIES**:

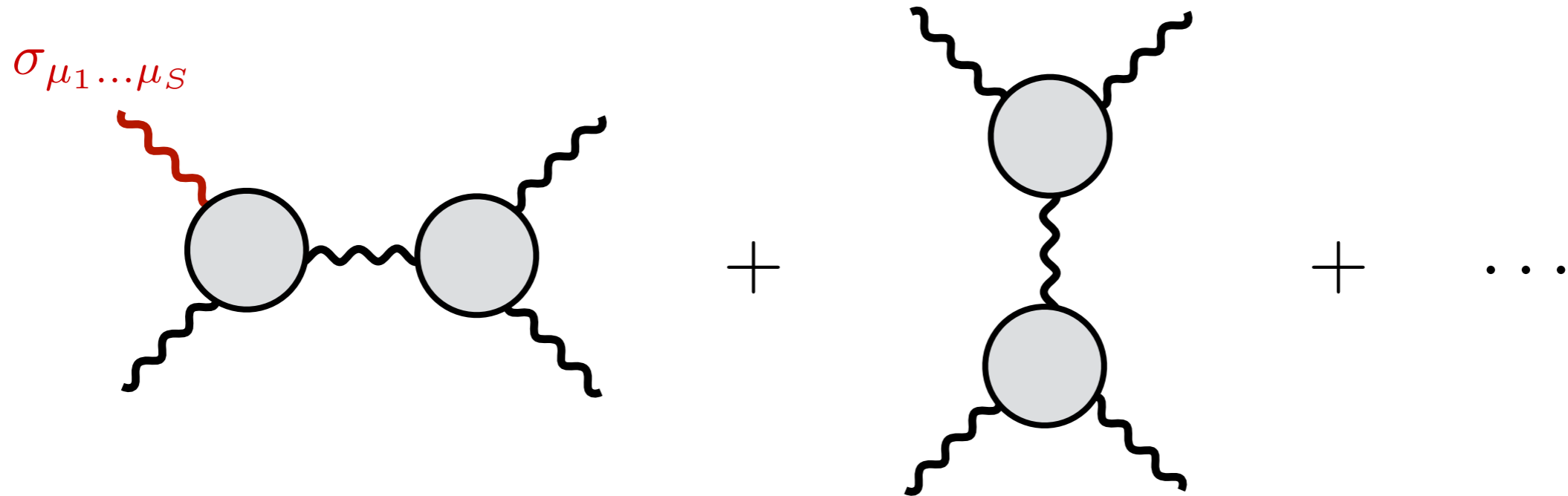


- Amplitudes have **poles** when intermediate particles go on-shell.
- On these poles, the amplitudes **factorize**.

The different singularities are connected by (Lorentz) **SYMMETRY**.

S-matrix Bootstrap

Consistent factorization is very constraining for massless particles:



→ Only consistent for spins

$$S = \left\{ 0, \frac{1}{2}, 1, \frac{3}{2}, 2 \right\}$$

YM SUSY GR

Arrows indicate the correspondence between theories and spin values: an orange arrow points from 'YM' to the value 1; a green arrow points from 'SUSY' to the value 3/2; and a purple arrow points from 'GR' to the value 2.

A Success Story

The modern amplitudes program has been very successful:

1. New Computational Techniques

- Recursion relations
- Generalized unitarity
- Soft theorems

2. New Mathematical Structures

- Positive geometries
- Amplituhedrons
- Associahedrons

3. New Relations Between Theories

- Color-kinematics duality
- BCJ double copy

4. New Constraints on QFTs

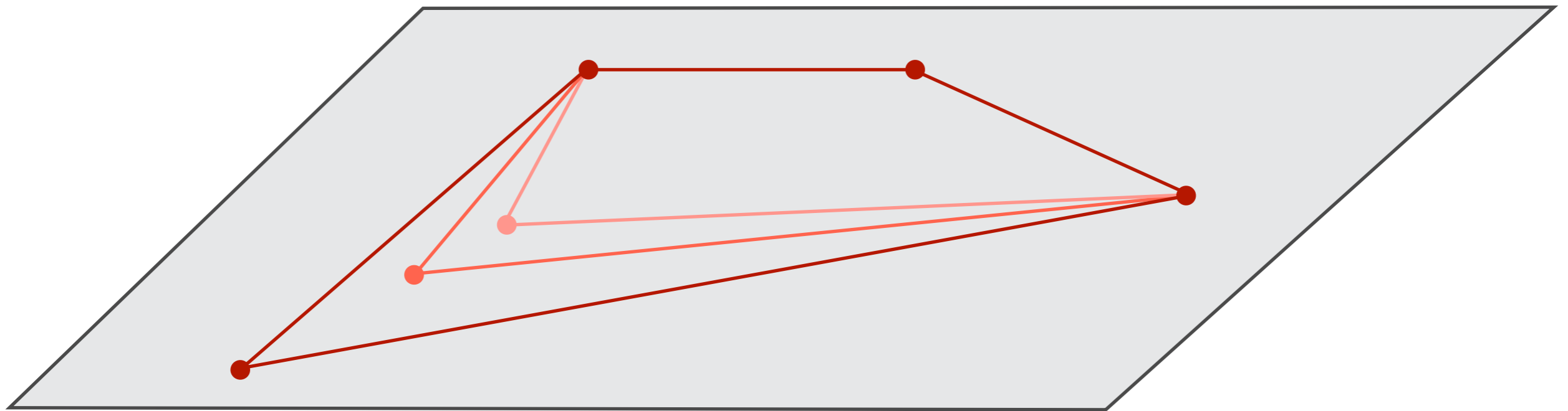
- Positivity bounds
- EFThedron

5. New Applications

- Gravitational wave physics
- Cosmology

Cosmological Bootstrap

If inflation is correct, then all cosmological correlations can be traced back to the future boundary of an approximate **de Sitter spacetime**:



We want to **bootstrap** these boundary correlations directly, without explicit reference to the bulk dynamics.

Cosmological Bootstrap

The logic for bootstrapping these correlators is similar to that for amplitudes:

- **SINGULARITIES:**

Correlators have characteristic singularities as a function of the external energies.

→ Analog of resonance singularities for amplitudes.

- **SYMMETRIES:**

These singularities are connected by causal time evolution, which is constrained by the symmetries of the bulk spacetime.

→ Analog of Lorentz symmetry for amplitudes.

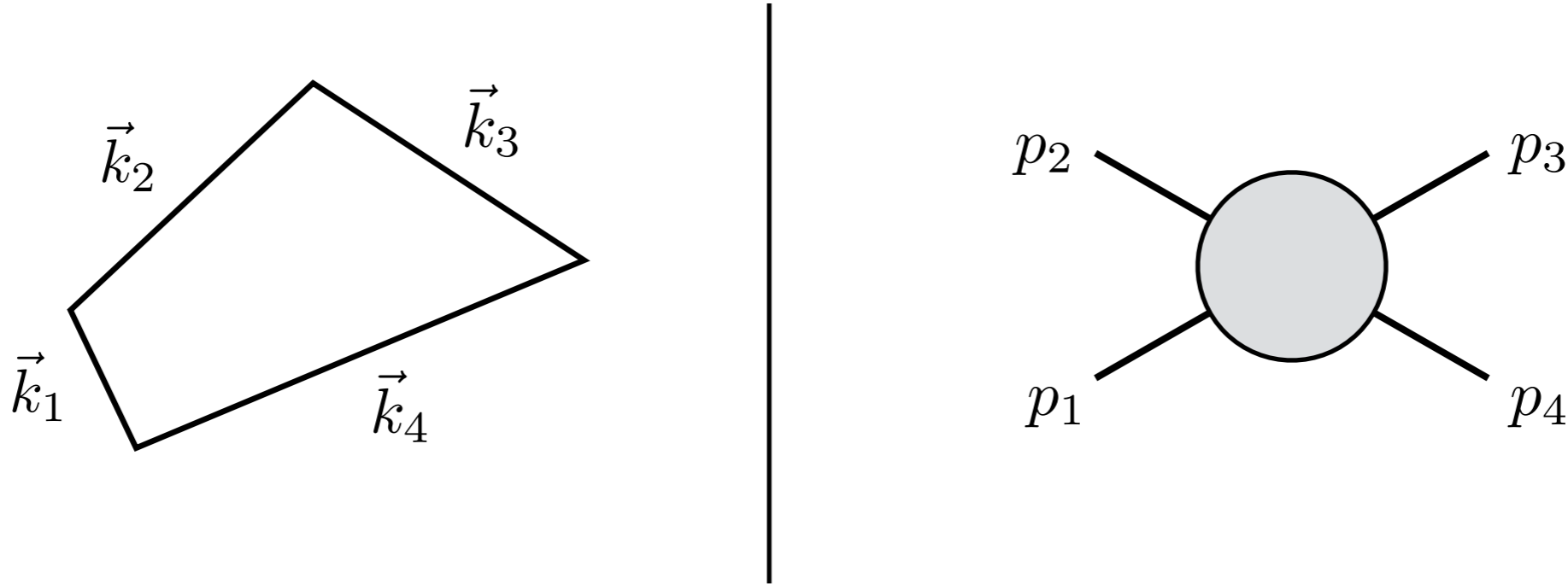
Arkani-Hamed, DB, Lee and Pimentel [2018]

DB, Duaso Pueyo, Joyce, Lee and Pimentel [2020]

DB, Chen, Duaso Pueyo, Joyce, Lee and Pimentel [2021]

Singularities

Correlators depend on the same external data as scattering amplitudes:



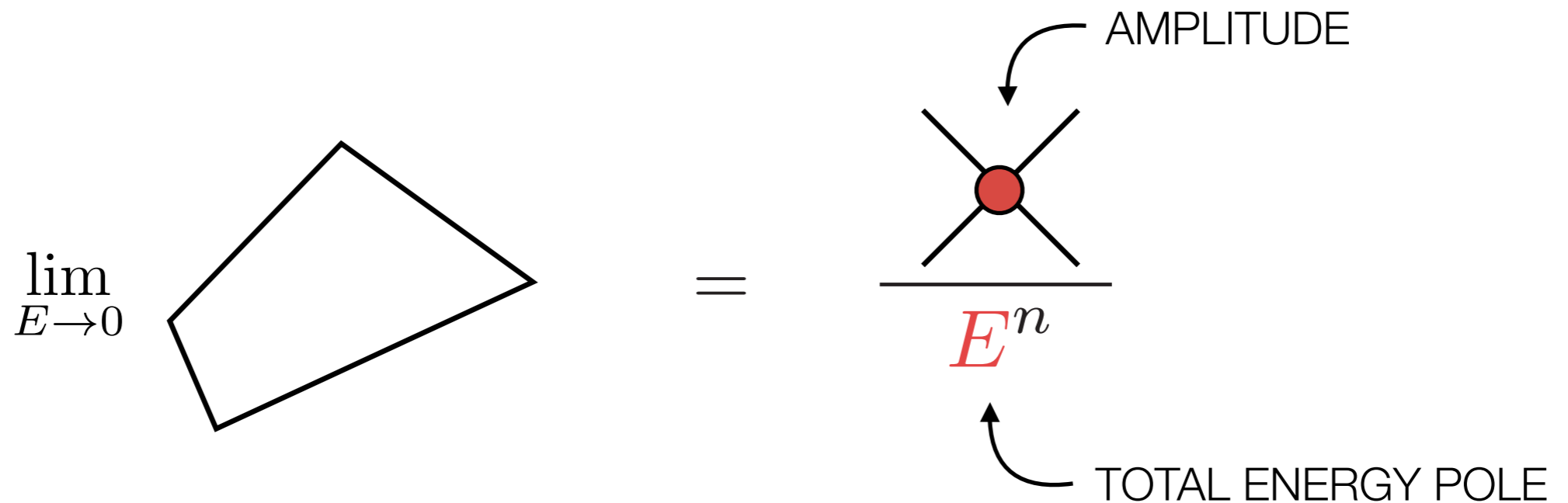
Two important differences:

- All energies must be positive: $E_n \equiv |\vec{k}_n| > 0$
- The total energy is **not** conserved in cosmology:

$$E \equiv |\vec{k}_1| + |\vec{k}_2| + |\vec{k}_3| + |\vec{k}_4| \neq 0$$

Singularities

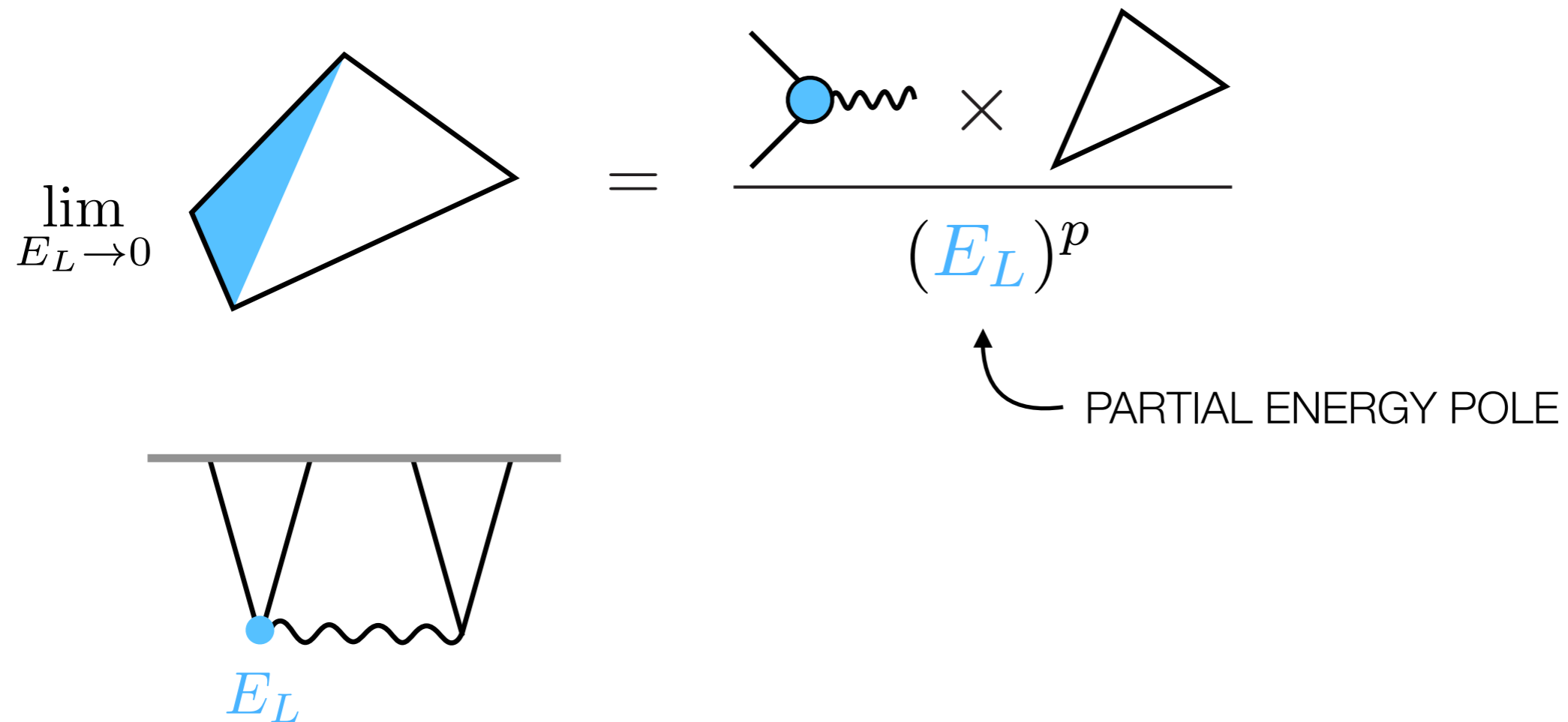
Something interesting happens in the limit of would-be energy conservation:



→ Amplitudes live inside correlators.

Singularities

Additional singularities arise when the energy of a subgraph is conserved:



Arkani-Hamed, Benincasa and Postnikov [2017]

Arkani-Hamed, DB, Lee and Pimentel [2018]

DB, Duaso Pueyo, Joyce, Lee and Pimentel [2020]

Singularities

Additional singularities arise when the energy of a subgraph is conserved:

The diagram illustrates a limit process in quantum field theory. On the left, a limit is taken as $E_L \rightarrow 0$ of a diagram consisting of a blue triangle attached to a larger orange triangle. This is equal to the product of a blue vertex with a wavy line and an orange triangle, all divided by $(E_L)^p$. An arrow labeled $E_R \rightarrow 0$ points to the right, where the wavy line and orange triangle are now connected to a single orange vertex, and the entire expression is divided by $(E_R)^q$.

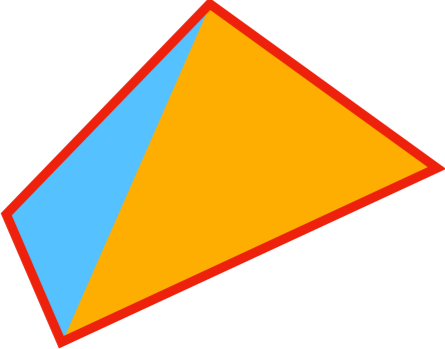
$$\lim_{E_L \rightarrow 0} \text{[Diagram]} = \frac{\text{[Diagram]} \times \text{[Diagram]}}{(E_L)^p} \xrightarrow{E_R \rightarrow 0} \frac{\text{[Diagram]}}{(E_R)^q}$$

→ Amplitudes are the building blocks of correlators.

Arkani-Hamed, Benincasa and Postnikov [2017]
Arkani-Hamed, DB, Lee and Pimentel [2018]
DB, Duaso Pueyo, Joyce, Lee and Pimentel [2020]

Singularities

At tree level, these are the only singularities of the correlator:

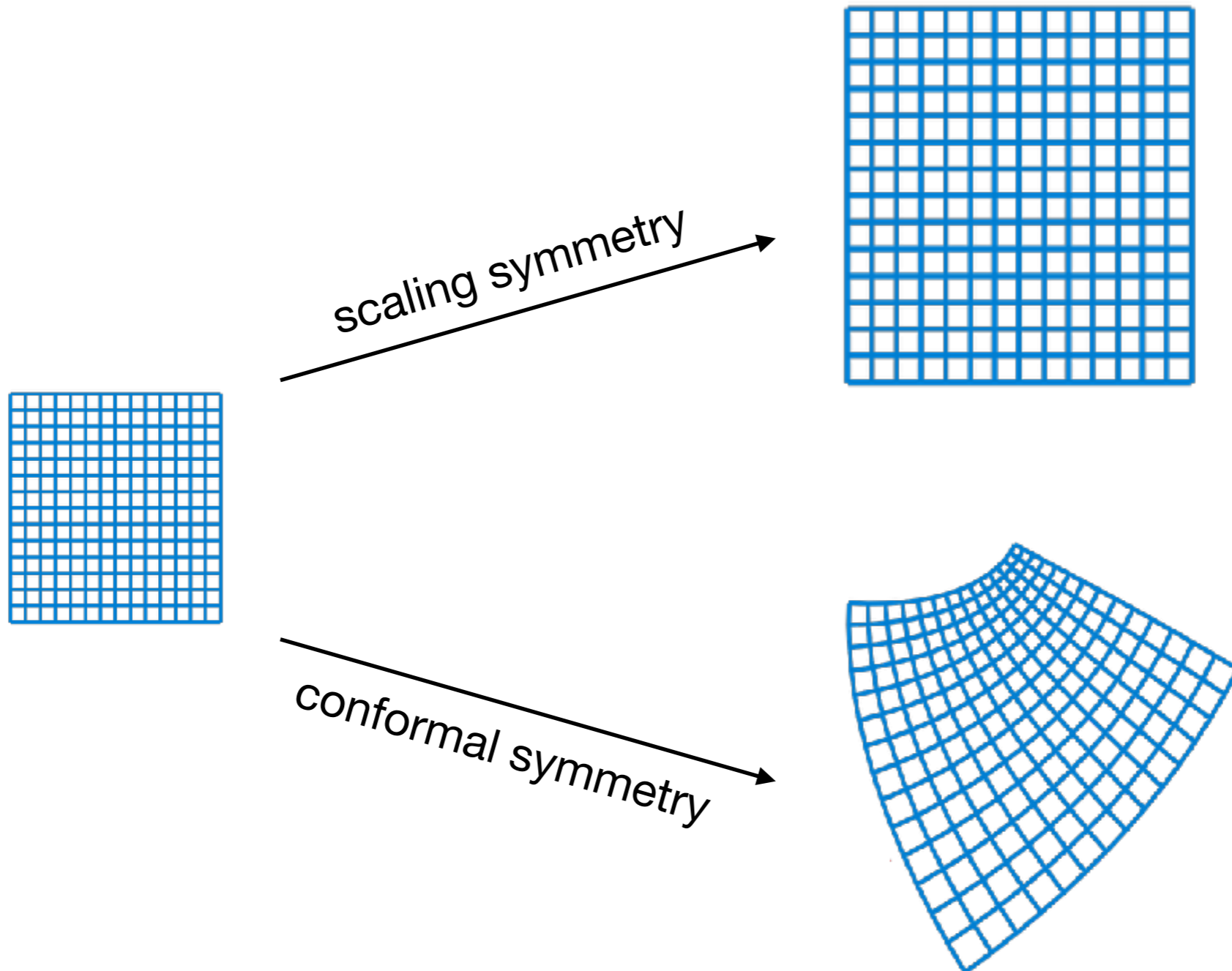

$$\begin{aligned} &= \frac{\text{Diagram 1}}{E^n} + \dots \\ &= \frac{\text{Diagram 2} \times \text{Diagram 3}}{(E_L)^p} + \dots \\ &= \frac{\text{Diagram 4} \times \text{Diagram 5}}{(E_R)^q} + \dots \end{aligned}$$

The diagrams are:
1. A wavy line connecting two vertices of a four-point tree-level diagram.
2. A vertex with a wavy line and two external lines, multiplied by an orange triangle.
3. A blue triangle.
4. A blue triangle multiplied by a vertex with a wavy line and two external lines.
5. A yellow triangle.

To determine the full correlator, we must connect these singular limits.

Symmetries

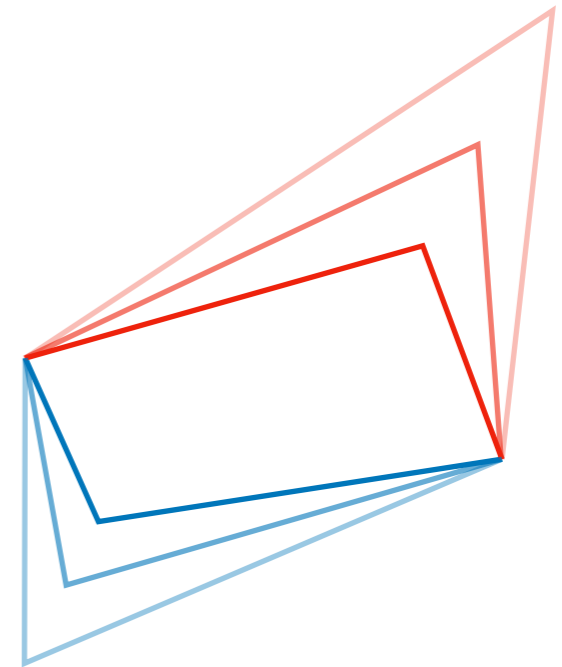
The isometries of the bulk de Sitter spacetime become rescaling symmetries of the correlators on the boundary:



Symmetries

These symmetries lead to a set of differential equations that control the amplitude of the correlations as we deform the quadrilateral

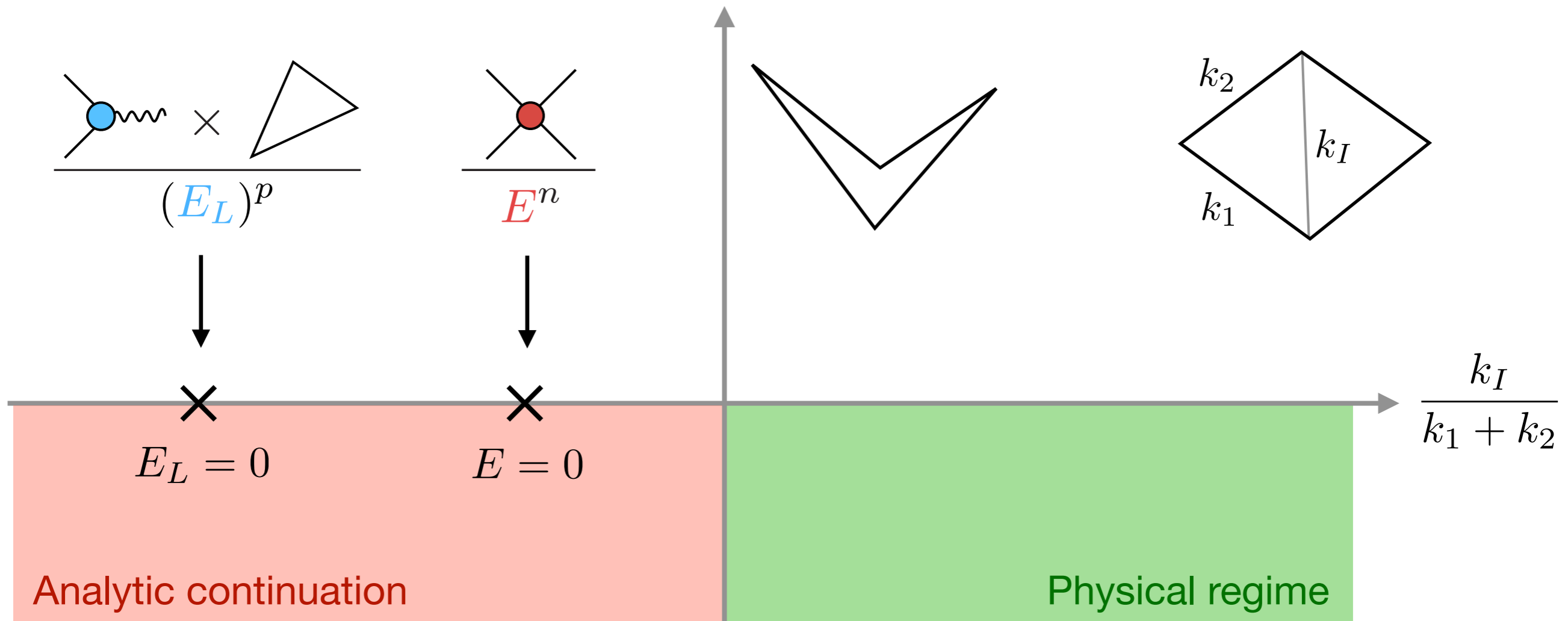
$$0 = \sum_{n=1}^4 \left[\vec{k}_n \partial_{\vec{k}_n}^2 - 2(\vec{k}_n \cdot \partial_{\vec{k}_n}) \partial_{\vec{k}_n} - 2\partial_{\vec{k}_n} \right]$$



The boundary conditions of this equation are the energy singularities.

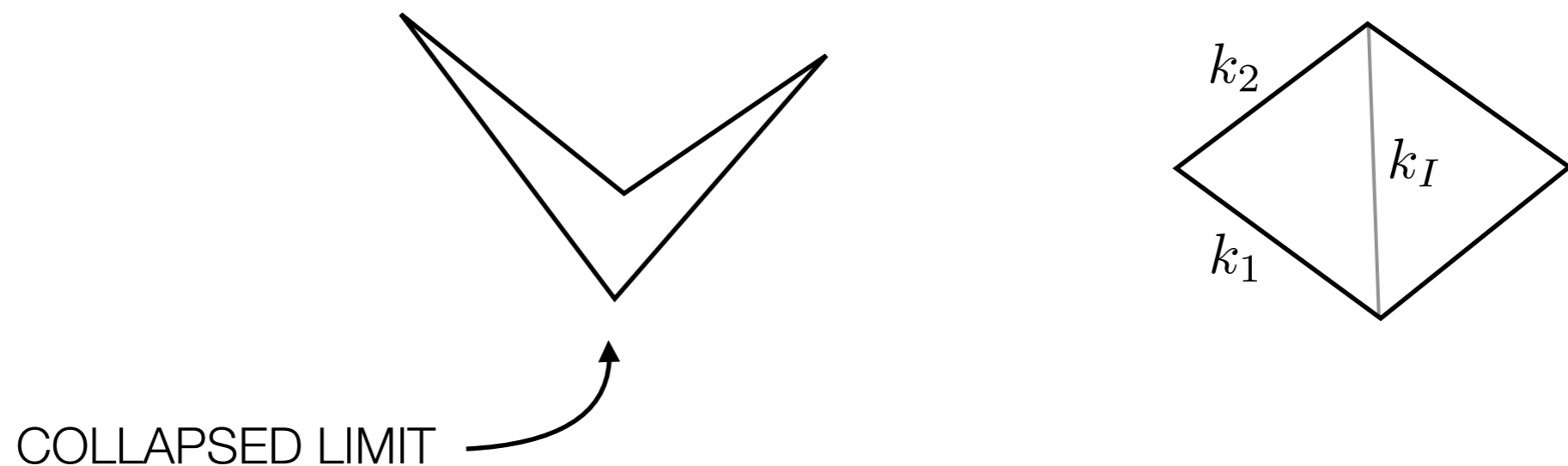
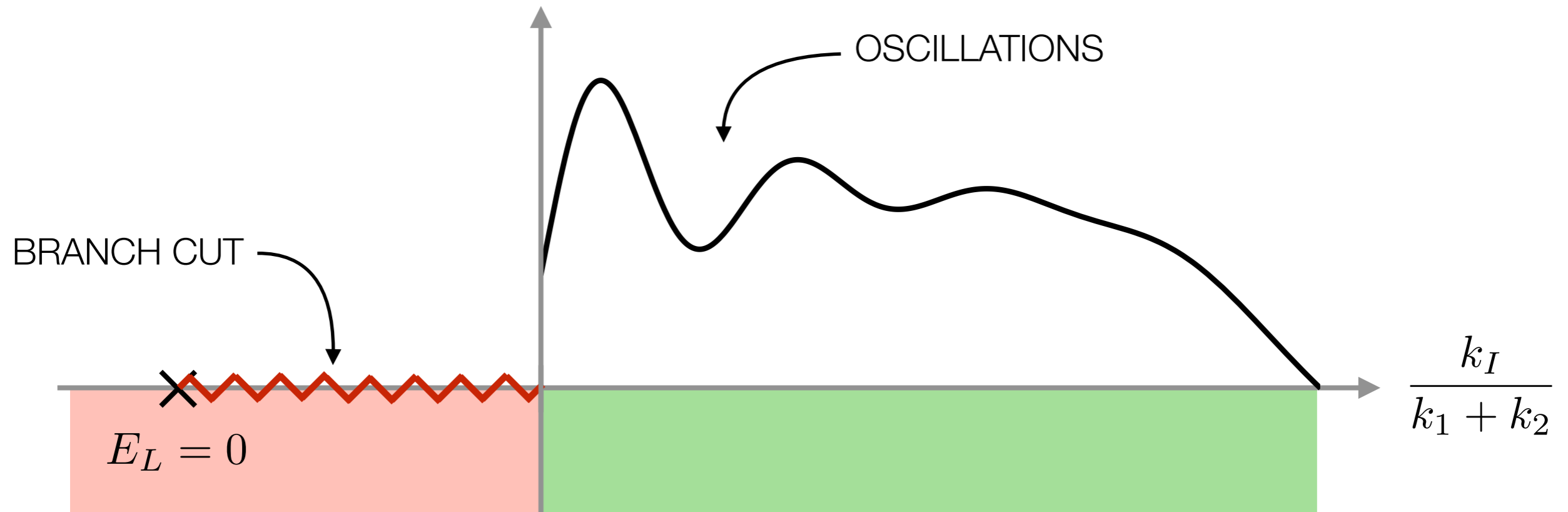
Symmetries

Symmetry relates the singularities in the unphysical region to the form of the correlator in the physical region:



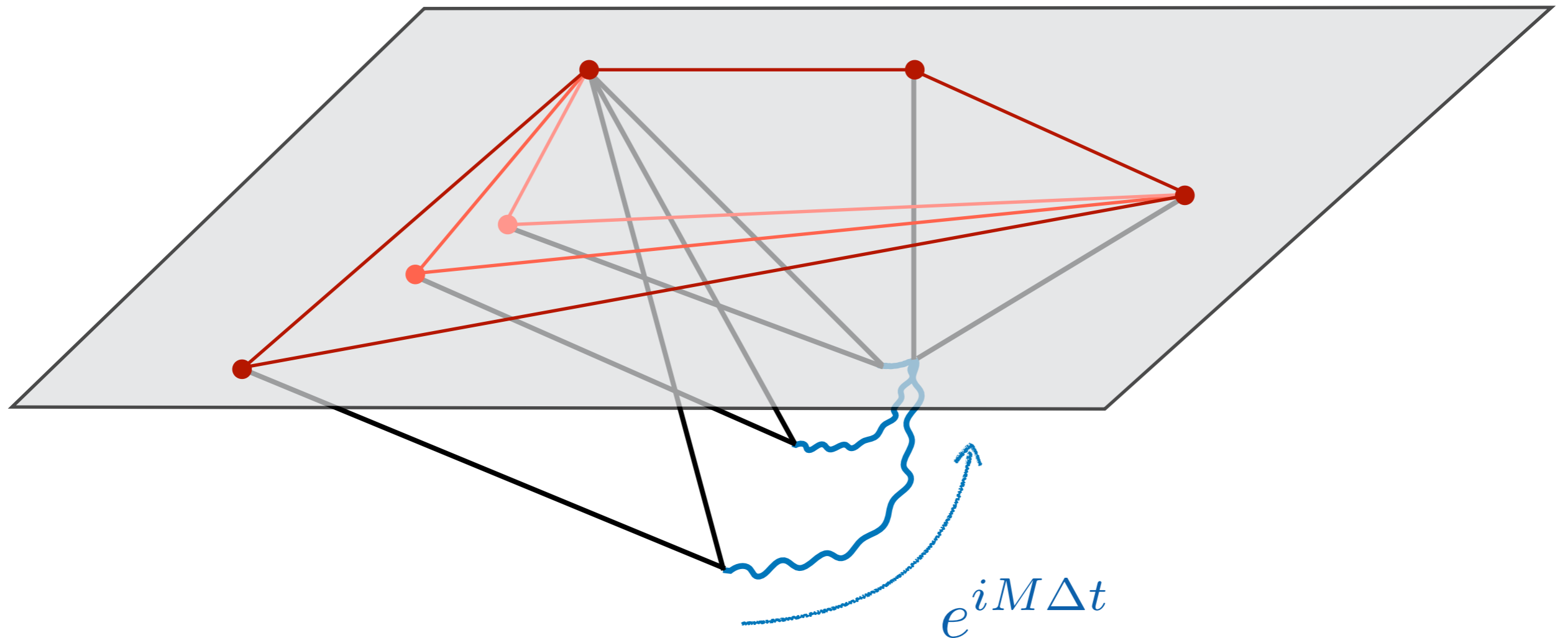
Particle Production

If the theory contains **massive particles**, then the correlator is forced to have an oscillatory feature in the physical regime:



Time Without Time

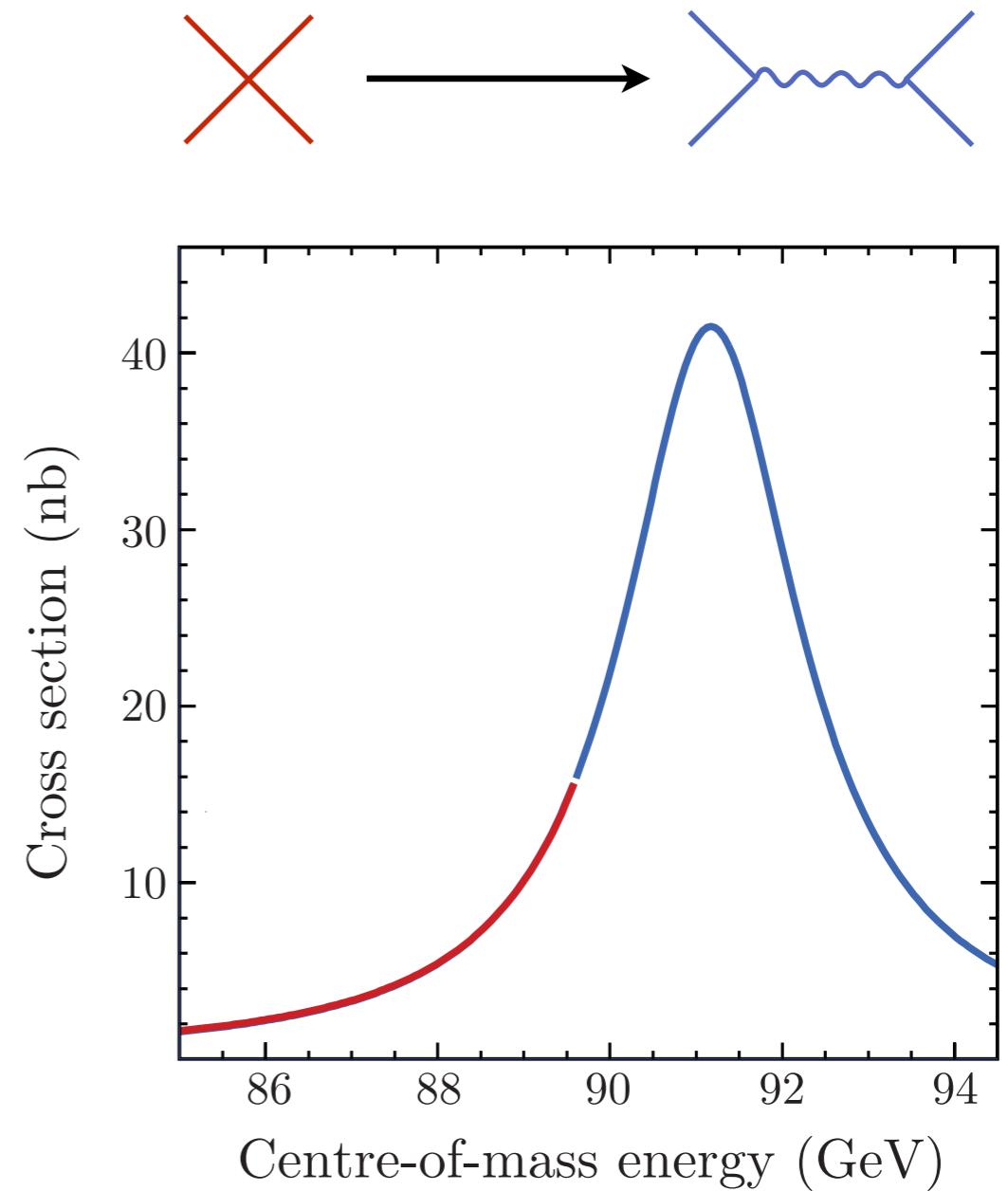
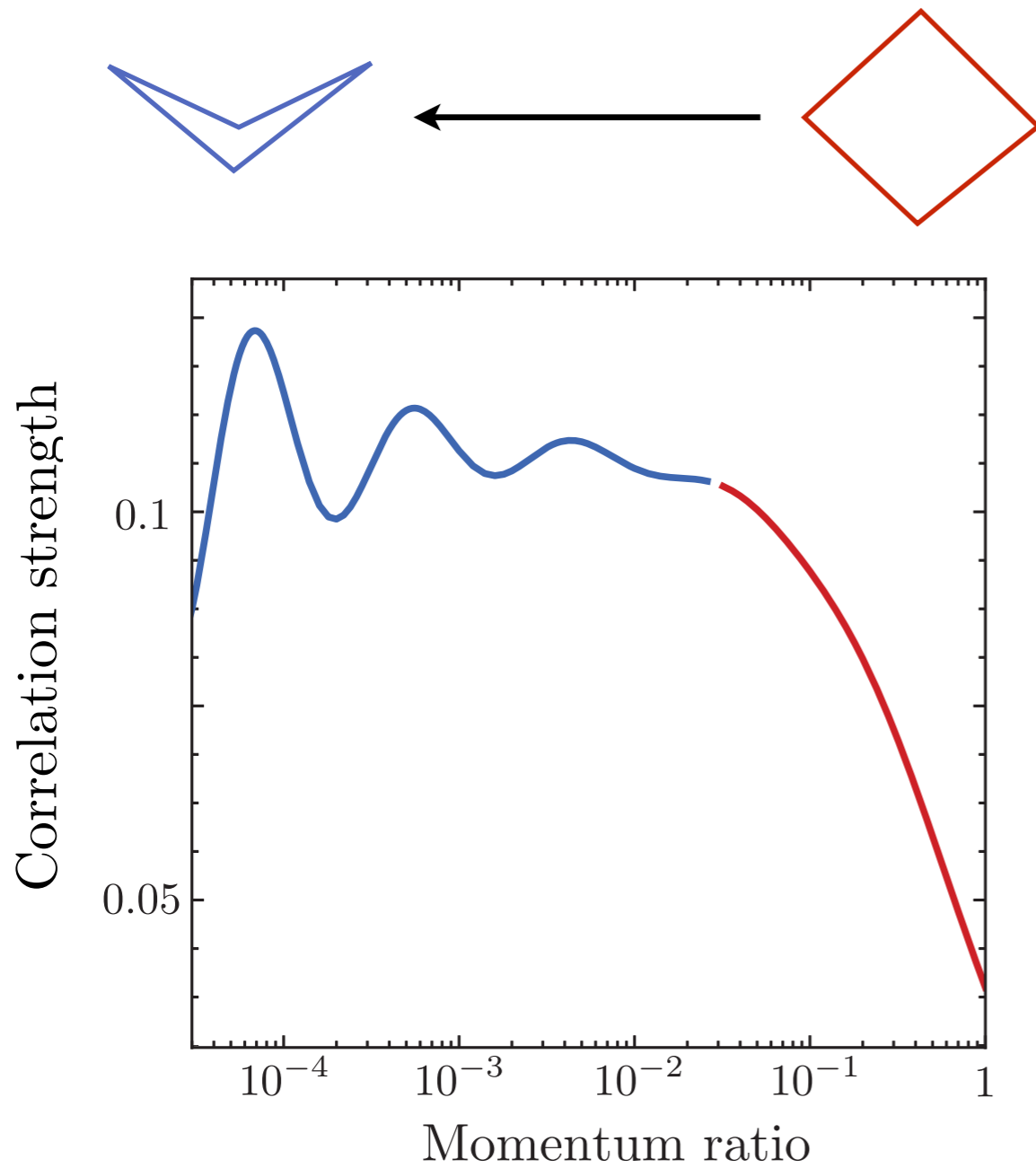
These oscillations reflect the evolution of the massive particles during inflation:



Time-dependent effects emerge in the solution of the time-independent bootstrap constraints.

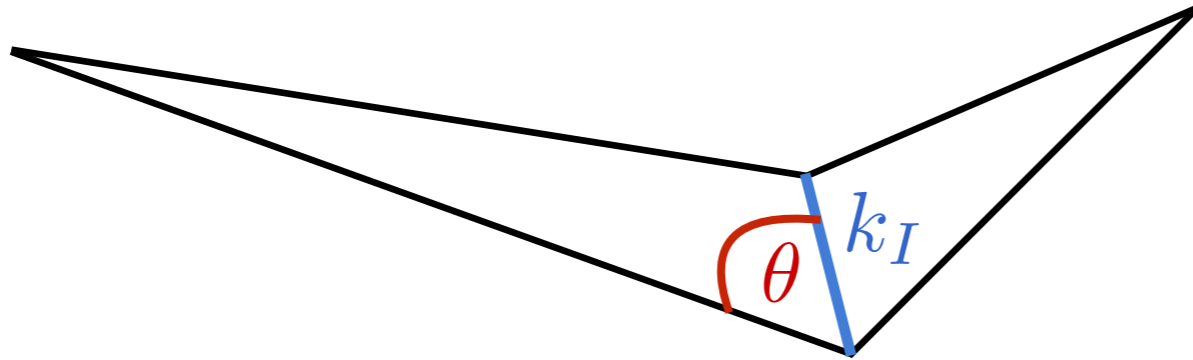
Cosmological Collider Physics

The oscillatory feature is the analog of a **resonance** in collider physics:



Particle Spectroscopy

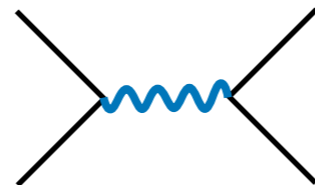
The frequency of the oscillations depends on the **mass** of the particles:



$$\propto \sin[M \log k_I] P_S(\cos \theta)$$

The angular dependence of the signal depends on the **spin** of the particles.

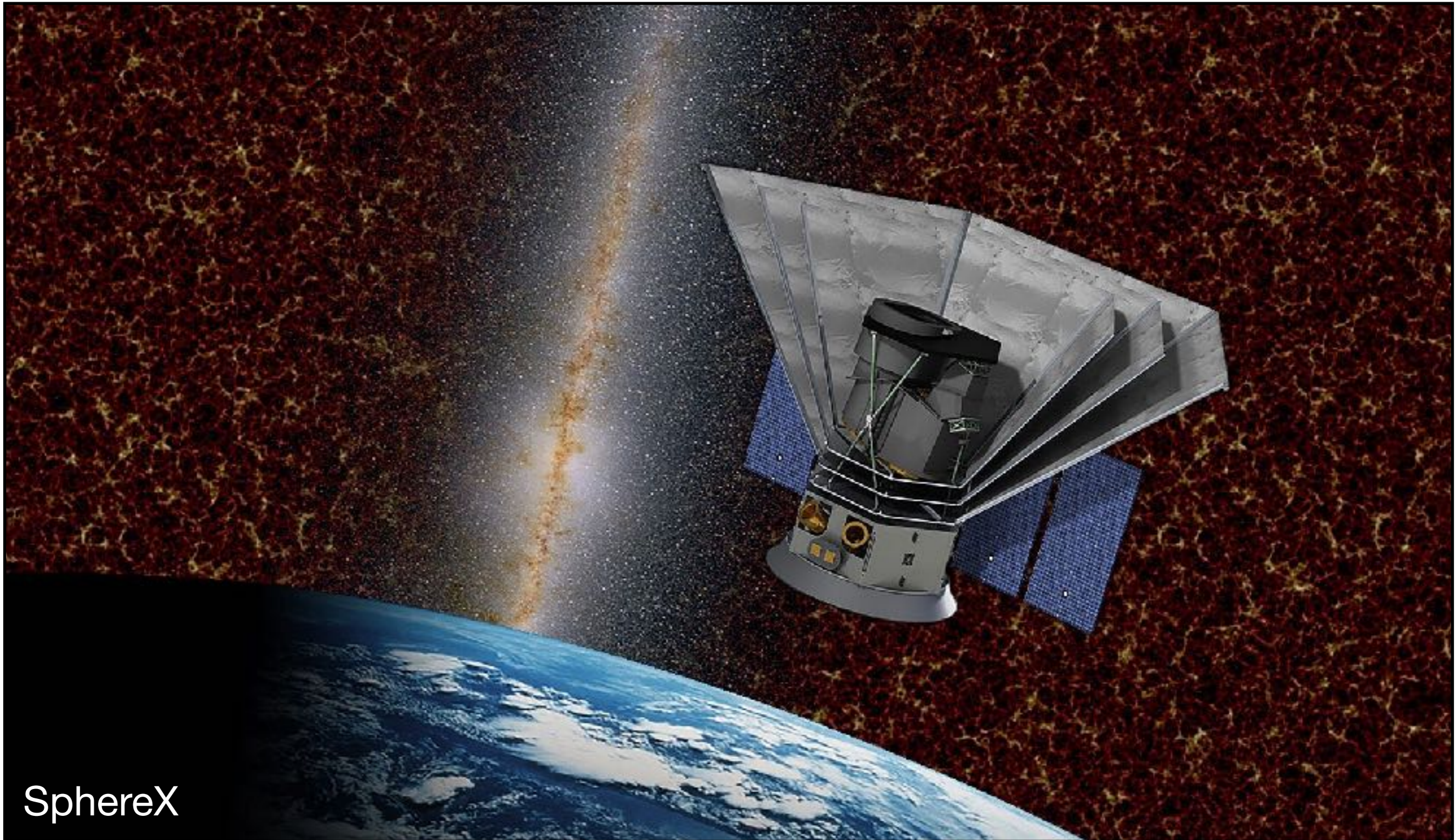
This is very similar to what we do in collider physics:

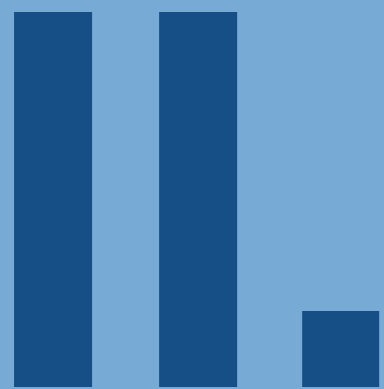


$$= \frac{g^2}{s - M^2} P_S(\cos \theta)$$

Observational Prospects

This signal will be an interesting target for future galaxy surveys:





Recent Progress

Recently, there has been significant progress in understanding the role of **unitarity** and **transmutation** in the cosmological bootstrap.

Goodhew, Jazayeri and Pajer [2020]

Cespedes, Davis and Melville [2020]

Meltzer and Sivaramakrishnan [2020]

Jazayeri, Pajer and Stefanyszyn [2021]

Melville and Pajer [2021]

Goodhew, Jazayeri, Lee and Pajer [2021]

DB, Chen, Duaso Pueyo, Joyce, Lee and Pimentel [2021]

Unitarity

An important constraint in quantum mechanics is

$$U^\dagger U = 1$$

*“What goes in,
must come out”*

For scattering amplitudes, this implies the **optical theorem**

$$\text{Im} \left(\text{Diagram} \right) = \sum_X \text{Diagram}_X \left(\text{Diagram} \right)^*$$

and the Cutkosky **cutting rules**.

→ What are the constraints of unitarity for cosmological correlators?

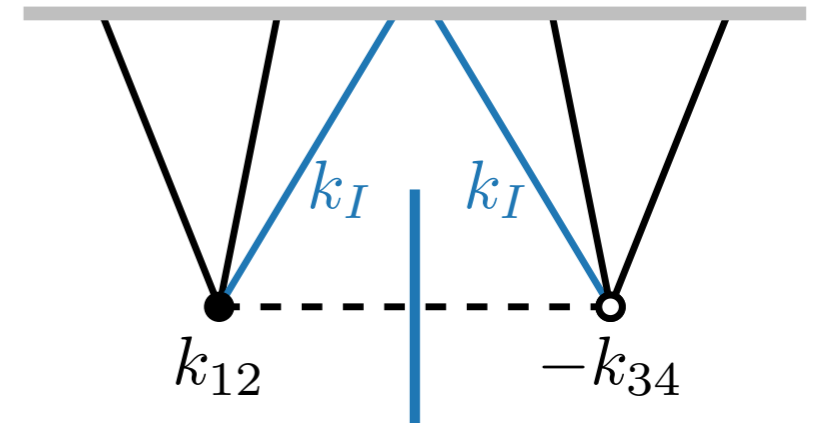
Unitarity

Goodhew, Jazayeri and Pajer derived the **cosmological optical theorem**

$$F_4(k_n) + F_4^*(-k_n) = \tilde{F}_3^L \otimes \tilde{F}_3^R$$

and associated **cutting rules**:

$$F_N(k_n) + F_N^*(-k_n) = - \sum_{\text{cuts}} F_N$$



- Unitarity enforces consistent factorization away from the energy poles.
- In some cases, this fixes the correlator without using de Sitter symmetry.

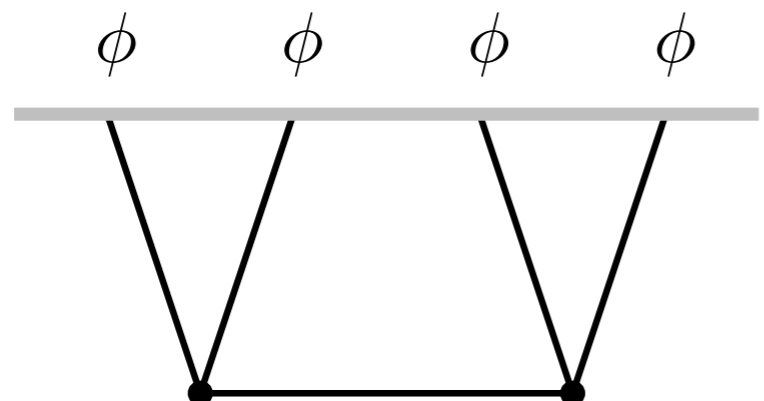
Goodhew, Jazayeri and Pajer [2020]

Jazayeri, Pajer and Stefanyszyn [2021]

DB, Chen, Duaso Pueyo, Joyce, Lee and Pimentel [2021]

Transmutation

Correlators in flat space are much easier to bootstrap than their counterparts in de Sitter space:



A Feynman diagram representing a correlator in flat space. It consists of a horizontal grey line at the top with four external legs labeled ϕ . Two vertices are connected by a horizontal line. Each vertex is connected to two external legs, forming two triangles that share the horizontal line between them.

$$= \frac{g^2}{EE_L E_R}$$

Only simple poles, with residues fixed by amplitudes.

Correlators in de Sitter space typically have higher-order poles, which makes them harder to bootstrap.

Transmutation

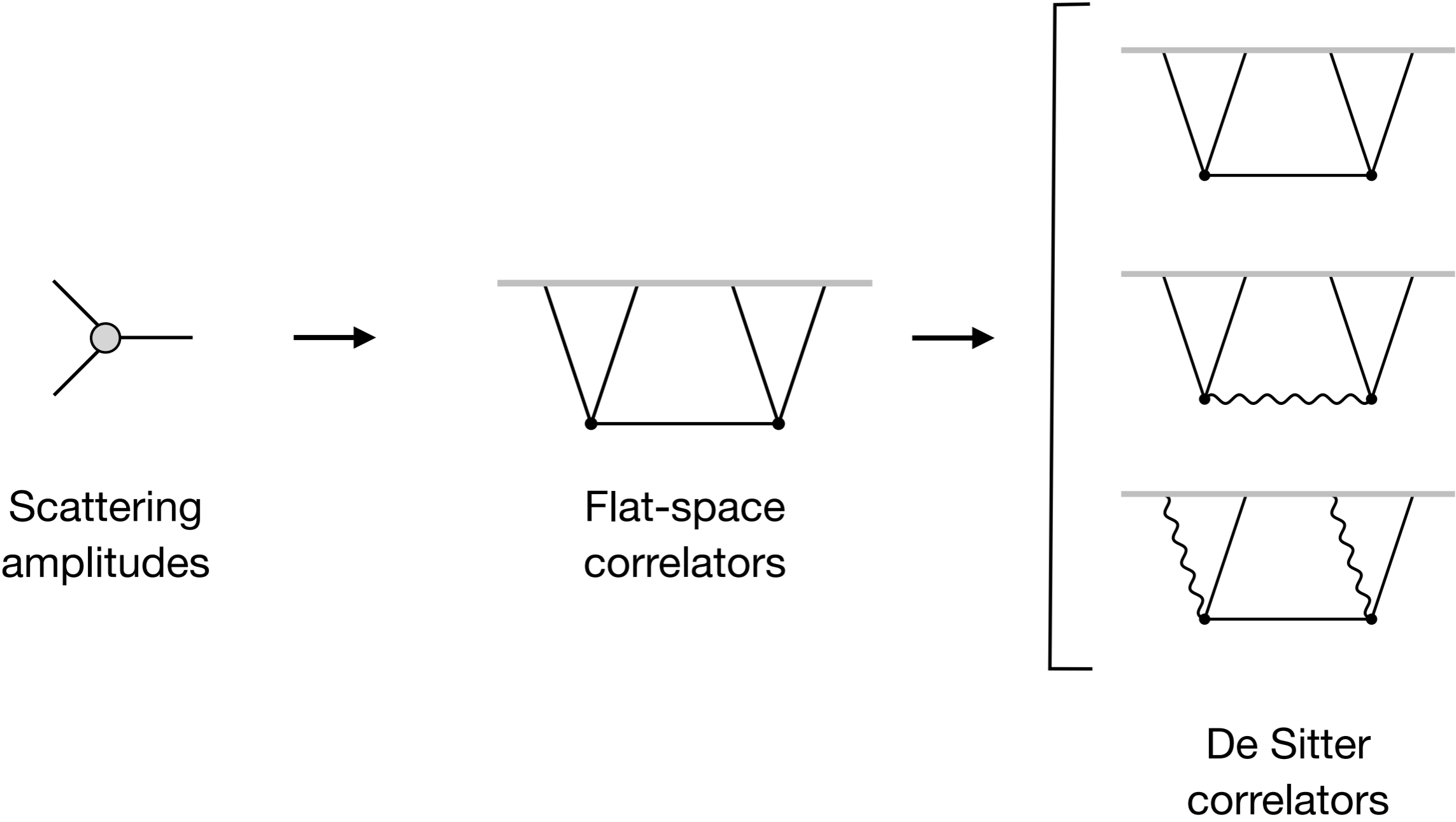
Recently, we showed that **complicated dS correlators** can be obtained by acting with **transmutation operators** on **simple flat-space seeds**:

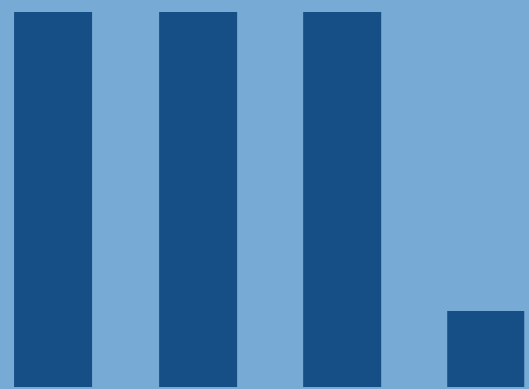
$$F_N^{(\text{dS})} = \mathcal{D} F_N^{(\text{flat})}$$

- Transmutation can also change the **spin** of the fields.
- It can also account for **symmetry breaking** interactions.

Charting the Space of dS Correlators

This has allowed us to bootstrap a large number of de Sitter correlators (using scattering amplitudes and flat-space correlators as input):





Future Directions

Cosmological correlation functions have a rich structure:

- Much of this structure is controlled by **singularities**.
- Behavior away from singularities captures **local bulk dynamics**.

The bootstrap approach has led to new **conceptual insights**:

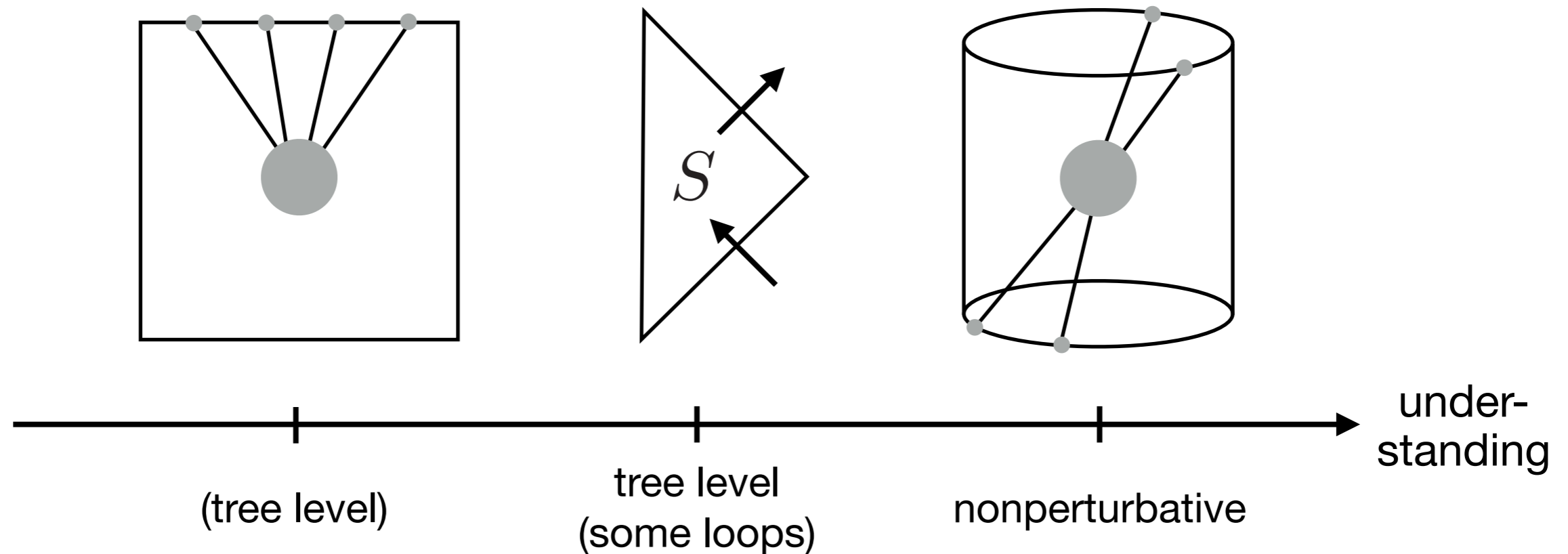
- Complex correlators can be derived from simpler seed functions:



It also has **practical applications**:

- Signatures of massive particles during inflation can be classified.

Only the beginning of a systematic exploration of cosmological correlators:



Important open questions are

- Are there **hidden structures** to be discovered?
- Can we go **beyond Feynman diagrams**?
- What is the analytic structure **beyond tree level**?
- What is the space of **UV-complete correlators**?
- What are its **observational signatures**?



NTU, Taiwan

Thank you for your attention!